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ASYMPTOTICS FOR THE RESOLVENT EQUATION ASSOCIATED TO THE GAME-THEORETIC p -LAPLACIAN

DIEGO BERTI AND ROLANDO MAGNANINI

ABSTRACT. We consider the (viscosity) solution u^ε of the elliptic equation $\varepsilon^2 \Delta_p^G u = u$ in a domain (not necessarily bounded), satisfying $u = 1$ on its boundary. Here, Δ_p^G is the *game-theoretic or normalized p -laplacian*. We derive asymptotic formulas for $\varepsilon \rightarrow 0^+$ involving the values of u^ε , in the spirit of Varadhan's work [Va], and its q -mean on balls touching the boundary, thus generalizing that obtained in [MS1] for $p = q = 2$. As in a related parabolic problem, investigated in [BM], we link the relevant asymptotic behavior to the geometry of the domain.

1. INTRODUCTION

In this paper we consider (viscosity) solutions u^ε of the following one-parameter family of problems:

$$(1.1) \quad u - \varepsilon^2 \Delta_p^G u = 0 \quad \text{in } \Omega,$$

$$(1.2) \quad u = 1 \quad \text{on } \Gamma,$$

where Ω is a domain in \mathbb{R}^N , $N \geq 2$, Γ is its boundary and $\varepsilon > 0$. The operator Δ_p^G in (1.1) is the *game-theoretic (or normalized or homogeneous) p -laplacian*, that is formally defined by

$$(1.3) \quad \Delta_p^G u = \frac{1}{p} |\nabla u|^{2-p} \operatorname{div} \{ |\nabla u|^{p-2} \nabla u \}$$

if $p \in [1, \infty)$, and by its limit as $p \rightarrow \infty$,

$$(1.4) \quad \Delta_\infty^G u = \frac{\langle \nabla^2 u \nabla u, \nabla u \rangle}{|\nabla u|^2},$$

in the extremal case $p = \infty$. Note that $\Delta_2^G = \frac{1}{2} \Delta$, where Δ is the classical (linear) Laplace operator. Otherwise, Δ_p^G is a non-linear operator in non-divergence form, with possibly discontinuous coefficients at points in which the gradient ∇u is zero. Nevertheless, differently from the case of the variational p -laplacian $\Delta_p u = \operatorname{div} \{ |\nabla u|^{p-2} \nabla u \}$, Δ_p^G inherits from Δ , for any $p \in (1, \infty]$, its 1-homogeneity and structure of second-order uniformly elliptic operator. The lack of continuity of the coefficients at the critical points of u and the non-divergence form of Δ_p^G make the use of the theory of *viscosity solutions* preferable to that of weak (variational) solutions. On this issue, besides the classical treatises [CGG] and [CIL], we mention a list of more

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recent references that specifically concern Δ_p^G : [AP], [APR], [BG1], [BG2], [Do], [JK], [KH], [KKK], [MPR] and [MPR2].

Problems involving Δ_p^G appear in a variety of applications. First of all, as the name evokes, the solution has a natural game-theoretic interpretation: the solution of problem (1.1)-(1.2) represents the limiting value of a two-player, zero-sum game, called *tug-of-war* (for $p = \infty$) or *tug-of-war with noise* (for $1 < p < \infty$), see for instance [PS],[PSSW], [MPR] and [MPR2].

In recent years there has been a growing interest for equations involving Δ_p^G in relation to numerical methods for *image enhancement or restoration* (see [Do] and [BSA]). Typically, for a possibly corrupted image represented by a function u_0 , it is considered an evolutionary process based on Δ_p^G with initial data u_0 and homogeneous Neumann boundary conditions. As explained in [Do], the different choice of p affects in which direction the brightness evolves; the 1-homogeneity of Δ_p^G ensures that such an evolution does not depend on the brightness of the image. The relation between the solution of the parabolic problem and the parametrized elliptic equation (1.1) is examined in [BSA] (where the backward version of the same problem is considered) for the classical p -laplacian, and can be extended to the case of Δ_p^G in hand.

More classical applications were considered in the linear case ($p = 2$). Indeed, in the context of *large deviations theory* (see [FW], [Va], [EI]), random differential equations with small noise intensities are considered and the profile for small values of ε of the solutions of (1.1)-(1.2) is related to the behaviour of the exit time of a certain stochastic process. In geometrical terms, that behavior is encoded in the following formula:

$$(1.5) \quad \lim_{\varepsilon \rightarrow 0^+} \varepsilon \log u^\varepsilon(x) = -\sqrt{2} d_\Gamma(x),$$

uniformly on a bounded domain $\bar{\Omega}$, where $d_\Gamma(x)$ denotes the distance of a point $x \in \Omega$ to Γ .

Further evidence of the influence of geometry on the asymptotic behavior of solutions of elliptic and parabolic equations for small values of the relevant parameter was given by the second author of this note and S. Sakaguchi in a series of papers both in the linear case ([MS1], [MS3], [MS5], [MS7], [MM]) and in certain non-linear contexts ([MS2], [MS4], [MS6], [Sa1]), concerning both initial-boundary value problems ([MS1], [MS2], [MS4]) and initial-value problems ([MPeS], [MPrS]), and even for two-phase problems ([Sa2], [Sa3], [CMS]). For instance, in [MS1, Theorem 2.3], in the case $p = 2$ for problem (1.1)-(1.2), the following formula involving the mean value of u^ε was established:

$$(1.6) \quad \lim_{\varepsilon \rightarrow 0^+} \left(\frac{R}{\varepsilon} \right)^{\frac{N-1}{2}} \int_{\partial B_R(x)} u^\varepsilon(y) dS_y = \frac{c_N}{\sqrt{\Pi_\Gamma(y_x)}};$$

c_N is a constant that can be derived from [MS1, Theorem 2.3]. Here, for a given $x \in \Omega$ with $d_\Gamma(x) = R$, $B_R(x)$ is assumed to be a ball contained in Ω and such that $(\mathbb{R}^N \setminus \Omega) \cap \partial B_R(x) = \{y_x\}$, $\kappa_1, \dots, \kappa_{N-1}$ denote the principal

curvatures (with respect to the inward normal) of Γ at points of Γ , and

$$(1.7) \quad \Pi_\Gamma(y_x) = \prod_{j=1}^{N-1} \left[1 - R \kappa_j(y_x) \right].$$

In this paper, we shall derive asymptotic formulas similar to (1.5) and (1.6) for the solution u^ε of (1.1)-(1.2), when $p \neq 2$.

First of all, in Section 2 we shall prove for $p \in (1, \infty]$ the validity of the following analog of (1.5):

$$(1.8) \quad \lim_{\varepsilon \rightarrow 0^+} \varepsilon \log [u^\varepsilon(x)] = -\sqrt{p'} d_\Gamma(x), \quad \text{for any } x \in \bar{\Omega}.$$

Here p' denotes the conjugate exponent to p (i.e. $1/p + 1/p' = 1$) and Ω is required to merely satisfy the regularity assumption: $\Gamma = \partial(\mathbb{R}^N \setminus \bar{\Omega})$.

