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## Harnack type estimates and Hölder continuity for non-negative solutions to certain sub-critically singular parabolic partial differential equations

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**Abstract.** A two-parameter family of Harnack type inequalities for non-negative solutions of a class of singular, quasilinear, homogeneous parabolic equations is established, and it is shown that such estimates imply the Hölder continuity of solutions. These classes of singular equations include p-Laplacean type equations in the sub-critical range  $1 and equations of the porous medium type in the sub-critical range <math>0 < m \le \frac{(N-2)_+}{N}$ .

#### 1. Introduction and main results

Let E be an open set in  $\mathbb{R}^N$  and for T > 0 let  $E_T$  denote the cylindrical domain  $E \times (0, T]$ . Consider quasi-linear, parabolic differential equations of the form

$$u \in C_{\text{loc}}\left(0, T; L_{\text{loc}}^{2}(E)\right) \cap L_{\text{loc}}^{p}\left(0, T; W_{\text{loc}}^{1, p}(E)\right)$$

$$u_{t} - \text{div } \mathbf{A}(x, t, u, Du) = 0 \quad \text{weakly in } E_{T}$$

$$(1.1)$$

where the function  $\mathbf{A}: E_T \times \mathbb{R}^{N+1} \to \mathbb{R}^N$  is only assumed to be measurable and subject to the structure conditions

$$\begin{cases} \mathbf{A}(x, t, u, Du) \cdot Du \ge C_o |Du|^p \\ |\mathbf{A}(x, t, u, Du)| \le C_1 |Du|^{p-1} \end{cases} \text{ a.e. } (x, t) \in E_T$$
 (1.2)

where  $C_0$  and  $C_1$  are given positive constants, and p is in the sub-critical range

$$1$$

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The homogeneous prototype of such a class of parabolic equations is

$$u_t - \text{div} |Du|^{p-2} Du = 0$$
 weakly in  $E_T$ . ((1.1)<sub>o</sub>)

The parameters  $\{N, p, C_o, C_1\}$  are the data, and we say that a generic constant  $\gamma = \gamma(N, p, C_o, C_1)$  depends upon the data, if it can be quantitatively determined a priori only in terms of the indicated parameters. For  $\rho > 0$  let  $B_\rho$  be the ball of center the origin on  $\mathbb{R}^N$  and radius  $\rho$  and for  $y \in \mathbb{R}^N$  let  $B_\rho(y)$  denote the homothetic ball centered at y. For  $\tau > 0$  and for  $(y, s) \in \mathbb{R}^N \times \mathbb{R}$  set also

$$Q_{\rho}(\tau) = B_{\rho} \times (-\tau, 0], \quad (y, s) + Q_{\rho}(\tau) = B_{\rho}(y) \times (s - \tau, s].$$

Let *u* be a non-negative weak solution of (1.1–1.3). Having fixed  $(x_o, t_o) \in E_T$ , and  $B_{4\rho}(x_o) \subset E$ , introduce the quantities

$$\int_{B_{\rho}(x_o)} u^q(x, t_o) dx, \quad \delta \stackrel{\text{def}}{=} \left[ \varepsilon \left( \int_{B_{\rho}(x_o)} u^q(\cdot, t_o) dx \right)^{\frac{1}{q}} \right]^{2-p} \rho^p \tag{1.4}$$

where  $\varepsilon \in (0, 1)$  is to be chosen, and  $q \ge 1$  is arbitrary. If  $\delta > 0$ , set also

$$\eta \stackrel{\text{def}}{=} \left[ \frac{\left( f_{B_{\rho}(x_o)} u^q(\cdot, t_o) dx \right)^{\frac{1}{q}}}{\left( f_{B_{4\rho}(x_o)} u^r(\cdot, t_o - \delta) dx \right)^{\frac{1}{r}}} \right]^{\frac{r\rho}{\lambda_r}}$$
(1.5)

where  $r \geq 1$  is any number such that

$$\lambda_r \stackrel{\text{def}}{=} N(p-2) + rp > 0. \tag{1.6}$$

**Theorem 1.1.** Let u be a non-negative, locally bounded, local, weak solution of (1.1–1.3). Introduce  $\delta$  as in (1.4) and assume that  $\delta > 0$ . There exist constants  $\varepsilon \in (0,1)$ , and  $\gamma > 1$ , depending only on the data and the parameters q, r, and a constant  $\beta > 1$ , depending only upon the data and independent of q, r, such that

$$\inf_{(x_{o},t_{o})+Q_{\rho}(\frac{1}{2}\delta)} u \geq \gamma \left[ \frac{\left( \int_{B_{\rho}(x_{o})} u^{q}(\cdot,t_{o}) dx \right)^{\frac{1}{q}}}{\left( \int_{B_{4\rho}(x_{o})} u^{r}(\cdot,t_{o}-\delta) dx \right)^{\frac{1}{r}}} \right]^{\beta \frac{rp}{\lambda_{r}}} \sup_{(x_{o},t_{o})+Q_{\rho}(\delta)} u \quad (1.7)$$

provided  $q \ge 1$  and  $r \ge 1$  satisfies (1.6) and  $(x_o, t_o) + Q_{8\rho}(\delta) \subset E_T$ . The constants  $\varepsilon \to 0$ , and  $\gamma \to \infty$  as either  $\lambda_r \to 0$  or  $\lambda_r \to \infty$ .

Remark 1.1. The estimate is vacuous if  $\delta = 0$ . This does occur for certain solutions of (1.1) for  $t_0$  larger than the extinction time ([8]).

