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A compact and efficient strontium oven for laser-cooling experiments

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Here we describe a compact and efficient strontium oven well suited for laser-cooling experiments. Novel design solutions allowed us to produce a collimated strontium atomic beam with a flux of $1.0 \times 10^{13} \text{ s}^{-1} \text{ cm}^{-2}$ at the oven temperature of 450°C , reached with an electrical power consumption of 36 W. The oven is based on a stainless-steel reservoir, filled with 6 g of metallic strontium, electrically heated in a vacuum environment by a tantalum wire threaded through an alumina multi-bore tube. The oven can be hosted in a standard DN40CF cube and has an estimated continuous operation lifetime of 10 years. This oven can be used for other alkali and alkaline earth metals with essentially no modifications. © 2012 American Institute of Physics. [<http://dx.doi.org/10.1063/1.4756936>]

I. INTRODUCTION

Laser cooling is today an established technique to produce samples of atoms with temperature approaching the zero Kelvin limit. Ultra-cold strontium samples are largely studied, with experiments ranging from quantum degeneracy, quantum computation, to applications as quantum sensors of force and highly accurate optical clocks. These experiments are all based on an oven, which provides the needed flux of atoms in an ultra high vacuum (UHV) environment (pressure $<10^{-9}$ mbar). At room temperature the vapor pressure of alkaline earth metals (as strontium) is lower than that of alkali at the same temperature, for this reason relatively high temperature ovens are typically needed ($T \gtrsim 450^\circ\text{C}$ for Sr). In particular the application of alkaline earth metals for optical clocks is today prompting significant scientific and technological efforts for the realization of compact, reliable, and transportable apparatus, in the perspective of their future use in space.¹⁻⁴ This background is at the basis of our study on an efficient strontium oven, which represents one of the critical parts in an optical clock in terms of power consumption, size, and heat management.

Here we describe a compact and efficient strontium oven essentially based on a heater placed in vacuum. The novelty of the proposed oven concerns the synthesis of the solutions proposed so far for high efficient atomic beam sources (for a general review of atomic sources, see Ref. 5), such as capillaries and thermal shield, into a compact design, for both horizontal and vertical operations, without the complication of independent heaters for capillaries and crucible and without water cooling of the thermal shield. Additionally, reduced size and high efficiency (high flux at lower temperature and power consumption) imply a simplified heat management of the test chamber where clock spectroscopy is performed, with important consequences on the uncertainty budget of Sr clock measurement, nowadays limited by the lack of control of the

homogeneity of the blackbody environment.⁶⁻⁸ The use of this oven can be extended to other alkali and alkaline earth metals not reacting with stainless-steel at high temperature.

II. OVEN DESIGN

The oven design is shown in Fig. 1. About 6 g of metallic strontium is placed in a stainless-steel reservoir, with internal length and diameter, respectively, of 34 mm and 10 mm. The reservoir is electrically heated in vacuum by a 0.3 mm diameter tantalum wire threaded through an alumina (AL23) multi-bore tube. This material has been chosen since it ensures both a high electrical insulation (specific resistance $10^{10} \Omega \text{ cm}$ at 500°C) and a good thermal conductivity ($11 \text{ W m}^{-1} \text{ K}^{-1}$ at 500°C). The reservoir is closed by a DN16CF flange (removable for strontium refilling) sealed by a standard nickel gasket. The opposite side of the reservoir is terminated by a nozzle (diameter 4 mm, length 8 mm) filled with $N_{\text{cap}} \simeq 120$ stainless-steel capillaries $L = 8$ mm long and with an internal radius $a = 100 \mu\text{m}$. The high temperature assembly “reservoir + closing flange” is supported by three stainless-steel rods (diameter 3 mm) welded on a DN40CF flange. The thermal contact between the high temperature assembly and the supporting rods is reduced by using steatite split bushes UHV cleaned, having a thermal conductivity of $2.5 \text{ W m}^{-1} \text{ K}^{-1}$ at 500°C , about one order of magnitude below that of 304 stainless-steel at the same temperature. The oven can be hosted in a standard DN40CF cube, which can be connected to an ion pump and to a generic vacuum system for laser-cooling experiments. The vacuum environment prevents the convection mechanism of heat transfer. The remaining heat transfer channel is represented by the thermal radiation (blackbody radiation) of the hot reservoir at 450°C . In order to shield this radiation, the oven is surrounded by a 1 mm thick aluminum cylinder. Aluminum has been chosen for its low emissivity coefficient $\varepsilon \simeq 0.2$.