Asymptotics (1.8) is obtained, first, by considering the solutions of (1.1)-(1.2) in the two basic radial cases in which Ω is either a ball or the exterior of its closure (see Subsection 2.1) and, secondly, by employing them as barriers, by virtue of the comparison principle and taking advantage of the fact that in the radial cases the solutions of (1.1)-(1.2) are obtained explicitly. In fact, equation (1.1) can be reduced to a *linear* ordinary differential equation when u^ε is radially symmetric. The use of radial barriers, that is natural in the case of a C^2 domain, will also be adapted to a quite general case.

In Section 2 we shall also investigate uniform estimates related to (1.8). An interesting feature of such estimates is that their quality seems to depend on the parameter p . In fact, while for a quite large class of regular domains (that we call $C^{0,\omega}$ domains and these include the scale of those with Hölder regularity) the following formula always holds uniformly on any compact subset of $\bar{\Omega}$

$$(1.9) \quad \varepsilon \log(u^\varepsilon) + \sqrt{p'} d_\Gamma = O(\varepsilon |\log \psi(\varepsilon)|) \quad \text{as } \varepsilon \rightarrow 0^+$$

(here $\psi(\varepsilon) = \psi_\omega(\varepsilon)$ vanishes with ε depending on the modulus of continuity ω), we see that, the right-hand side of (1.9) improves to $O(\varepsilon \log |\log \psi(\varepsilon)|)$ if $p = N$, $O(\varepsilon \log \varepsilon)$ for $p \in (N, \infty)$, and $O(\varepsilon)$ if $p = \infty$ (see Theorem 2.6).

In Section 3, we shall derive a formula that greatly generalizes (1.6): in fact, in Theorem 3.5 it is shown that for $1 < p \leq \infty$ and $1 < q < \infty$ it holds that

$$(1.10) \quad \lim_{\varepsilon \rightarrow 0^+} \left(\frac{R}{\varepsilon} \right)^{\frac{N+1}{2(q-1)}} \mu_{q,\varepsilon}(x) = \frac{c_{N,q}}{\left\{ (p')^{\frac{N+1}{2}} \Pi_\Gamma(y_x) \right\}^{\frac{1}{2(q-1)}}}$$

(the exact value of the constant $c_{N,q}$ is given in Theorem 3.5). Here, $\mu_{q,\varepsilon}(x)$ is the q -mean of u^ε on a ball $B_R(x)$ touching Γ at only one point y_x ; $\mu_{q,\varepsilon}(x)$ is the unique $\mu \in \mathbb{R}$ such that

$$(1.11) \quad \|u^\varepsilon - \mu\|_{L^q(B_R(x))} \leq \|u^\varepsilon - \lambda\|_{L^q(B_R(x))} \quad \text{for all } \lambda \in \mathbb{R};$$

this can be defined for every $q \in [1, \infty]$ (see [IMW]). Notice that $\mu_{2,\varepsilon}(x)$ is the standard mean value of u^ε on $B_R(x)$.

If $q = \infty$ we simply have that $\mu_{\infty,\varepsilon}(x) \rightarrow 1/2$, which does not give any geometrical information.

Also to obtain (1.10), we first show that an appropriate formula holds for suitable spherical barriers for u^ε and hence we use the fact that q -means are monotonic with respect to (almost everywhere) partial order of functions on $B_R(x)$.

Formulas (1.8) and (1.10) are elliptic versions of similar formulas obtained in [BM] for the viscosity solution $u = u(x, t)$ of the parabolic equation

$$u_t - \Delta_p^G u = 0 \quad \text{in } \Omega \times (0, \infty),$$

subject to the boundary-initial data:

$$u = 0 \quad \text{on } \Omega \times \{0\} \quad \text{and} \quad u = 1 \quad \text{on } \Gamma \times (0, \infty).$$

Indeed, in [BM] we obtained:

$$\lim_{t \rightarrow 0^+} 4t \log u(x, t) = -p' d_\Gamma(x)^2,$$

and

$$\lim_{t \rightarrow 0^+} \left(\frac{R^2}{t} \right)^{\frac{N+1}{4(q-1)}} \mu_{q,t}(x) = \frac{C_{N,q}}{\left\{ (p')^{\frac{N+1}{2}} \Pi_\Gamma(y_x) \right\}^{\frac{1}{2(q-1)}}},$$

for some explicit constant $C_{N,q}$, where $\mu_{q,t}(x)$ is the q -mean of $u(\cdot, t)$ on $B_R(x)$. It should be noticed that in the case $p = 2$, due to the linearity of Δ , the last formula and (1.10) can be obtained from one another, since $u^\varepsilon(x)$ and $u(x, t)$ are related by a Laplace transformation. When $p \neq 2$, this is no longer possible and the elliptic and parabolic cases must be treated separately. Moreover, in the elliptic case, due to the availability of explicit barriers, we obtain the more accurate uniform estimate (1.9) (compare with [BM, Eq. (2.18)]).

We conclude this introduction by mentioning that the linearity of Δ was used in [MS1] to derive radial symmetry of the so-called *stationary isothermic surfaces*, that is those level surfaces of the temperature which are invariant in time. In fact, it was shown that the mean value $\mu_{q,t}(x)$ or $\mu_{q,\varepsilon}(x)$ does not depend on x if this lies on a stationary isothermic surface, and hence, for instance, (1.6) gives that Π_Γ must be constant on Γ , so that radial symmetry then ensues from Alexandrov's Soap Bubble Theorem for Weingarten surfaces (see [Al]). For $p \neq 2$, this approach is no longer possible. However, an approach based on the method of moving planes, as considered in [MS4] and [CiMS], may still be possible.

2. LARGE-DEVIATIONS ASYMPTOTICS

The proof of the asymptotic formulas (1.8) and (1.9) will be carried out in two steps.

2.1. Asymptotics in the radial cases. In this subsection we collect the relevant properties of the explicit solutions for problem (1.1)-(1.2) when Ω is a ball or the complement of a ball. We shall denote by B_R the ball with radius R centered at the origin.

In [BM][Lemma 2.1], we have shown that the (viscosity) solution of (1.1)-(1.2) in B_R is given by

$$(2.1) \quad u^\varepsilon(x) = \begin{cases} \frac{\int_0^\pi e^{\sqrt{p'} \frac{|x|}{\varepsilon} \cos \theta} (\sin \theta)^\alpha d\theta}{\int_0^\pi e^{\sqrt{p'} \frac{R}{\varepsilon} \cos \theta} (\sin \theta)^\alpha d\theta} & \text{if } 1 < p < \infty, \\ \frac{\cosh(|x|/\varepsilon)}{\cosh(R/\varepsilon)} & \text{if } p = \infty, \end{cases}$$

for $x \in \overline{B_R}$; here

$$(2.2) \quad \alpha = \frac{N-p}{p-1} = -1 + \frac{N-1}{p-1}.$$

For the reader's convenience, we observe that (2.1) can be verified, away from the origin of B_R , by direct differentiation. A technical lemma ([BM, Lemma A.6]) then makes sure that (2.1) is the viscosity solution of (1.1) in B_R , that is, also at its unique critical point. For $1 < p < \infty$, one has to take the first and second derivatives of the function in (2.1) and hence compare them, after integrating by parts the resulting formula for the first derivative. Alternatively, one can manipulate formulas for the *modified Bessel functions* (see [AS, Section 9.6; formulas 9.6.10 and 9.6.18]).