Remark 1.2. Inequality (1.7) is not a Harnack inequality per se, since  $\eta$  depends upon the solution itself. It would reduce to a Harnack inequality if  $\eta \geq \eta_o > 0$  for some absolute constant  $\eta_o$  depending only upon the data. This however cannot occur since a Harnack inequality for solutions of (1.1–1.3) does not hold, as shown by the counterexamples of [7]. Further comments in this direction are in Remark 4.1.

Remark 1.3. An estimate similar to (1.7) has been derived in [3] for non-negative solutions of the prototype equation  $(1.1)_o$ , by means of maximum and comparison principles, and some asymptotic estimates of [8]. However the Harnack inequality is a structural property of a parabolic equation, unrelated to comparison and maximum principles. This emerges from the pioneering work of Moser [9, 10], and the results of [1,4,6,11]. Theorem 1.1 is in this direction.

Remark 1.4. Inequality (1.7) actually holds for non-negative solutions of (1.1–1.2) for all  $1 , provided <math>r \ge 1$  satisfies (1.6). For super-critical values of  $p > \frac{2N}{N+1}$  one has  $\lambda = \lambda_1 > 0$ , and (1.6) can be realized for r = 1. However, for  $\lambda > 0$  the strong form of a Harnack estimate holds ([7]). Therefore (1.7), while true for all 1 , holds significance only for critical and sub-critical values <math>1 . In this sense (1.7) can be regarded as a "weak" form of a Harnack estimate. Neverthless (1.7) is sufficient to establish the local Hölder continuity of locally bounded, weak solutions of (1.1–1.2), irrespective of their sign, as we show in Sect. 4.

#### 2. Components of the proof of Theorem 1.1

2.1.  $L_{loc}^r$ - $L_{loc}^{\infty}$  Estimates For  $r \ge 1$  Such That  $\lambda_r > 0$ 

**Proposition 2.1.** Let u be a non-negative, locally bounded, local, weak solution to (1.1–1.3), and assume that  $u \in L^r_{loc}(E_T)$  for some  $r \ge 1$ , satisfying (1.6). There exists a positive constant  $\gamma_r$  depending only upon the data, and r, such that

$$\sup_{B_{\rho}(y)\times[s,t]}u\leq \frac{\gamma_r}{(t-s)^{\frac{N}{\lambda_r}}}\left(\int_{B_{2\rho}(y)}u^r(x,2s-t)dx\right)^{\frac{p}{\lambda_r}}+\gamma_r\left(\frac{t-s}{\rho^p}\right)^{\frac{1}{2-p}} \tag{2.1}$$

for all cylinders

$$B_{2\rho}(y) \times [s - (t - s), s + (t - s)] \subset E_T.$$
 (2.2)

The constant  $\gamma_r \to \infty$  if either  $\lambda_r \to 0$  or  $\lambda_r \to \infty$ .

*Remark 2.1.* The values of u in the upper part of the cylinder (2.2) are estimated by the integral on the lower base of the cylinder.

Remark 2.2. The local boundedness of a weak solution is insured by the integrability  $u \in L^r_{loc}(E_T)$  for some  $r \ge 1$  satisfying (1.6). For  $\frac{2N}{N+2} , such an integrability condition is a consequence of the notion of weak solution. Indeed$ 

by parabolic embedding (see [5], Chapter I, Proposition 3.1),  $u \in L^m_{loc}(E_T)$  with  $m = \frac{N+2}{N}p$ , and  $\lambda_m > 0$ . For 1 this is no longer the case, and the integrability requirement is an extra assumption imposed on the notion of weak solution, to insure its local boundedness. Indeed for <math>1 there exist unbounded, local, weak solutions to (1.1) ([5]).

The proof of Proposition 2.1 follows arguments similar to those in of [5] Chap. V, with minor modifications outlined in Appendix A.

#### 2.2. Expansion of positivity

**Proposition 2.2.** Let u be a non-negative, local, weak solution to (1.1-1.2), for 1 , satisfying

$$\left| [u(\cdot, t) > M] \cap B_{\rho}(y) \right| > \alpha |B_{\rho}| \tag{2.3}$$

for all times

$$s - \epsilon M^{2-p} \rho^p \le t \le s \tag{2.4}$$

for some M > 0, and some  $\alpha, \epsilon \in (0, 1)$ . Assume moreover that

$$B_{8\rho} \times (s - \epsilon M^{2-p} \rho^p, s) \subset E_T$$
.

There exists  $\sigma \in (0, 1)$  that can be determined a priori, quantitatively only in terms of the data, and the numbers  $\alpha$  and  $\epsilon$ , independent of M, such that

$$u(x,t) \ge \sigma M \quad for \ all \ \ x \in B_{2\rho}(y)$$
 (2.5)

for all times

$$s - \frac{1}{2}\epsilon M^{2-p}\rho^p < t \le s. \tag{2.6}$$

Remark 2.3. Thus measure-theoretical information on the measure of the "positivity set" in  $B_{\rho}(y)$  for all times in (2.4) implies that such a positivity set actually expands to  $B_{2\rho}(y)$  for comparable times. This is the main underlying structural fact of a Harnack inequality.

*Remark 2.4.* The proof, given in [5], Chap. IV, and in [7], shows that the functional dependence of  $\sigma$  on  $\epsilon$  and  $\alpha$  is of the form

$$\sigma(\epsilon, \alpha) \approx a^{1/\epsilon^b \alpha^c}$$
 (2.7)

for constants  $a \in (0, 1)$  and b, c > 1 depending only upon the data.

Remark 2.5. Proposition 2.2 holds for all 1 , irrespective of <math>p belonging to the sub-critical or super-critical range.