Electrical power is applied by means of two UHV electrical feedthroughs. Here two stainless-steel clamps are used to electrically connect the two feedthrough pins to the tantalum

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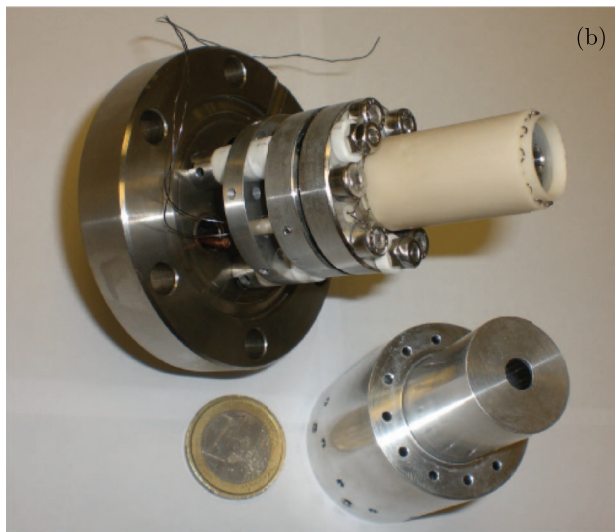
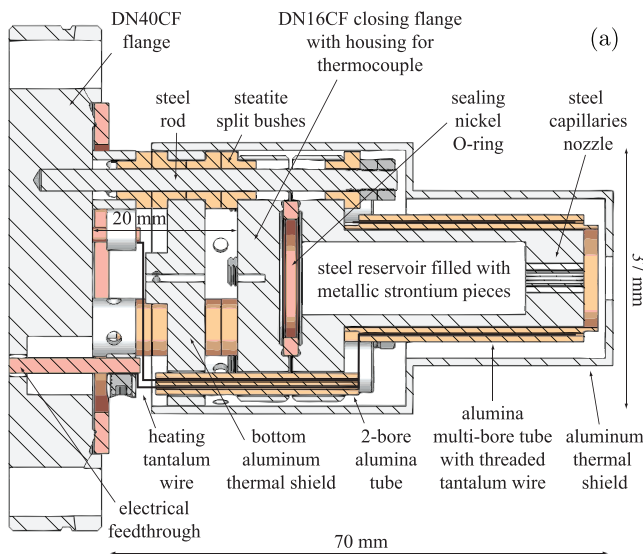


FIG. 1. (a) Section view of the compact strontium oven showing the internal components and layout. (b) Picture of the oven with the aluminum thermal shield. For the sake of comparison, a one euro coin is displayed.

heating wire. Short circuits are prevented by beading the two tantalum wire terminations with alumina tubes (0.4 mm internal diameter). Tantalum wires go through the DN16CF flange and reservoir and finally into the multi-bore tube heater by means of an alumina tube 28 mm long and with two 0.8 mm diameter holes.

The temperature of the oven is measured by a thermocouple in thermal contact with the rear side of the DN16CF flange. An aluminum plate is placed between the DN40CF and the DN16CF flange to shield the blackbody radiation. The thermocouple is electrically connected through the DN40CF flange by means of an UHV feedthrough.

III. EXPERIMENTAL RESULTS

A. Thermal properties

In our experimental setup (see Fig. 2), the oven is placed in a custom vacuum chamber evacuated by a 40 l/s ion pump. An aperture with radius $r = 2$ mm is placed at a distance