Moreover, formula (1.8) holds with $\Omega = B_R$ and the next lemma gives a quantitative estimate of the convergence in (1.8).

Lemma 2.1 (Uniform asymptotics in a ball). *Let $1 < p \leq \infty$ and let u^ε be the solution of (1.1)-(1.2) in B_R .*

Then

$$(2.3) \quad \varepsilon \log u^\varepsilon + d_\Gamma = \begin{cases} O(\varepsilon |\log \varepsilon|) & \text{if } 1 < p < \infty, \\ O(\varepsilon) & \text{if } p = \infty, \end{cases}$$

uniformly on $\overline{B_R}$ as $\varepsilon \rightarrow 0^+$.

Proof. The case $p = \infty$ follows at once, since (2.1) gives:

$$\varepsilon \log \{u^\varepsilon(x)\} + d_\Gamma(x) = \varepsilon \log \left[\frac{1 + e^{-2\frac{|x|}{\varepsilon}}}{1 + e^{-2\frac{R}{\varepsilon}}} \right].$$

If $1 < p < \infty$, by (2.1) we have that

$$\varepsilon \log \{u^\varepsilon(x)\} + \sqrt{p'} d_\Gamma(x) = \varepsilon \log \left[\frac{\int_0^\pi e^{-\sqrt{p'}(1-\cos \theta) \frac{|x|}{\varepsilon}} (\sin \theta)^\alpha d\theta}{\int_0^\pi e^{-\sqrt{p'}(1-\cos \theta) \frac{R}{\varepsilon}} (\sin \theta)^\alpha d\theta} \right]$$

and the right-hand side is decreasing in $|x|$, so that

$$0 \leq \varepsilon \log \{u^\varepsilon(x)\} + \sqrt{p'} d_\Gamma(x) \leq \varepsilon \log \left[\frac{\int_0^\pi (\sin \theta)^\alpha d\theta}{\int_0^\pi e^{-\sqrt{p'}(1-\cos \theta) \frac{R}{\varepsilon}} (\sin \theta)^\alpha d\theta} \right].$$

This formula gives (2.3), since we have that

$$\int_0^\pi e^{-\sqrt{p'}(1-\cos\theta)\frac{R}{\varepsilon}} (\sin\theta)^\alpha d\theta = 2^{\frac{\alpha-1}{2}} \Gamma\left(\frac{\alpha+1}{2}\right) \left(\frac{R\sqrt{p'}}{\varepsilon}\right)^{-\frac{\alpha+1}{2}} [1 + O(\varepsilon)]$$

as $\varepsilon \rightarrow 0^+$, by using Lemma A.1. \square

Next, we consider the complement of a ball.

Lemma 2.2 (Uniform asymptotics in the complement of a ball). *Set $1 < p \leq \infty$, $\Omega = \mathbb{R}^N \setminus \overline{B_R}$, and let u^ε be the bounded solution of (1.1)-(1.2).*

Then we have that

$$(2.4) \quad u^\varepsilon(x) = \begin{cases} \frac{\int_0^\infty e^{-\sqrt{p'} \cosh\theta \frac{|x|}{\varepsilon}} (\sinh\theta)^\alpha d\theta}{\int_0^\infty e^{-\sqrt{p'} \cosh\theta \frac{R}{\varepsilon}} (\sinh\theta)^\alpha d\theta} & \text{if } 1 < p < \infty, \\ e^{-\frac{|x|-R}{\varepsilon}} & \text{if } p = \infty. \end{cases}$$

In particular,

$$(2.5) \quad \varepsilon \log \{u^\varepsilon(x)\} + \sqrt{p'} d_\Gamma(x) = O(\varepsilon) \quad \text{as } \varepsilon \rightarrow 0^+,$$

uniformly on every compact subset of $\overline{\Omega}$.

Proof. Verifying (1.2) is immediate, as well as that (1.1) is satisfied for $p = \infty$. As for $1 < p < \infty$, similarly to the case of the ball B_R , (1.1) is verified by taking derivatives and integrating by parts. Differently from the case of the ball, here u^ε does not have critical points. Thus, it is a classical and hence viscosity solution of (1.1).

Now, notice that $d_\Gamma(x) = |x| - R$ for $|x| \geq R$. If $p = \infty$, (2.5) holds exactly as

$$\varepsilon \log \{u^\varepsilon(x)\} + d_\Gamma(x) \equiv 0.$$

If $1 < p < \infty$, we write that

$$\varepsilon \log \{u^\varepsilon(x)\} + \sqrt{p'} d_\Gamma(x) = \frac{\int_0^\infty e^{-\sqrt{p'} \frac{|x|}{\varepsilon} (\cosh\theta - 1)} (\sinh\theta)^\alpha d\theta}{\int_0^\infty e^{-\sqrt{p'} \frac{R}{\varepsilon} (\cosh\theta - 1)} (\sinh\theta)^\alpha d\theta}$$

and hence, by monotonicity, we have that

$$\varepsilon \log \left\{ \frac{\int_0^\infty e^{-\sqrt{p'} \frac{R'}{\varepsilon} (\cosh\theta - 1)} (\sinh\theta)^\alpha d\theta}{\int_0^\infty e^{-\sqrt{p'} \frac{R}{\varepsilon} (\cosh\theta - 1)} (\sinh\theta)^\alpha d\theta} \right\} \leq \varepsilon \log \{u^\varepsilon(x)\} + \sqrt{p'} d_\Gamma(x) \leq 0,$$

for every x such that $R \leq |x| \leq R'$, with $R' > R$. Our claim then follows by an inspection on the left-hand side, after applications for $\sigma = \sqrt{p'} R'$ and $\sigma = \sqrt{p'} R$ of Lemma A.1. \square

2.2. Asymptotics in a general domain. We begin by recalling that for equation (1.1) the *comparison principle* holds, as noted in [CGG, Remark 4.6] or in [AP, Appendix D] and as shown in [Sat, Theorem 2.1] in the general case of bounded solutions on unbounded domains, that is if u and v are viscosity solutions of (1.1) in Ω such that $u \leq v$ on Γ , then $u \leq v$ on $\overline{\Omega}$.

The next two lemmas give explicit barriers for the solution in a general domain Ω . We observe that no regularity assumption on Ω is needed.

Lemma 2.3 (Control from above). *Let $1 < p \leq \infty$ and u^ε be the bounded (viscosity) solution of (1.1)-(1.2).*

Then, we have that

$$\varepsilon \log \{u^\varepsilon(x)\} + \sqrt{p'} d_\Gamma(x) \leq \varepsilon \log E_p^\varepsilon(d_\Gamma(x)),$$

for every $x \in \overline{\Omega}$, where

$$(2.6) \quad E_p^\varepsilon(\sigma) = \begin{cases} \frac{\int_0^\pi (\sin \theta)^\alpha d\theta}{\int_0^\pi e^{-\sqrt{p'} \frac{1-\cos \theta}{\varepsilon} \sigma} (\sin \theta)^\alpha d\theta} & \text{if } 1 < p < \infty, \\ \frac{2}{1 + e^{-\frac{2\sigma}{\varepsilon}}} & \text{if } p = \infty. \end{cases}$$

In particular, it holds that

$$\varepsilon \log E_p^\varepsilon(d_\Gamma) = \begin{cases} O(\varepsilon \log \varepsilon) & \text{if } 1 < p < \infty, \\ O(\varepsilon) & \text{if } p = \infty, \end{cases}$$

as $\varepsilon \rightarrow 0^+$, on every subset of Ω in which d_Γ is bounded.