#### 2.3. $L_{loc}^r$ estimates backward in time

**Proposition 2.3.** Let u be a non-negative, local, weak solution to (1.1-1.2), for  $1 , and assume that <math>u \in L^r_{loc}(E_T)$  for some  $r \ge 1$ . There exists a constant  $\bar{\gamma}_r$  depending only upon the data and r, such that for all cylinders  $B_{2\rho}(y) \times [\tau, t] \subset E_T$ 

$$\sup_{\tau \le s \le t} \int_{B_{\rho}(y)} u^r(x, s) dx \le \bar{\gamma}_r \int_{B_{2\rho}(y)} u^r(x, \tau) dx + \bar{\gamma}_r \left[ \frac{(t - \tau)^r}{\rho^{\lambda_r}} \right]^{\frac{1}{2-p}} \tag{2.8}$$

where  $\lambda_r$  is defined in (1.6), but it is not required to be positive.

The proof is in Appendix A. If r=1 this estimate can be given the form of a Harnack inequality in the  $L^1_{loc}$  topology.

**Proposition 2.4.** Let u be a non-negative, local, weak solution to (1.1–1.2), for  $1 . There exists a positive constant <math>\bar{\gamma}$  depending only upon the data, such that for all cylinders  $B_{2\rho}(y) \times [\tau, t] \subset E_T$ 

$$\sup_{\tau \le s \le t} \int_{B_{\rho}(y)} u(x,\tau) dx \le \bar{\gamma} \inf_{\tau \le s \le t} \int_{B_{2\rho}(y)} u(x,\tau) dx + \bar{\gamma} \left( \frac{t-\tau}{\rho^{\lambda}} \right)^{\frac{1}{2-p}}$$
 (2.9)

where  $\lambda = \lambda_1$  is defined in (1.6), but it is not required to be positive.

If  $p_* then <math>\lambda > 0$ , whereas  $1 implies <math>\lambda \le 0$ . However (2.9) holds true for all  $1 and accordingly, <math>\lambda$  could be of either sign. The constant  $\bar{\gamma} = \bar{\gamma}(p) \to \infty$  as either  $p \to 2$  or  $p \to 1$ . The proof is in [7].

#### 3. Estimating the positivity set of the solutions

Having fixed  $(x_o, t_o) \in E_T$ , assume it coincides with the origin, write  $B_\rho(0) = B_\rho$  and introduce the quantity  $\delta$  as in (1.4), which is assumed to be positive. From (2.8) and the definition of  $\delta$ 

$$\int_{B_0} u^q(\cdot,0) dx \le \bar{\gamma}_q \int_{B_{20}} u^q(\cdot,\tau) dx + \bar{\gamma}_q \, \varepsilon^q \int_{B_0} u^q(\cdot,0) dx$$

for all  $q \ge 1$  and for all  $\tau \in (-\delta, 0]$ . Choosing  $\bar{\gamma}_q \varepsilon^q \le \frac{1}{2}$  yields

$$\int_{B_{2\rho}} u^q(\cdot, \tau) dx \ge \frac{1}{2\bar{\gamma}_q} \int_{B_\rho} u^q(\cdot, 0) dx \quad \text{for all } \tau \in (-\delta, 0]. \tag{3.1}$$

Next apply the sup-estimate (2.1) over the cylinder  $B_{2\rho} \times (-\frac{1}{2}\delta, 0]$  with  $r \ge 1$  such that  $\lambda_r > 0$ , to get

$$\begin{split} \sup_{B_{2\rho}\times(-\frac{1}{2}\delta,0]} u &\leq \frac{\gamma_r[\omega_N(4\rho)^N]^{\frac{\rho}{\lambda_r}}}{\delta^{\frac{N}{\lambda_r}}} \left( \int_{B_{4\rho}} u^r(\cdot,-\delta) dx \right)^{\frac{1}{r}\frac{r\rho}{\lambda_r}} + \gamma_r \left( \frac{\delta}{\rho^p} \right)^{\frac{1}{2-p}} \\ &\leq \frac{\gamma_r'}{\varepsilon^{\frac{N(2-p)}{\lambda_r}}} \frac{1}{\eta} \left( \int_{B_\rho} u^q(\cdot,0) dx \right)^{\frac{1}{q}} + \gamma_r' \varepsilon \left( \int_{B_\rho} u^q(\cdot,0) dx \right)^{\frac{1}{q}} \\ &= \gamma_r' \varepsilon \left( 1 + \frac{1}{\eta \varepsilon^{\frac{r\rho}{\lambda_r}}} \right) \left( \int_{B_\rho} u^q(\cdot,0) dx \right)^{\frac{1}{q}} \end{split}$$

for a constant  $\gamma'_r$  depending only upon the data and r. One verifies that  $\gamma'_r \to \infty$  as either  $\lambda_r \to 0$  or  $\lambda_r \to \infty$ .

Assume momentarily that  $0 < \eta < 1$  so that in the round brackets containing  $\eta$ , the second term dominates the first. In such a case

$$\sup_{B_{2\rho}\times(-\frac{1}{2}\delta,0]} u \le M \stackrel{\text{def}}{=} \frac{1}{\varepsilon'\eta} \left( \int_{B_{\rho}} u^q(\cdot,0) dx \right)^{\frac{1}{q}} \quad \text{where } \varepsilon' = \frac{\varepsilon^{\frac{N(2-\rho)}{\lambda_r}}}{2\gamma_r'}. \quad (3.2)$$

and therefore

$$\varepsilon'\eta M = \left(\int_{B_{\rho}} u^{q}(\cdot, 0) dx\right)^{\frac{1}{q}}.$$
 (3.3)