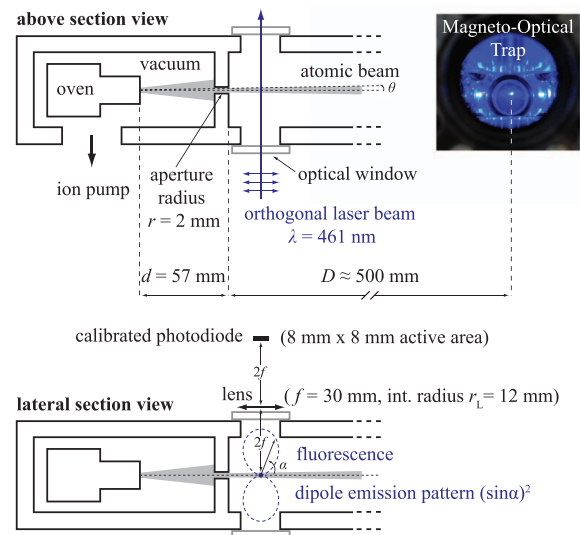


FIG. 2. Schematic view of the experimental setup. The oven is placed in vacuum and the atomic beam is further collimated by an aperture with radius r placed at distances d and D , respectively, from the nozzle and magneto-optical trapped atoms. An orthogonal laser beam resonant with the strontium transition at $\lambda = 461$ nm and horizontally polarized is used to estimate flux and divergence of the atomic beam by measuring the atomic fluorescence.

$d = 57$ mm from the nozzle of the oven to further increase the collimation of the atomic beam ($r/d \approx 35$ mrad) and to avoid the effusion of strontium toward the optical windows of the vacuum chamber. The equilibrium pressure in the vacuum chamber at the working oven temperature of 450°C is $\sim 10^{-8}$ mbar. In such environment, this temperature is reached with a power consumption of 36 W (2.38 A–15.2 V). Figure 3 shows that the measured equilibrium temperature T is not linear with the electrical power consumption P , since the thermal properties of the used materials depend on temperature and the loss of energy by thermal radiation increases with T^4 . However, it is possible to define a total

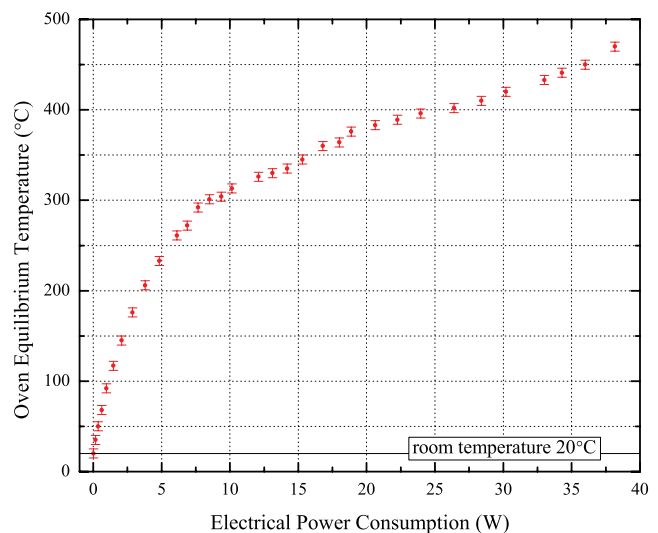


FIG. 3. Oven equilibrium temperature as a function of the electrical power consumption. The oven working temperature of 450°C is reached with 36 W of power consumption. The vertical error bars are given by the temperature uncertainty of the thermocouple ($\pm 5^\circ\text{C}$).

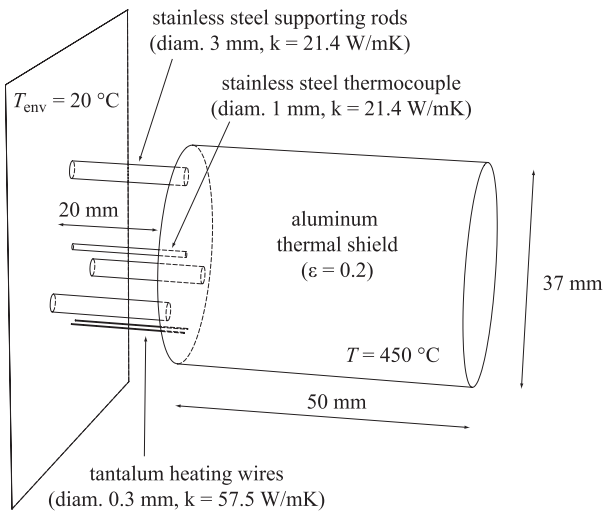


FIG. 4. Simplified thermal model of the proposed oven to trace out the origin and nature of the heat losses (see the text).

thermal resistance $R_{\text{th}}(T)$ given by

$$T - T_{\text{env}} = R_{\text{th}}(T) \times P, \quad (1)$$

where $T_{\text{env}} = 20^\circ\text{C}$ is the room temperature. At the working temperature of 450°C , we experimentally find a thermal resistance $R_{\text{th}} \simeq 12^\circ\text{C}/\text{W}$.