Proof. For a fixed $x \in \Omega$, we consider the ball $B^x = B_R(x)$ with $R = d_\Gamma(x)$ and denote by $u_{B^x}^\varepsilon$ the solution of (1.1)-(1.2) with $\Omega = B^x$. The comparison principle gives that

$$u^\varepsilon \leq u_{B^x}^\varepsilon \quad \text{on } \overline{B^x}$$

and, in particular,

$$(2.7) \quad u^\varepsilon(x) \leq u_{B^x}^\varepsilon(x).$$

Observe that the uniqueness of the solution of (1.1)-(1.2) and the scaling properties of Δ_p^G imply that

$$u_{B^x}^\varepsilon(x) = u_B^{\varepsilon/R}(0),$$

where u_B^ε is the solution of (1.1)-(1.2) with $\Omega = B$, the unit ball. The explicit expression in (2.1) and (2.7) then yield the desired claim, since $R = d_\Gamma(x)$.

The last claim follows from Lemma 2.1. \square

Lemma 2.4 (Control from below). *Let $1 < p \leq \infty$ and u^ε be the bounded (viscosity) solution of (1.1)-(1.2). Pick $z \in \mathbb{R}^N \setminus \overline{\Omega}$.*

Then, we have that

$$\varepsilon \log \{u^\varepsilon(x)\} + \sqrt{p'} \{|x - z| - d_\Gamma(z)\} \geq \varepsilon \log e_{p,z}^\varepsilon(x) \quad \text{for any } x \in \overline{\Omega},$$

where

$$(2.8) \quad e_{p,z}^\varepsilon(x) = \begin{cases} \frac{\int_0^\infty e^{-\sqrt{p'} \frac{\cosh \theta - 1}{\varepsilon} |x-z|} (\sinh \theta)^\alpha d\theta}{\int_0^\infty e^{-\sqrt{p'} \frac{\cosh \theta - 1}{\varepsilon} d_\Gamma(z)} (\sinh \theta)^\alpha d\theta} & \text{if } 1 < p < \infty, \\ 1 & \text{if } p = \infty. \end{cases}$$

Proof. We consider the ball $B = B_R(z)$ with radius $R = d_\Gamma(z)$ and let v^ε be the bounded solution of (1.1)-(1.2) relative to $\mathbb{R}^N \setminus \overline{B} \supset \Omega$. From the fact that $z \in \mathbb{R}^N \setminus \overline{\Omega}$, we have that $\Gamma \subset \mathbb{R}^N \setminus B$, which implies that

$$v^\varepsilon \leq 1 \quad \text{on } \Gamma,$$

by the explicit expression of v^ε given in (2.4). Thus, by the comparison principle, we infer that $v^\varepsilon \leq u^\varepsilon$ on $\overline{\Omega}$. The desired claim then follows by easy manipulations on (2.4). \square

Theorem 2.5 (Pointwise convergence). *Let $1 < p \leq \infty$ and Ω be a domain satisfying $\Gamma = \partial(\mathbb{R}^N \setminus \overline{\Omega})$; assume that u^ε is the bounded (viscosity) solution of (1.1)-(1.2).*

Then, it holds that

$$(2.9) \quad \lim_{\varepsilon \rightarrow 0^+} \varepsilon \log \{u_p^\varepsilon(x)\} = -\sqrt{p'} d_\Gamma(x) \quad \text{for any } x \in \overline{\Omega}.$$

Proof. Given $z \in \mathbb{R}^N \setminus \overline{\Omega}$ and $\varepsilon > 0$, combining Lemmas 2.3 and 2.4 gives at $x \in \overline{\Omega}$ that

$$(2.10) \quad \sqrt{p'} \{-|x-z| + d_\Gamma(x) + d_\Gamma(z)\} + \varepsilon \log e_{p,z}^\varepsilon(x) \leq \\ \varepsilon \log \{u^\varepsilon(x)\} + \sqrt{p'} d_\Gamma(x) \leq \varepsilon \log E_p^\varepsilon(d_\Gamma(x)).$$

Letting $\varepsilon \rightarrow 0^+$ then gives that

$$\sqrt{p'} \{-|x-z| + d_\Gamma(x) + d_\Gamma(z)\} \leq \\ \liminf_{\varepsilon \rightarrow 0^+} \left[\varepsilon \log \{u^\varepsilon(x)\} + \sqrt{p'} d_\Gamma(x) \right] \leq \\ \limsup_{\varepsilon \rightarrow 0^+} \left[\varepsilon \log \{u^\varepsilon(x)\} + \sqrt{p'} d_\Gamma(x) \right] \leq 0,$$

where we have used Lemma 2.3 and the fact that $\varepsilon \log e_{p,z}^\varepsilon(x)$ vanishes, as $\varepsilon \rightarrow 0^+$, by applying Lemma A.1 to (2.8).

Now, since z is arbitrary in $\mathbb{R}^N \setminus \overline{\Omega}$, we obtain (2.9), by taking the limit for $z \rightarrow y$, where $y \in \Gamma$ is a point realizing $|x-y| = d_\Gamma(x)$. \square

2.3. Uniform asymptotics. For a domain of class C^0 , we mean that its boundary is locally the graph of a continuous function. For the sequel, it is convenient to specify the modulus of continuity, by the following definition (introduced in [BM]). Let $\omega : (0, \infty) \rightarrow (0, \infty)$ be a strictly increasing continuous function such that $\omega(\tau) \rightarrow 0$ as $\tau \rightarrow 0^+$. We say that a domain Ω is of class $C^{0,\omega}$, if there exists a number $r > 0$ such that, for every point $x_0 \in \Gamma$, there is a coordinate system $(y', y_N) \in \mathbb{R}^{N-1} \times \mathbb{R}$, and a function $\phi : \mathbb{R}^{N-1} \rightarrow \mathbb{R}$ such that

- (i) $B_r(x_0) \cap \Omega = \{(y', y_N) \in B_r(x_0) : y_N < \phi(y')\}$;
- (ii) $B_r(x_0) \cap \Gamma = \{(y', y_N) \in B_r(x_0) : y_N = \phi(y')\}$

(iii) $|\phi(y') - \phi(z')| \leq \omega(|y' - z'|)$ for all $(y', \phi(y')), (z', \phi(z')) \in B_r(x_0) \cap \Gamma$.
In the sequel, it will be useful the function defined for $\varepsilon > 0$ by

$$\psi(\varepsilon) = \min_{0 \leq s \leq r} \sqrt{s^2 + [\omega(s) - \varepsilon]^2}$$

— this is the distance of the point $z_\varepsilon = (0', \varepsilon) \in \mathbb{R}^{N-1} \times \mathbb{R}$ from the graph of the function ω . Notice that $\psi(\varepsilon) = \varepsilon$ if $\phi \in C^k$ with $k \geq 2$ and, otherwise, $\psi(\varepsilon) \sim C\omega^{-1}(\varepsilon)$, for some positive constant C , where ω^{-1} is the inverse function of ω .