Let  $\nu \in (0, 1)$  to be chosen. Combining (3.3) with (3.1) gives

$$\begin{split} &(\varepsilon'\eta M)^q \leq 2^{N+1}\bar{\gamma}_q \int_{B_{2\rho}} u^q(\cdot,\tau) dx \\ &\leq 2^{N+1}\bar{\gamma}_q v^q (\eta M)^q + 2^{N+1}\bar{\gamma}_q M^q \frac{|[u(\cdot,\tau) > v\eta M] \cap B_{2\rho}|}{|B_{2\rho}|} \end{split}$$

for all  $\tau \in (-\frac{1}{2}\delta, 0]$ . From this

$$|[u(\cdot, \tau) > \nu \eta M] \cap B_{2\rho}| \ge \alpha \eta^q |B_{2\rho}| \quad \text{where } \alpha = \frac{\varepsilon'^q - \nu^q \, 2^{N+1} \bar{\gamma}_q}{2^{N+1} \bar{\gamma}_q} \quad (3.4)$$

for all  $\tau \in (-\frac{1}{2}\delta, 0]$ . By choosing  $\nu \in (0, 1)$  sufficiently small, only dependent on the data and  $\bar{\gamma}_q$ , we can insure that  $\alpha \in (0, 1)$  depends only upon the data and q, and is independent of  $\eta$ .

**Proposition 3.1.** Let u be a non-negative, locally bounded, local, weak solution of (1.1-1.2) for  $1 . Fix <math>(x_o, t_o) \in E_T$ , let  $B_{4\rho}(x_o) \subset E$  and let  $\delta$  and  $\eta$  be defined by (1.4-1.6) for some  $\varepsilon \in (0, 1)$ . For every  $r \geq 1$  satisfying (1.6) and every  $q \geq 1$ , there exist constants  $\varepsilon$ , v,  $\alpha \in (0, 1)$ , depending only upon the data and q and r, such that

$$|[u(\cdot,t) > \nu \eta M] \cap B_{2\rho}(x_o)| \ge \alpha \eta^q |B_{2\rho}| \quad \text{for all } t \in (t_o - \frac{1}{2}\delta, t_o].$$
 (3.5)

#### 3.1. A first form of the Harnack inequality

The definition of (1.4) of  $\delta$  and the parameters in (3.2–3.4), imply that

$$\frac{1}{2}\delta = \epsilon (\nu \eta M)^{2-p} \rho^p$$
 where  $\epsilon = \frac{1}{2} \left(\frac{\varepsilon \varepsilon'}{\nu}\right)^{2-p}$ . (3.6)

Therefore by Proposition 2.2 with M replaced by  $\nu \eta M$  and  $\alpha$  replaced by  $\alpha \eta^q$ 

$$u(\cdot,t) > \sigma(\alpha \eta^q, \epsilon) \nu \eta M$$
 in  $B_{4\rho}(x_o)$ , for all  $t \in (t_o - \frac{1}{4}\delta, t_o]$ .

**Proposition 3.2.** Let u be a non-negative, locally bounded, local, weak solution of (1.1-1.3). Fix  $(x_o, t_o) \in E_T$ , let  $B_{4\rho}(x_o) \subset E$  and let  $\delta$  and  $\eta$  be defined by (1.4-1.6) for some  $\varepsilon \in (0, 1)$ . For every  $r \geq 1$  satisfying (1.6) and every  $q \geq 1$ , there exist a constant  $\varepsilon$ , depending only upon the data and q and r, and a continuous, increasing function  $\eta \to f(\eta)$  defined in  $\mathbb{R}^+$  and vanishing at  $\eta = 0$ , that can be quantitatively determined a priori only in terms of the data, such that

$$\inf_{B_{4\rho}(x_o)} u(\cdot, t) \ge f(\eta) \sup_{(x_o, t_o) + Q_{2\rho}(\frac{1}{4}\delta)} u, \quad \text{for all } t \in (t_o - \frac{1}{4}\delta, t_o].$$
 (3.7)

provided  $(x_o, t_o) + Q_{8\rho}(\delta) \subset E_T$ .

Remark 3.1. In view of (2.7) the function  $f(\cdot)$  can be taken of the form

$$f(\eta) \approx \eta \, B^{-\frac{1}{\eta^d}}$$

for constants B, d > 1 depending only upon the data and q and r.

*Remark 3.2.* The function  $f(\cdot)$  depends on  $\delta$  only through the parameter  $\varepsilon$  in the definition (1.4) of  $\delta$ .

Remark 3.3. The inequality (3.7) is a Harnack type estimate of the same form as that established in [7], where however the constant  $f(\eta)$  depends on the solution itself, through  $\eta$  defined in (1.5), as a proper quotient of the  $L_{\text{loc}}^q$  and  $L_{\text{loc}}^r$  averages of u, respectively at time  $t = t_0$  on ball  $B_{\rho}(x_0)$ , and at time  $t = t_0 - \delta$  on ball  $B_{4\rho}(x_0)$ .

Remark 3.4. The inequality (3.7) has been derived by assuming that  $0 < \eta < 1$ . If  $\eta \ge 1$  the same proof gives (3.7) where  $f(\eta) \ge f(1)$ , thereby establishing a strong form of the Harnack estimate for these solutions. As shown in [7] such a strong form is false for p in the sub-critical range (1.3).

It turns out that (3.7) is actually sufficient to establish that any locally bounded, possibly of variable sign, local, weak solutions of (1.1–1.2) for  $1 , is locally Hölder continuous in <math>E_T$ . In turn, such a Hölder continuity permits one to improve the lower bound in (3.7) by estimating  $f(\cdot)$  to a power of its argument, as indicated in (1.7).