Maximum temperature is limited by the thermal shield in aluminum having the melting point at 660°C . However, by using stainless-steel it is possible to increase the maximum operating temperature at the $\sim 1000^\circ\text{C}$ level.

In order to evaluate whether the main heat losses are due to conduction or blackbody radiation (at 450°C), we consider a simplified thermal model of the oven (as shown in Fig. 4) with the aluminum thermal shield at 450°C in contact with the environment at 20°C (at distance of 20 mm) through three stainless-steel supporting rods, one stainless-steel thermocouple, and two tantalum heating wires. From the simple expression $P_{\text{cond}} = kA\Delta T/l$, where k is the thermal conductivity, A is the area of the conducting surface, $\Delta T = T - T_{\text{env}}$ is the temperature difference, and l is the thickness of the conducting surface separating the two temperatures, we find that the heat conduction loss due to the supporting rods is 10 W, thermocouple and heating wires together are at level of 0.5 W. The main loss is due to the blackbody radiation emitted by the aluminum thermal shield at 450°C , given by

$$P_{\text{rad}} = S_{\text{tot}}\sigma\varepsilon(T^4 - T_{\text{env}}^4) \simeq 24 \text{ W}, \quad (2)$$

where $S_{\text{tot}} = 80 \text{ cm}^2$ is total surface area of the thermal shield, σ is the Stefan-Boltzmann constant, and $\varepsilon \simeq 0.2$ is the emissivity of aluminum. The total estimated power loss is below 35 W, close to the measured value 36 W. The latter value can be now put in perspective and we can conclude that the thermal properties of the oven can be further improved by choosing for the thermal shield a material or a surface finishing (lapping or plating) with reduced emissivity ($\varepsilon < 0.2$).

Additionally, compact size and high efficiency of the oven ensure a better control on the blackbody radiation (BBR) coming from the hot furnace into the atoms trapping region. In our experimental configuration, atoms are magneto-

optical trapped at a distance $D \simeq 500 \text{ mm}$ from the cold (at room temperature) aperture, leading to a total BBR frequency shift due to the oven, on the strontium clock transition $5s^2\ ^1S_0 - 5s5p\ ^3P_0$, given in fractional units by^{9,10}

$$\left| \frac{\delta\nu_{\text{BBR}}^{(\text{oven})}}{\nu_0} \right| = 5.484 \times 10^{-15} \times \left(\frac{T}{300 \text{ K}} \right)^4 \times \frac{\Theta}{4\pi} < 10^{-18}, \quad (3)$$

where $\nu_0 \simeq 429 \text{ THz}$ is the frequency of the strontium clock transition, $T = 723 \text{ K}$ is the temperature of the oven, and $\Theta = \pi(r/D)^2$ is the solid angle under which the atoms see the cold aperture. Remarkably, this estimation makes the BBR frequency shift produced by the oven in this configuration completely negligible in the state of art strontium clocks corrections and uncertainty budget.⁶⁻⁸ The latter is nowadays dominated by the BBR effect of the environment at room temperature and a compact and thermally insulated oven can additionally simplify the heat management of the test chamber (where the clock spectroscopy is performed) and improve its temperature homogeneity with a great benefit in the final uncertainty budget of strontium clock.