Theorem 2.6 (Uniform convergence). *Let $1 < p \leq \infty$ and Ω be a domain of class C^0 . Suppose that u^ε is the bounded (viscosity) solution of (1.1)-(1.2).*

Then, as $\varepsilon \rightarrow 0^+$, we have that

$$(2.11) \quad \varepsilon \log \{u^\varepsilon(x)\} + d_\Gamma(x) = \begin{cases} O(\varepsilon) & \text{if } p = \infty, \\ O(\varepsilon \log \varepsilon) & \text{if } p > N. \end{cases}$$

Moreover, if Ω is a $C^{0,\omega}$ domain, it holds that

$$(2.12) \quad \varepsilon \log \{u^\varepsilon(x)\} + \sqrt{p'} d_\Gamma(x) = \begin{cases} O(\varepsilon \log |\log \psi(\varepsilon)|) & \text{if } N = p, \\ O(\varepsilon \log \psi(\varepsilon)) & \text{if } 1 < p < N. \end{cases}$$

The formulas (2.11) and (2.12) hold uniformly on the compact subsets of $\bar{\Omega}$.

In particular, if $\varepsilon \log \psi(\varepsilon) \rightarrow 0$ as $\varepsilon \rightarrow 0^+$, then the convergence in (2.9) is uniform on every compact subset of $\bar{\Omega}$.

Proof. For any fixed compact subset K of $\bar{\Omega}$ we let d be the positive number, defined as

$$d = \max_{x' \in K} \{d_\Gamma(x'), |x'|\}.$$

To obtain the uniform convergence in (2.9) we will choose $z = z_\varepsilon$ independently on $x \in K$, as follows.

If Ω is a $C^{0,\omega}$ domain, fix $x \in K$, take $y \in \Gamma$ minimizing the distance to x , and consider a coordinate system in $\mathbb{R}^{N-1} \times \mathbb{R}$ such that $y = (0', 0)$. If we take $z_\varepsilon = (0', \varepsilon)$, then $z_\varepsilon \in \mathbb{R}^N \setminus \bar{\Omega}$ when ε is sufficiently small. With this choice, (2.10) reads as

$$\begin{aligned} \sqrt{p'} \{-|x - z_\varepsilon| + d_\Gamma(x) + d_\Gamma(z_\varepsilon)\} + \varepsilon \log e_{p,z_\varepsilon}^\varepsilon(x) &\leq \\ \varepsilon \log u_p^\varepsilon(x) + \sqrt{p'} d_\Gamma(x) &\leq \varepsilon \log E_p^\varepsilon(d_\Gamma(x)). \end{aligned}$$

Hence, we get:

$$(2.13) \quad -\sqrt{p'} \varepsilon + \varepsilon \log e_{p,z_\varepsilon}^\varepsilon(x) \leq \varepsilon \log u_p^\varepsilon(x) + \sqrt{p'} d_\Gamma(x) \leq \varepsilon \log E_p^\varepsilon(d_\Gamma(x)),$$

since $d_\Gamma(z_\varepsilon) \geq 0$ and $|x - z_\varepsilon| \leq d_\Gamma(x) + \varepsilon$.

Thus, if $p = \infty$, Lemmas 2.3 and (2.4) give that

$$-\varepsilon \leq \varepsilon \log \{u^\varepsilon(x)\} + d_\Gamma(x) \leq \varepsilon \log \left\{ \frac{2}{1 + e^{-\frac{d}{\varepsilon}}} \right\},$$

being $d_\Gamma(x) \leq d$, and (2.11) follows at once.

Next, if $1 < p < \infty$, we recall that $\varepsilon \log E_p^\varepsilon(d_\Gamma) = O(\varepsilon \log \varepsilon)$ on K as $\varepsilon \rightarrow 0^+$, by Lemma 2.3. On the other hand, by observing that $d_\Gamma(z_\varepsilon) \geq \psi(\varepsilon)$,

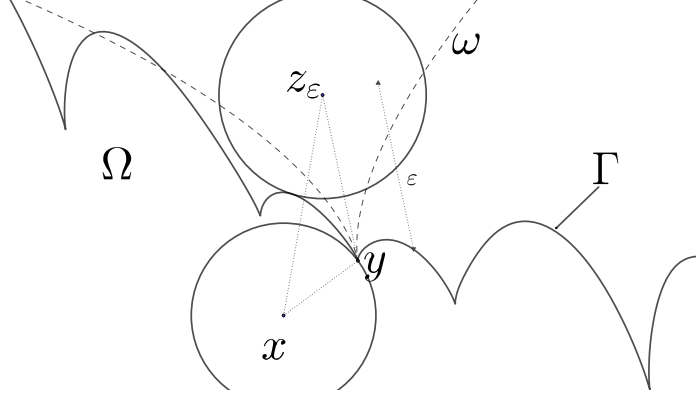


FIGURE 1. The geometric description of the approximating scheme in the proof of Theorem 2.6.

by our assumption on Ω , and that also $|x - z_\varepsilon| \leq 2d$ for $\varepsilon \leq d$, (2.8) gives on K that

$$e_{p,z_\varepsilon}^\varepsilon \geq \frac{\int_0^\infty e^{-\frac{2d\sqrt{p'}}{\varepsilon}(\cosh \theta - 1)} (\sinh \theta)^\alpha d\theta}{\int_0^\infty e^{-\frac{\sqrt{p'}\psi(\varepsilon)}{\varepsilon}(\cosh \theta - 1)} (\sinh \theta)^\alpha d\theta}.$$

Now, to this formula we apply Lemma A.1 with $\sigma = 2d\sqrt{p'}/\varepsilon$ at the numerator and $\sigma = \sqrt{p'}\psi(\varepsilon)/\varepsilon$ at the denominator. Thus, since by (2.2) the sign of α is that of $N - p$, on K we have as $\varepsilon \rightarrow 0$ that

$$\begin{aligned} \varepsilon \log(e_{p,z_\varepsilon}^\varepsilon) &\geq \alpha \varepsilon \log \psi(\varepsilon) - \frac{\alpha - 1}{2} \varepsilon \log \varepsilon + O(\varepsilon) = \\ &\qquad \qquad \qquad \alpha \varepsilon \log \psi(\varepsilon) + O(\varepsilon \log \varepsilon), \end{aligned}$$

if $p < N$,

$$\varepsilon \log(e_{p,z_\varepsilon}^\varepsilon) \geq -\varepsilon \log |\log \psi(\varepsilon)| + O(\varepsilon \log \varepsilon),$$

if $p = N$, and

$$\varepsilon \log(e_{p,z_\varepsilon}^\varepsilon) \geq \frac{\alpha + 1}{2} \varepsilon \log \varepsilon + O(\varepsilon),$$

if $p > N$. □

3. ASYMPTOTICS FOR THE q -MEANS OF A BALL

Throughout this subsection we assume that Ω is a (not necessarily bounded) domain that satisfies both the uniform exterior and interior ball conditions, i.e. there exist $r_i, r_e > 0$ such that every $y \in \Gamma$ has the property that there exist $z_i \in \Omega$ and $z_e \in \mathbb{R}^N \setminus \bar{\Omega}$ for which

$$(3.1) \quad B_{r_i}(z_i) \subset \Omega \subset \mathbb{R}^N \setminus \bar{B}_{r_e}(z_e) \quad \text{and} \quad \bar{B}_{r_i}(z_i) \cap \bar{B}_{r_e}(z_e) = \{y\}.$$