#### 4. The first form of the Harnack inequality implies the Hölder continuity of u

Let u be a locally bounded, possibly of variable sign, local, weak solution of (1.1-1.3) in  $E_T$ . It is shown in [5] (Chap. IV, Proposition 2.1 and Lemma 2.1), that u is locally Hölder continuous in  $E_T$  if there exist constants  $\theta \in (0, 1)$  and C, A > 1, depending only upon the data and independent of u, such that, for every  $(x_o, t_o) \in E_T$ , constructing the sequences

$$R_o = R$$
,  $R_n = \frac{R}{C^n}$ ;  $\omega_o = \omega$ ,  $\omega_{n+1} = \theta \omega_n$  for  $n = 0, 1, 2, ...$ 

for positive R and  $\omega$ , and the cylinders

$$Q_n = B_{R_n}(x_o) \times \left(t_o - \left(\frac{\omega_n}{A}\right)^{2-p} R_n^p, t_o\right] \quad \text{for } n = 1, 2, \dots$$

there holds

$$Q_{n+1} \subset Q_n \subset Q_o \subset E_T$$
 and ess osc  $u \leq \omega_n$ .

We will show that (3.7) permits one to construct such sequences for an arbitrary  $(x_o, t_o) \in E_T$ . Having fixed  $(x_o, t_o) \in E_T$  assume it coincides with the origin of  $\mathbb{R}^{N+1}$  and for  $\rho > 0$  set

$$R_o = 4\rho$$
 and  $Q = B_{4\rho} \times (-(4\rho)^p, 0]$  (4.1)

where  $\rho$  is so small that  $Q \subset E_T$ . Set also

$$\mu_o^+ = \operatorname{ess \, sup} u, \quad \mu_o^- = \operatorname{ess \, inf} u, \quad \omega_o = \mu_o^+ - \mu_o^- = \operatorname{ess \, osc} u.$$

Since u is locally bounded in  $E_T$ , without loss of generality we may assume that  $\omega_o \le 1$  so that

$$Q_o \stackrel{\text{def}}{=} B_{4\rho} \times \left( -\left(\frac{\omega_o}{A}\right)^{2-p} (4\rho)^p, 0 \right] \subset Q \subset E_T \quad \text{and } \operatorname{ess osc} u \leq \omega_o$$

for a number  $A \ge 1$  to be chosen. Now set

$$\mu^+ = \operatorname{ess\,sup} u, \qquad \mu^- = \operatorname{ess\,inf} u, \qquad \bar{\omega} = \operatorname{ess\,osc} u$$

and introduce the two functions defined in  $Q_o$ 

$$v_{+} = \mu^{+} - u, \quad v_{-} = u - \mu^{-}.$$

Without loss of generality may assume that

$$\mu^{+} - \frac{1}{4}\omega_{o} \ge \mu^{-} + \frac{1}{4}\omega_{o}. \tag{4.2}$$

Indeed otherwise  $\bar{\omega} \leq \frac{1}{2}\omega_o$  and thus passing from Q to any smaller cylinder the essential oscillation of u is reduced by a factor  $\frac{1}{2}$ , and there is nothing to prove. Then either

$$\left| \left[ v_{-}(\cdot, 0) \ge \frac{1}{4}\omega_{o} \right] \cap B_{\rho} \right| \ge \frac{1}{2}|B_{\rho}| \quad \text{or}$$

$$\left| \left[ v_{+}(\cdot, 0) \ge \frac{1}{4}\omega_{o} \right] \cap B_{\rho} \right| > \frac{1}{2}|B_{\rho}|.$$

$$(4.3)$$

Indeed by virtue of (4.2)

$$\left[u \leq \mu^+ - \frac{1}{4}\omega_o\right] \cap B_\rho \supset \left[u \leq \mu^- + \frac{1}{4}\omega_o\right] \cap B_\rho.$$

Therefore if the first of (4.3) is violated, then

$$\left| \left[ u \le \mu^+ - \frac{1}{4} \omega_o \right] \cap B_\rho \right| > \frac{1}{2} |B_\rho|.$$

Compute and estimate the values  $\delta_{\pm}$ , as defined by (1.4), relative to the functions  $v_{\pm}$ , over  $B_{\rho}$  at the time level t=0. Assuming the first of (4.3) holds

$$\omega_o^q \ge \frac{1}{|B_\rho|} \int_{B_\rho} \left( u(\cdot, 0) - \mu^- \right)^q dx$$

$$\ge \frac{1}{|B_\rho|} \int_{B_\rho \cap \{v_- > \frac{1}{2}\omega_0\}} [u(\cdot, 0) - \mu^-]^q dx \ge \frac{1}{2} \left( \frac{\omega_o}{4} \right)^q.$$

Therefore if the first of (4.3) holds

$$\frac{1}{2^{\frac{2-p}{q}}} \left(\frac{\omega_o}{4A_o}\right)^{2-p} \rho^p \le \delta_- \le \left(\frac{\omega_o}{A_o}\right)^{2-p} \rho^p \quad \text{for } A_o^{-1} = \varepsilon$$
 (4.4)

and there holds the inclusion

$$B_{4\rho} \times (-\delta_-, 0] \subset B_{4\rho} \times \left(-\left(\frac{\omega_o}{A_o}\right)^{2-p} \rho^p, 0\right].$$