B. Orthogonal crossed-beam spectroscopy for laser frequency stabilization

In order to characterize the atomic beam produced by the oven, spectroscopy on the strontium transition $5s^2\ ^1S_0 \leftrightarrow 5s5p\ ^1P_1$ at $\lambda = 461 \text{ nm}$ (saturation intensity $I_s = 42.7 \text{ mW/cm}^2$) is performed. A linearly polarized probe beam, with a diameter $\phi = 2.35 \text{ mm}$ ($1/e^2$) and an optical power of $165 \mu\text{W}$ (saturation parameter $s \equiv III_s \simeq 0.09$), is sent orthogonally with respect to the direction of the atomic beam (see Fig. 2). In this linear absorption regime with a single probe beam, the atomic fluorescence signal is measured as a function of the probe frequency (see Fig. 5). This spectrum is fitted with the sum of six Voigt functions, one for each bosonic isotope (^{84}Sr , ^{86}Sr , ^{88}Sr) and hyperfine components of the fermionic ^{87}Sr ($F = 9/2, 11/2, 7/2$). By knowing the natural linewidth ($\gamma = 32 \text{ MHz}$), the isotope shifts, hyperfine splittings, relative strengths, and the natural abundances,¹¹ it is possible to obtain from the fit the remaining five free parameters given by an overall amplitude and offset, a center frequency offset, the gaussian width (not calibrated), and more importantly the time/frequency scaling factor to calibrate in frequency our scan. With this calibration, we obtain a full-width-half-maximum of $(92.2 \pm 0.2) \text{ MHz}$ for the ^{88}Sr peak. This value essentially depends on the transverse velocity distribution of atoms (or equivalently to the collimation of the atomic beam). Indeed, the limited transit-time of atoms across the probe beam $\Delta t = \phi/v_{\text{beam}}$, with v_{beam} the most probable longitudinal speed in an effusive atomic beam,

$$v_{\text{beam}} = \sqrt{3k_B T/M} \simeq 450 \text{ m/s} \quad (T = 450^\circ\text{C}) \quad (4)$$

(k_B Boltzmann constant, M strontium mass), corresponds to a Fourier-limited frequency width $1/\Delta t$ of 190 kHz, that can be neglected. Now considering for simplicity just one isotope

peak, by assuming a normal transverse velocity distribution,

$$f_t(v_t) \propto \exp\left(-\frac{v_t^2}{2\sigma_t^2}\right), \quad (5)$$

where σ_t is the velocity standard deviation, the fluorescence signal is given by the Voigt profile,

$$F(\delta\omega) \propto \int_{-\infty}^{+\infty} \frac{f_t(v_t)}{(\delta\omega - |\mathbf{k}|v_t)^2 + \Gamma^2/4} dv_t, \quad (6)$$

where $\delta\omega = \omega - \omega_0$ is the angular frequency detuning from the resonance ω_0 , $|\mathbf{k}| = 2\pi/\lambda$ is the modulus of the probe beam wavevector (here we assumed a perfect orthogonality) and $\Gamma = 2\pi\gamma$ is the natural linewidth (in angular frequency). We find from the frequency axis calibration that the fit value of the standard deviation σ_t is (14.41 ± 0.04) m/s, from which we obtain the experimental value of the atomic beam collimation,

$$\theta \equiv \frac{\sigma_t \sqrt{2 \ln 2}}{v_{\text{beam}}} \simeq 37 \text{ mrad}, \quad (7)$$

where $\sigma_t \sqrt{2 \ln 2}$ is the half-width-half-maximum of the transverse velocity distribution in Eq. (5). The measured value of the beam collimation is very close to the angular selection given by the presence of the aperture ($r/d \simeq 35$ mrad).

Such collimated atomic beam is a good tool to frequency stabilize a laser at the MHz level on the strontium transition $5s^2 \ ^1S_0 \leftrightarrow 5s5p \ ^1P_1$ (error-signal in Fig. 5) and is thus an alternative to hollow cathode lamps, sealed vapor cells,^{12,13} and split-photodiode technique.¹⁴ As the transverse Doppler width is comparable to the natural linewidth, the Doppler-free saturation spectroscopy on the atomic beam has no advantage with respect to single beam excitation. For this application,