3.1. Enhanced barriers. We begin by refining the barriers given in Lemma 2.3 and 2.4, at least in a strip near the boundary Γ . To this aim, we will use the notation:

$$\Omega_\rho = \{y \in \Omega : d_\Gamma(y) \leq \rho\}, \quad \rho > 0.$$

We will also use the two families of probability measures on the intervals $[0, \infty)$ and $[0, \pi]$ with densities defined, respectively, by

$$d\nu^\sigma(\theta) = \frac{e^{-\sigma(\cosh\theta-1)}(\sinh\theta)^\alpha}{\int_0^\infty e^{-\sigma(\cosh\theta-1)}(\sinh\theta)^\alpha d\theta} d\theta,$$

$$d\mu^\sigma(\theta) = \frac{e^{-\sigma(1-\cos\theta)}(\sin\theta)^\alpha}{\int_0^\pi e^{-\sigma(1-\cos\theta)}(\sin\theta)^\alpha d\theta} d\theta.$$

Lemma 3.1. *Let u^ε be the bounded (viscosity) solution of (1.1)-(1.2).*

If $p \in (1, \infty)$, we set for $\sigma_\varepsilon = \sqrt{p'}r_e/\varepsilon$:

$$U^\varepsilon(\tau) = \int_0^\infty e^{-\tau \cosh\theta} d\nu^{\sigma_\varepsilon}(\theta), \quad \tau \geq 0,$$

and

$$V^\varepsilon(\tau) = \begin{cases} \int_0^\pi e^{-\tau \cos\theta} d\mu^{\sigma_\varepsilon}(\theta) & \text{if } 0 \leq \tau < \sigma_\varepsilon, \\ \left\{ \int_0^\pi e^{-\tau \cos\theta} d\mu^0(\theta) \right\}^{-1} & \text{if } \tau \geq \sigma_\varepsilon. \end{cases}$$

If $p = \infty$, we set $U^\varepsilon(\tau) = e^{-\tau}$ and

$$V^\varepsilon(\tau) = \begin{cases} \frac{\cosh(\sigma_\varepsilon - \tau)}{\cosh \sigma_\varepsilon} & \text{if } 0 \leq \tau < \sigma_\varepsilon, \\ 1/\cosh \tau & \text{if } \tau \geq \sigma_\varepsilon. \end{cases}$$

Then, we have that

$$(3.2) \quad U^\varepsilon \left(\frac{d_\Gamma(x)}{\varepsilon/\sqrt{p'}} \right) \leq u^\varepsilon(x) \leq V^\varepsilon \left(\frac{d_\Gamma(x)}{\varepsilon/\sqrt{p'}} \right),$$

for any $x \in \bar{\Omega}$.

Proof. Let $p \in (1, \infty)$. For any $x \in \Omega$ we can consider $y \in \Gamma$ such that $|x - y| = d_\Gamma(x)$. From the assumptions on Ω there exists $z_e \in \mathbb{R}^N \setminus \bar{\Omega}$ such that (3.1) holds for y . As seen in the proof of Lemma 2.4, by using the comparison principle and the explicit expression (2.4), we obtain

$$u^\varepsilon(x) \geq \frac{\int_0^\infty e^{-\sqrt{p'}|x-z_e|/\varepsilon \cosh\theta} (\sinh\theta)^\alpha d\theta}{\int_0^\infty e^{-\sqrt{p'}r_e/\varepsilon \cosh\theta} (\sinh\theta)^\alpha d\theta}.$$

Thus, the fact that $|x - z_e| = d_\Gamma(x) + r_e$ gives the first inequality in (3.2), by recalling the definition of U^ε .

To obtain the second inequality in (3.2) we proceed differently whether $x \in \Omega_{r_i}$ or not. Indeed, if $x \in \Omega_{r_i}$, there exists $z_i \in \Omega$ such that (3.1) holds for some $y \in \Gamma$ and $x \in B_{r_i}(z_i)$; moreover, since $\partial B_{d_\Gamma(x)}(x) \cap \partial B_{r_i}(z_i) = \{y\}$, we observe that x lies in the segment joining y to z_i , and hence $|x - z_i| = r_i - d_\Gamma(x)$. Again, by using the comparison principle and the expression in (2.1), we get that

$$u^\varepsilon(x) \leq \frac{\int_0^\pi e^{\sqrt{p'} \cos\theta \frac{|x-z_i|}{\varepsilon}} (\sin\theta)^\alpha d\theta}{\int_0^\pi e^{\sqrt{p'} \cos\theta \frac{r_i}{\varepsilon}} (\sin\theta)^\alpha d\theta},$$

that, by using the definition of V^ε and the fact that $|x - z_i| = r_i - d_\Gamma(x)$, leads to the second inequality in (3.2).

If $x \in \Omega \setminus \overline{\Omega}_{r_i}$, we just note that the expression of V^ε was already obtained in Lemma 2.3.

The case $p = \infty$ can be treated with similar arguments. \square

3.2. Asymptotics for q -means. From now on, in order to use the function Π_Γ defined in (1.7), we assume that Ω is a domain of class C^2 (not necessarily bounded).

First, we recall from [MS2, Lemma 2.1] the following geometrical lemma.

Lemma 3.2. *Let $x \in \Omega$ and assume that, for $R > 0$, there exists $y_x \in \Gamma$ such that $\overline{B_R(x)} \cap (\mathbb{R}^N \setminus \Omega) = \{y_x\}$ and that $\kappa_j(y_x) < 1/R$ for $j = 1, \dots, N-1$.*

Then, it holds that

$$\lim_{s \rightarrow 0^+} \frac{\mathcal{H}_{N-1}(\Gamma_s \cap B_R(x))}{s^{\frac{N-1}{2}}} = \frac{\omega_{N-1} (2R)^{\frac{N-1}{2}}}{(N-1)} [\Pi_\Gamma(y_x)]^{-\frac{1}{2}},$$

where \mathcal{H}_{N-1} denotes $(N-1)$ -dimensional Hausdorff measure and ω_{N-1} is the surface area of a unit sphere in \mathbb{R}^{N-1} .

The next lemma gives the asymptotic formula for $\varepsilon \rightarrow 0^+$ for the q -mean on $B_R(x)$ of a quite general class of functions, which includes both the barriers U^ε and V^ε .

Lemma 3.3. *Set $1 < q < \infty$. Let $x \in \Omega$ and assume that, for $R > 0$, there exists $y_x \in \Gamma$ such that $\overline{B_R(x)} \cap (\mathbb{R}^N \setminus \Omega) = \{y_x\}$ and that $\kappa_j(y_x) < 1/R$ for $j = 1, \dots, N-1$.*

Let $\{\xi_n\}_{n \in \mathbb{N}}$ and $\{f_n\}_{n \in \mathbb{N}}$ be sequences such that

- (i) $\xi_n > 0$ and $\xi_n \rightarrow 0$ as $n \rightarrow \infty$;
- (ii) $f_n : [0, \infty) \rightarrow [0, \infty)$ are decreasing functions;
- (iii) f_n converges to a function f almost everywhere as $n \rightarrow \infty$;
- (iv) it holds that

$$(3.3) \quad \lim_{n \rightarrow \infty} \int_0^\infty f_n(\tau)^{q-1} \tau^{\frac{N-1}{2}} d\tau = \int_0^\infty f(\tau)^{q-1} \tau^{\frac{N-1}{2}} d\tau,$$

and the last integral converges.