Similar estimates hold for  $\delta_+$  if the second of (4.3) is in force. By the structure conditions (1.2) both  $v_\pm$  are solutions of (1.1–1.6) for the same constants  $C_o$  and  $C_1$  and hence the Harnack-type inequality (3.7) holds for either  $v_-$  or  $v_+$ , i.e.,

$$\inf_{Q_{4\rho}(\frac{1}{4}\delta_{\pm})} v_{\pm} \ge f(\eta_{\pm}) \sup_{Q_{2\rho}(\frac{1}{2}\delta_{\pm})} v_{\pm}. \tag{4.5}$$

where  $\eta_{\pm}$  are defined as in (1.5) for  $v_{\pm}$ . By virtue of (4.4), which holds for either  $\delta_{-}$  or  $\delta_{+}$ , and Remark 3.2, the function  $f(\cdot)$  can be taken to be the same. Assume now that the first of (4.3) holds true. Then as shown before

$$\int_{B_\rho} v_-^q(\cdot,0) dx \geq \frac{1}{|B_\rho|} \int_{B_\rho \cap [v \geq \frac{1}{4}\omega_\rho]} v_-^q(x,0) dx \geq \frac{1}{2} \left(\frac{\omega_\rho}{4}\right)^q.$$

On the other hand

$$\int_{B_{4o}} v_-^r(x, -\delta_-) dx \le \omega_o^r$$

and therefore recalling the definition (1.5) of  $\eta_{-}$ 

$$f(\eta_{-}) \ge f\left(\left(\frac{1}{2^{1/q}4}\right)^{\frac{pr}{\lambda_r}}\right) \stackrel{\text{def}}{=} 1 - \theta$$

for  $\theta \in (0, 1)$  depending only on the data and q and r. This and (4.5) imply

$$\inf_{B_{4\rho} \times (-\frac{1}{4}\delta, 0]} v_{-} \ge (1 - \theta) \sup_{B_{2\rho} \times (-\frac{1}{2}\delta, 0]} v_{-}$$
(4.6)

from which

$$\operatorname{ess} \operatorname{osc} u \leq \omega_1 \stackrel{\text{def}}{=} \theta \omega_0$$

where

$$Q_1 = B_\rho \times \left( -\left(\frac{\omega_o}{A}\right)^{2-p} \rho^p, 0 \right] \quad \text{and } A = 2^{1/q} \, 4^{1+\frac{1}{2-p}} A_o.$$

This and (4.4) determine A depending only upon the data and q, r. Taking into account (4.1) the cylinder  $Q_1$  is determined from  $Q_o$  by the indicated choice of A and for C = 4. A similar argument holds if the second of (4.3) is in force. This process can now be iterated and continued to yield:

**Proposition 4.1.** Let u be a locally bounded, local, weak solution of (1.1-1.2) for  $1 , in <math>E_T$ . There exist constants  $\bar{\gamma} > 1$  and  $\epsilon_o \in (0, 1)$ , depending only upon the data and r and q, such that for all  $(x_o, t_o) \in E_T$ , setting

$$M = \underset{(x_o, t_o) + Q_R(R^p)}{\operatorname{ess \, sup}} u \quad for \ (x_o, t_o) + Q_R(R^p) \subset E_T, \tag{4.7}$$

there holds

$$\operatorname*{ess\,osc}_{(x_o,t_o)+Q_\rho(\delta_M)} u \leq \bar{\gamma} M \left(\frac{\rho}{R}\right)^{\epsilon_o} \quad where \ \delta_M = \left(\frac{M}{A}\right)^{2-p} \rho^p \qquad (4.8)$$

for all  $0 < \rho < R$  and all cylinders

$$(x_o, t_o) + Q_o(\delta_M) \subset (x_o, t_o) + Q_R(R^p) \subset E_T$$
.

Remark 4.1. Returning to Remark 1.2, the previous arguments show that either  $\eta_+$  or  $\eta_-$  are bounded below by an absolute, positive constant  $\eta_o$ . Thus (4.5) implies that either  $\mu^+ - u$  or  $u - \mu^-$  satisfy a strong form of the Harnack Inequality. By the results of [7] a strong form of the Harnack estimate need not hold simultaneously for  $\mu^+ - u$  and  $u - \mu^-$ .

#### 5. Proof of Theorem 1.1 concluded

Assume  $(x_o, t_o)$  coincides with the origin of  $\mathbb{R}^{N+1}$ . Returning to (3.3) observe that by (3.2) and the same argument leading to (3.4)

$$|[u(\cdot,0) > \nu \eta M] \cap B_{\rho}| \ge \alpha \eta^q |B_{\rho}|$$
 and  $\sup_{B_{2\rho} \times (-\frac{1}{2}\delta,0]} u \le M$ 

for the same values of  $\nu$  and  $\alpha$  and with  $\delta$  given by (3.7). Since u is locally Hölder continuous, there exists  $x_1 \in B_{\rho}$  such that

$$u(x_1, 0) = v \eta M$$
.

Using the parameter A claimed by Proposition 4.1, construct the cylinder with "vertex" at  $(x_1, 0)$ 

$$(x_1,0) + Q_{2r} \left[ \left( \frac{v\eta M}{A} \right)^{2-p} r^p \right] \subset B_{2\rho} \times \left( -\frac{1}{4}\delta, 0 \right].$$

In view of (3.7) and the choice (4.4-4.6) of the parameter A, such an inclusion can be realized by possibly increasing A by a fixed quantitative factor depending only on the data, and by choosing r sufficiently small. Assuming the choice of r has been made, by Proposition 4.1

$$|u(x,t) - u(x_1,0)| \le \bar{\gamma} M \left(\frac{r}{\rho}\right)^{\epsilon_0}$$

for all

$$(x,t) \in \tilde{Q}_1 \stackrel{\text{def}}{=} (x_1,0) + Q_r \left[ \left( \frac{v \eta M}{A} \right)^{2-p} r^p \right].$$