For some $1 < q < \infty$, let $\mu_{q,n}(x)$ be the q -mean of $f_n(d_\Gamma/\xi_n)$ on $B_R(x)$.

Then we have:

$$(3.4) \quad \lim_{n \rightarrow \infty} \left(\frac{R}{\xi_n} \right)^{\frac{N+1}{2(q-1)}} \mu_{q,n}(x) = \left\{ \frac{2^{-\frac{N+1}{2}} N!}{\Gamma\left(\frac{N+1}{2}\right)^2} \int_0^\infty f(\tau)^{q-1} \tau^{\frac{N-1}{2}} d\tau \right\}^{\frac{1}{q-1}} [\Pi_\Gamma(y_x)]^{-\frac{1}{2(q-1)}}.$$

Proof. From [IMW], we know that $\mu_n = \mu_{q,n}(x)$ is the only root of the equation

$$(3.5) \quad \int_{B_R(x)} [f_n(d_\Gamma/\xi_n) - \mu_n]_+^{q-1} dy = \int_{B_R(x)} [\mu_n - f_n(d_\Gamma/\xi_n)]_+^{q-1} dy,$$

where we mean $[t]_+ = \max(0, t)$.

Thus, if we set

$$\Gamma_\sigma = \{y \in B_R : d_\Gamma(y) = \sigma\},$$

by the co-area formula we get that

$$\int_{B_R(x)} [f_n(d_\Gamma/\xi_n) - \mu_n]_+^{q-1} dy = \int_0^{2R} [f_n(\sigma/\xi_n) - \mu_n]_+^{q-1} \mathcal{H}_{N-1}(\Gamma_\sigma) d\sigma,$$

that, after the change of variable $\sigma = \xi_n \tau$ and easy manipulations, leads to the formula:

$$\begin{aligned} \int_{B_R(x)} [f_n(d_\Gamma/\xi_n) - \mu_n]_+^{q-1} dy = \\ \xi_n^{\frac{N+1}{2}} \int_0^{2R/\xi_n} [f_n(\tau) - \mu_n]_+^{q-1} \tau^{\frac{N-1}{2}} \left[\frac{\mathcal{H}_{N-1}(\Gamma_{\xi_n \tau})}{(\xi_n \tau)^{\frac{N-1}{2}}} \right] d\tau. \end{aligned}$$

Therefore, since $\mu_n \rightarrow 0$ as $n \rightarrow \infty$, an inspection of the integrand at the right-hand side, assumptions (i)-(iv), and Lemma 3.2 make it clear that we can apply the generalized dominated convergence theorem (see [LL]) to infer that

$$(3.6) \quad \lim_{n \rightarrow \infty} \xi_n^{-\frac{N+1}{2}} \int_{B_R(x)} [f_n(d_\Gamma/\xi_n) - \mu_n]_+^{q-1} dy = \frac{(2R)^{\frac{N-1}{2}} \omega_{N-1}}{(N-1)\sqrt{\Pi_\Gamma(y_x)}} \int_0^\infty f(\tau)^{q-1} \tau^{\frac{N-1}{2}} d\tau.$$

Next, by employing again the co-area formula, the right-hand side of (3.5) can be re-arranged as

$$\int_{B_R(x)} [\mu_n - f_n(d_\Gamma/\xi_n)]_+^{q-1} dy = \mu_n^{q-1} \int_0^{2R} \left[1 - \frac{f_n(\sigma/\xi_n)}{\mu_n} \right]_+^{q-1} \mathcal{H}_{N-1}(\Gamma_\sigma) d\sigma,$$

that leads to the formula

$$(3.7) \quad \lim_{\varepsilon \rightarrow 0^+} \mu_\varepsilon^{1-q} \int_{B_R(x)} [\mu_n - f_n(d_\Gamma/\xi_n)]_+^{q-1} dy = |B_R|,$$

by dominated convergence theorem, if we can prove that

$$(3.8) \quad \frac{f_n(\sigma/\xi_n)}{\mu_n} \rightarrow 0 \quad \text{as } n \rightarrow \infty,$$

for almost every $\sigma \geq 0$. Then, after straightforward computations, (3.4) will follow by putting together (3.5), (3.6) and (3.7).

We now complete the proof by proving that (3.8) holds. From (3.5), (3.6), and the fact that

$$\int_{B_R(x)} [\mu_n - f_n(d_\Gamma/\xi_n)]_+^{q-1} dy \leq \mu_n^{q-1} |B_R|,$$

we have that there is a positive constant c such that

$$\mu_n^{1-q} \leq c \xi_n^{-\frac{N+1}{2}}.$$

Also, for every $\sigma > 0$ we have that

$$\begin{aligned} \int_{\sigma/2\xi_n}^{\infty} f_n(\tau)^{q-1} \tau^{\frac{N-1}{2}} d\tau &\geq \int_{\sigma/2\xi_n}^{\sigma/\xi_n} f_n(\tau)^{q-1} \tau^{\frac{N-1}{2}} d\tau \geq \\ &\frac{2(1 - 2^{-\frac{N+1}{2}})}{N+1} f_n(\sigma/\xi_n)^{q-1} \left(\frac{\sigma}{\xi_n}\right)^{\frac{N+1}{2}} \geq \\ &\frac{2(1 - 2^{-\frac{N+1}{2}})}{c(N+1)} \sigma^{\frac{N+1}{2}} \left\{ \frac{f_n(\sigma/\xi_n)}{\mu_n} \right\}^{q-1}. \end{aligned}$$

Thus, (3.8) follows, since the first term of this chain of inequalities converges to zero as $n \rightarrow \infty$, under our assumptions on f_n and ξ_n , in virtue of the generalized dominated convergence theorem. \square

Remark 3.4. The case $q = \infty$ is simpler. From [IMW] and then the monotonicity of f_n we obtain that:

$$\begin{aligned} \mu_{\infty,n}(x) &= \frac{1}{2} \left\{ \min_{B_R(x)} f_n(d_\Gamma/\xi_n) + \max_{B_R(x)} f_n(d_\Gamma/\xi_n) \right\} = \\ &\frac{1}{2} \{f_n(2R/\xi_n) + f_n(0)\}. \end{aligned}$$

Thus, if we replace the assumptions (iii) and (iv) by $f_n(0) \rightarrow f(0)$ as $n \rightarrow \infty$, we conclude that $\mu_{\infty,n}(x) \rightarrow f(0)/2$, since $f_n(2R/\xi_n) \rightarrow 0$ as $n \rightarrow \infty$.