From this

$$u(x,t) \ge \frac{1}{2} \nu \eta M$$
 for all  $(x,t) \in \tilde{Q}_1$ 

provided r is chosen to be so small that

$$\frac{\bar{\gamma}}{\nu\eta} \left(\frac{r}{\rho}\right)^{\epsilon_o} = \frac{1}{2} \quad \text{that is} \quad r = \varepsilon_1 \eta^{\frac{1}{\epsilon_o}} \rho \quad \text{where } \varepsilon_1 = \left(\frac{\nu}{2\,\bar{\gamma}}\right)^{\frac{1}{\epsilon_o}} \tag{5.1}$$

Therefore by Proposition 2.2

$$u \ge \sigma[\nu \eta M]$$
 in  $(x_1, 0) + Q_{2r} \left[ \left( \frac{\sigma[\nu \eta M]}{A} \right)^{2-p} (2r)^p \right]$ 

for  $\sigma \in (0, 1)$  depending only on A and p. This process can now be iterated to give

$$u \ge \sigma^n[\nu \eta M]$$
 in  $(x_1, 0) + Q_{2^n r} \left[ \left( \frac{\sigma^n[\nu \eta M]}{A} \right)^{2-p} (2^n r)^p \right]$ 

for all  $n \in \mathbb{N}$ . Choose n as the smallest integer for which

$$2^n r \ge 4\rho$$
 that is  $n \ge \log_2\left(\frac{4}{\varepsilon_1 \eta^{\frac{1}{\epsilon_0}}}\right)$ .

For such a choice

$$u \ge \gamma \eta^{\beta} M$$
 in  $Q_{2\rho} \left[ \left( \frac{\gamma \eta^{\beta} M}{A} \right)^{2-p} \rho^{p} \right]$ 

for some  $\beta = \beta(\text{data})$ .

#### 6. Equations of porous medium type

The techniques apply, by minor variants, to non-negative solutions of the class of quasi-linear, singular, parabolic equations of the porous-medium type. Precisely

$$u \in C_{\text{loc}}\left(0, T; L_{\text{loc}}^{2}(E)\right) \text{ such that } |u|^{\frac{m+1}{2}} \in L_{\text{loc}}^{2}\left(0, T; W_{\text{loc}}^{1,2}(E)\right)$$

$$u_{t} - \text{div } \mathbf{A}(x, t, u, Du) = 0 \text{ weakly in } E_{T}.$$

$$(6.1)$$

The functions  $\mathbf{A}: E_T \times \mathbb{R}^{N+1} \to \mathbb{R}^N$  are only assumed to be measurable and subject to the structure conditions

$$\begin{cases} \mathbf{A}(x, t, u, Du) \cdot Du \ge C_o |u|^{1-m} |D|u|^m|^2 \\ |\mathbf{A}(x, t, u, Du)| \le C_1 |D|u|^m| \end{cases}$$
 a.e. in  $E_T$  (6.2)

where  $C_o$  and  $C_1$  are given positive constants, and m is in the critical and sub-critical range

$$0 < m \le \frac{(N-2)_+}{N}. (6.3)$$

The homogeneous prototype of such a class is

$$u_t - \Delta |u|^{m-1} u = 0 \quad \text{weakly in } E_T. \tag{6.1}_o$$

For  $\tau > 0$  and for  $(y, s) \in \mathbb{R}^N \times \mathbb{R}$  set

$$Q_{\rho}(\tau) = B_{\rho} \times (-\tau, 0], \qquad (y, s) + Q_{\rho}(\tau) = B_{\rho}(y) \times (s - \tau, s].$$

Let *u* be a non-negative weak solution of (6.1–6.3). Having fixed  $(x_o, t_o) \in E_T$ , and  $B_{4\rho}(x_o) \subset E$ , introduce the quantities

$$\int_{B_{\rho}(x_o)} u^q(x, t_o) dx, \quad \delta \stackrel{\text{def}}{=} \left[ \varepsilon \left( \int_{B_{\rho}(x_o)} u^q(\cdot, t_o) dx \right)^{\frac{1}{q}} \right]^{1-m} \rho^2 \tag{6.4}$$

where  $\varepsilon \in (0, 1)$  is to be chosen, and  $q \ge 1$  is arbitrary. If  $\delta > 0$ , set also

$$\eta \stackrel{\text{def}}{=} \left[ \frac{\left( f_{B_{\rho}(x_o)} u^q(\cdot, t_o) dx \right)^{\frac{1}{q}}}{\left( f_{B_{4\rho}(x_o)} u^r(\cdot, t_o - \delta) dx \right)^{\frac{1}{r}}} \right]^{\frac{2r}{\lambda_r}}$$
(6.5)

where  $r \ge 1$  is any number such that

$$\lambda_r \stackrel{\text{def}}{=} N(m-1) + 2r > 0.$$
 (6.6)

**Theorem 6.1.** Let u be a non-negative, locally bounded, local, weak solution of (6.1–6.3). Introduce  $\delta$  as in (6.4) and assume that  $\delta > 0$ . There exist constants  $\varepsilon \in (0, 1)$ , and  $\gamma > 1$ , depending only on the data and the parameters q, r, and a constant  $\beta > 1$ , depending only upon the data and independent of q, r, such that

$$\inf_{(x_{o},t_{o})+Q_{\rho}(\frac{1}{2}\delta)} u \geq \gamma \left[ \frac{\left( \int_{B_{\rho}(x_{o})} u^{q}(\cdot,t_{o}) dx \right)^{\frac{1}{q}}}{\left( \int_{B_{4\rho}(x_{o})} u^{r}(\cdot,t_{o}-\delta) dx \right)^{\frac{1}{r}}} \right]^{\beta \frac{2r}{\lambda_{r}}} \sup_{(x_{o},t_{o})+Q_{\rho}(\delta)} u \quad (6.7)$$

provided  $q \ge 1$  and  $r \ge 1$  satisfies (6.6) and  $(x_o, t_o) + Q_{8\rho}(\delta) \subset E_T$ . The constants  $\varepsilon \to 0$ , and  $\gamma \to \infty$  as either  $\lambda_r \to 0$  or as  $\lambda_r \to \infty$ .