Theorem 3.5. *Set $1 < p \leq \infty$. Let $x \in \Omega$ be such that $B_R(x) \subset \Omega$ and $B_R(x) \cap (\mathbb{R}^N \setminus \Omega) = \{y_x\}$; suppose that $k_j(y_x) < \frac{1}{R}$, for every $j = 1, \dots, N-1$.*

Let u^ε be the bounded (viscosity) solution of (1.1)-(1.2) and, for $1 < q \leq \infty$, let $\mu_{q,\varepsilon}(x)$ be the q -mean of u^ε on $B_R(x)$.

Then, if $1 < q < \infty$, we have that (1.10) holds, that is

$$\lim_{\varepsilon \rightarrow 0^+} \left(\frac{\varepsilon}{R}\right)^{-\frac{N+1}{2(q-1)}} \mu_{q,\varepsilon}(x) = \frac{c_{N,q}}{\left\{ (p')^{\frac{N+1}{2}} \Pi_\Gamma(y_x) \right\}^{\frac{1}{2(q-1)}}},$$

where

$$c_{N,q} = \left\{ \frac{2^{-\frac{N+1}{2}} N!}{(q-1)^{\frac{N+1}{2}} \Gamma\left(\frac{N+1}{2}\right)} \right\}^{\frac{1}{q-1}}.$$

If $q = \infty$, we simply have that $\mu_{\infty,\varepsilon}(x) \rightarrow 1/2$ as $\varepsilon \rightarrow 0$.

Proof. We have that $\mu_{q,\varepsilon}^{U^\varepsilon}(x) \leq \mu_{q,\varepsilon}(x) \leq \mu_{q,\varepsilon}^{V^\varepsilon}(x)$ by the monotonicity properties of the q -means, where with $\mu_{q,\varepsilon}^{U^\varepsilon}$ and $\mu_{q,\varepsilon}^{V^\varepsilon}$ we denote the q -mean of $U^\varepsilon(d/\varepsilon)$ and $V^\varepsilon(d/\varepsilon)$ on $B_R(x)$. Hence, in order to prove (1.10), we only need to apply Lemma 3.3 to $f_n = U^{\varepsilon_n}$ and $f_n = V^{\varepsilon'_n}$, where the vanishing sequences ε_n and ε'_n are chosen so that the liminf and limsup of $(\varepsilon/R)^{-\frac{N+1}{2(q-1)}} \mu_{q,\varepsilon}(x)$ as $\varepsilon \rightarrow 0$ are attained along them, respectively.

By an inspection, it is not difficult to check that $f_n = U^{\varepsilon_n}$ and $f_n = V^{\varepsilon'_n}$, with $\xi_\varepsilon = \varepsilon/\sqrt{p'}$ and $f(\tau) = e^{-\tau}$, satisfy the relevant assumptions of Lemma 3.3, by applying, in particular, Lemma A.2 for (iii) and the dominated convergence theorem for (iv). \square

APPENDIX A. TECHNICAL LEMMAS

Here, we collect two useful lemmas.

Lemma A.1 (One-dimensional asymptotics). *For $\alpha > -1$ and $\sigma > 0$, let*

$$f(\sigma) = \int_0^\infty e^{-\sigma(\cosh \theta - 1)} (\sinh \theta)^\alpha d\theta.$$

Then, f is continuous in $(0, \infty)$ and

$$f(\sigma) = 2^{\frac{\alpha-1}{2}} \Gamma\left(\frac{\alpha+1}{2}\right) \sigma^{-\frac{\alpha+1}{2}} \{1 + O(1/\sigma)\} \text{ as } \sigma \rightarrow \infty.$$

Moreover, if $\sigma \rightarrow 0$, we have that

$$f(\sigma) = \begin{cases} \sigma^{-\alpha} \Gamma(\alpha) \{1 + o(1)\} & \text{if } \alpha > 0, \\ \log(1/\sigma) + O(1) & \text{if } \alpha = 0, \\ \frac{\sqrt{\pi}}{2 \sin(\alpha\pi/2)} \frac{\Gamma(\frac{\alpha+1}{2})}{\Gamma(\frac{\alpha}{2}+1)} + o(1) & \text{if } -1 < \alpha < 0. \end{cases}$$

Proof. By the change of variable $\tau = \sigma(\cosh \theta - 1)$ we get:

$$f(\sigma) = \frac{1}{\sigma} \int_0^\infty e^{-\tau} \left(\frac{2\tau}{\sigma} + \frac{\tau^2}{\sigma^2}\right)^{\frac{\alpha-1}{2}} d\tau.$$

When $\sigma \rightarrow \infty$, our claim follows by writing

$$f(\sigma) = 2^{\frac{\alpha-1}{2}} \sigma^{-\frac{\alpha+1}{2}} \int_0^\infty e^{-\tau} \left(\tau + \frac{\tau^2}{2\sigma}\right)^{\frac{\alpha-1}{2}} d\tau.$$

When $\sigma \rightarrow 0$ and $\alpha > 0$, our claim follows by writing

$$f(\sigma) = \sigma^{-\alpha} \int_0^\infty e^{-\tau} (\tau^2 + 2\sigma\tau)^{\frac{\alpha-1}{2}} d\tau,$$

since in this case the (limiting) integral converges near zero.

For $-1 < \alpha \leq 0$, we use the formula

$$f(\sigma) = \frac{1}{\sqrt{\pi}} \Gamma\left(\frac{\alpha+1}{2}\right) \left(\frac{\sigma}{2}\right)^{-\frac{\alpha}{2}} e^\sigma K_{\alpha/2}(\sigma),$$

where $K_{\alpha/2}(\sigma)$ is the modified Bessel's function of the second kind of order $\alpha/2$ (see [AS, Formula 9.6.23]). Then, [AS, Formula 9.6.9] and [AS, Formula 9.6.13] give our claims for $\alpha = 0$ and $-1 < \alpha < 0$, respectively. \square

Lemma A.2 (Mollifier). *Let $g : \mathbb{R} \rightarrow \mathbb{R}$ be a bounded and continuous function.*

Then, we have that

$$\lim_{\sigma \rightarrow \infty} \int_0^\infty g(\theta) \left[\frac{e^{-\sigma(\cosh \theta - 1)} (\sinh \theta)^\alpha}{\int_0^\infty e^{-\sigma(\cosh \theta - 1)} (\sinh \theta)^\alpha} \right] d\theta = g(0)$$

and

$$\lim_{\sigma \rightarrow \infty} \int_0^\pi g(\theta) \left[\frac{e^{-\sigma(1 - \cos \theta)} (\sin \theta)^\alpha}{\int_0^\pi e^{-\sigma(1 - \cos \theta)} (\sin \theta)^\alpha} d\theta \right] d\theta = g(0).$$

Proof. The two formulas in the statement follow by observing that in both cases the relevant integral can be written as

$$\int_0^\infty j_\sigma(\theta) g(\theta) d\theta,$$

where j_σ is a *mollifier*, that has the salient properties:

$$j_\theta \geq 0, \quad \int_0^\infty j_\sigma(\theta) d\theta = 1 \quad \text{for any } \sigma > 0 \quad \text{and}$$

$$\lim_{\sigma \rightarrow \infty} \int_\delta^\infty j_\sigma(\theta) d\theta = 0 \quad \text{for any } \delta > 0.$$

The last property easily follows from the substitution $\tau = \sigma (\cosh \theta - 1)$ or $\tau = \sigma (1 - \cos \theta)$. \square

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