Remark 6.1. An estimate similar to (6.7) has been derived in [2] for non-negative solutions of the prototype equation  $(6.1)_o$ , by means of maximum and comparison principles. The arguments for the classes  $(1.1)_o$  and  $(6.1)_o$  are conceptually and technically similar.

Remark 6.2. Inequality (6.7) actually holds for non-negative solutions of (6.1–6.2) for all 0 < m < 1, provided  $r \ge 1$  satisfies (6.6). For super-critical values of  $m > \frac{(N-2)_+}{N}$  one has  $\lambda = \lambda_1 > 0$ , and (6.6) can be realized for r = 1. However, for  $\lambda > 0$  the strong form of a Harnack estimate holds ([7]). Therefore (6.7), while true for all 0 < m < 1, holds significance only for critical and sub-critical values  $0 < m \le \frac{(N-2)_+}{N}$ . In this sense (6.7) can be regarded as a "weak" form of a Harnack estimate. Neverthless it can be shown (6.7) is sufficient to establish the local Hölder continuity of locally bounded, weak solutions of (6.1–6.2), irrespective of their sign.

### **Appendix**

#### A. Proof of Propositions 2.1 and 2.3

**Proposition A.1.** Let u be a non-negative, locally bounded, local, weak solution of (1.1-1.2) for  $1 . For every <math>r \ge 1$  satisfying (1.6), there exists a positive constant  $\bar{\gamma}_r$ , depending only upon the data and r, such that for all  $B_{2\rho}(y) \times [2s-t, t] \subset E_T$ , for s < t

$$\sup_{B_{\rho}(y)\times[s,t]} u \leq \tilde{\gamma}_r \left(\frac{\rho^p}{t-s}\right)^{\frac{N}{\lambda_r}} \left(\frac{1}{\rho^N(t-s)} \int_{2s-t}^t \int_{B_{2\rho}(y)} u^r dx d\tau\right)^{\frac{p}{\lambda_r}} + \tilde{\gamma}_r \left(\frac{t-s}{\rho^p}\right)^{\frac{1}{2-p}}. \tag{A.1}$$

The proof is in [5] Chap. V.

#### A.1. Proof of Proposition 2.3

If r=1 this follows from Proposition 2.4. Assume r>1, take (y,t)=(0,0), fix  $\sigma \in (0,1]$  and let  $\zeta$  be a non-negative piecewise smooth cutoff function in  $\mathbb{R}^N$  vanishing outside  $B_{(1+\sigma)\rho}$  and satisfying

$$0 \leq \zeta \leq 1 \text{ in } B_{(1+\sigma)\rho}; \quad \zeta = 1 \text{ in } B_{\rho}; \quad |D\zeta| \leq \frac{C}{\sigma\rho} \text{ in } B_{(1+\sigma)\rho}.$$

In the weak formulation of (1.1–1.2), take the testing function  $u^{r-1}\zeta^p$ , modulo a standard Steklov time averaging process, and integrate over the cylinder  $Q = B_{(1+\sigma)\rho} \times (0, s]$ . This gives

$$\begin{split} &\frac{1}{r} \iint_{Q} \zeta^{p} u_{\tau}^{r} dx d\tau + (r-1) \iint_{Q} u^{r-2} \zeta^{p} \mathbf{A}(x, t, u, Du) \cdot Du dx d\tau \\ &+ p \iint_{Q} \zeta^{p-1} u^{r-1} \mathbf{A}(x, t, u, Du) \cdot D\zeta dx d\tau = \frac{1}{r} T_{1} + (r-1)T_{2} + T_{3}. \end{split}$$

Compute

$$T_1 = \int_{B_{(1+\sigma)\rho}} u^r(x,s) \zeta^p dx - \int_{B_{(1+\sigma)\rho}} u^r(x,0) \zeta^p dx$$

and estimate

$$\iint_{Q} u^{r-2} \zeta^{p} \mathbf{A}(x, t, u, Du) \cdot Du \, dx d\tau \ge C_{o} \iint_{Q} u^{r-2} \zeta^{p} |Du|^{p} dx d\tau$$
$$|T_{3}| \le C_{o}(r-1) \iint_{Q} \zeta^{p} u^{r-2} |Du|^{p} dx d\tau + \frac{C}{(\sigma \rho)^{p}} \iint_{Q} u^{p-2+r} dx d\tau$$

for a constant C depending only upon the data and r and such that  $C \to \infty$  if either  $\lambda_r \to 0$  or  $\lambda_r \to \infty$ . Combining these estimate yields

$$\begin{split} \sup_{0 \leq s \leq t} \int_{B_{\rho}} u^r(x,s) dx &\leq \int_{B_{(1+\sigma)\rho}} u^r(x,0) dx + \frac{C}{(\sigma\rho)^p} \iint_{Q} u^{p-2+r} dx d\tau \\ &\leq \int_{B_{(1+\sigma)\rho}} u^r(x,0) dx + \frac{C}{\sigma^p} \left(\frac{t^r}{\rho^{\lambda_r}}\right)^{\frac{1}{r}} \left(\sup_{0 \leq s \leq t} \int_{B_{(1+\sigma)\rho}} u^r dx\right)^{\frac{p-2+r}{r}}. \end{split}$$

The proof is concluded by a standard interpolation argument as in Lemma 4.3 of Chap. I of [5].

#### A.2. Proof of Proposition 2.1

The proof of Proposition 2.1 follows by combining (A.1) and Proposition 2.3.

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