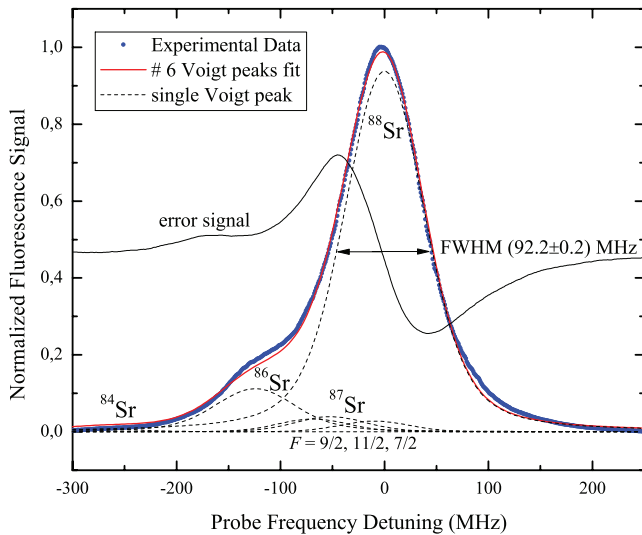


FIG. 5. Fluorescence spectrum of the strontium $5s^2 \ ^1S_0 \leftrightarrow 5s5p \ ^1P_1$ transition (blue dots) obtained in an orthogonal crossed-beam spectroscopy with a single probe beam. The fit function (red line) given by the sum of six Voigt functions (adjusted r-square of fit 0.998) is used to calibrate the frequency axis. Each Voigt peak (black dashed lines) represents the contribution of the three bosonic isotopes ^{84}Sr , ^{86}Sr , ^{88}Sr and the three hyperfine components of the fermionic ^{87}Sr . The error signal associated to the spectrum can be used to frequency stabilize the laser on the atomic transition.

the orthogonality between the probe and the atomic beam is critical to avoid any systematic shift $\delta\nu$ between the atomic resonance and the stabilized laser frequency. The sensitivity to this effect is given by the projection of the most probable longitudinal velocity onto the probing direction,

$$\delta\nu \sim \frac{v_{\text{beam}}}{\lambda} \varphi \simeq 0.98 \text{ MHz/mrad}, \quad (8)$$

where φ ($\ll 1$ rad) represents the angular deviation from the perpendicular to the atomic beam. The condition of orthogonality (and zero systematic frequency shift) is experimentally achieved by maximizing the peak and minimizing the width of the fluorescence signal. Additionally, an estimation of the residual systematic shift $\delta\nu$ can be obtained from the width variation when the probe beam is retroreflected and is given by the simple expression,

$$\delta\nu \simeq (\Delta\nu_{\text{refl.}} - \Delta\nu_{\text{single}})/2, \quad (9)$$

where $\Delta\nu_{\text{refl.}}$ and $\Delta\nu_{\text{single}}$ represent, respectively, the widths of the fluorescence signal with and without retro-reflection. In our experimental configuration, we find a systematic shift (modulus) lower than 2 MHz, well below the natural linewidth of the strontium $5s^2 \ ^1S_0 \leftrightarrow 5s5p \ ^1P_1$ transition ($\gamma = 32$ MHz).

C. Atomic beam flux, effusion regime, and duration

In order to estimate the atomic flux, the fluorescence peak produced in the orthogonal spectroscopy is measured by means of a calibrated photodiode. The flux of atoms \mathcal{F} (atoms $\text{s}^{-1} \text{cm}^{-2}$) is given by

$$\mathcal{F} = \rho v_{\text{beam}} \quad (10)$$

with ρ atomic density. The latter is related to the measured fluorescence optical peak power \mathcal{P}_{max} (W) through

$$\rho = \frac{1}{V_{\text{int}}} \times \mathcal{P}_{\text{max}} \times \frac{\Omega_{\text{tot}}}{\Omega_{\text{ph}}} \times \left(\frac{\hbar\omega_0\Gamma s}{2\sigma_t\sqrt{2\pi}} \int_{-\infty}^{+\infty} \frac{\exp(-v_t^2/2\sigma_t^2)}{1 + 4(|\mathbf{k}|v_t/\Gamma)^2} dv_t \right)^{-1}, \quad (11)$$

where $V_{\text{int}} \approx \pi(\phi/2)^2 2r$ is the interaction volume, $\Omega_{\text{tot}} = \int_0^\pi (\sin\alpha)^2 (2\pi \sin\alpha d\alpha) = 8\pi/3$ is total solid angle of dipole emission, $\Omega_{\text{ph}} = \pi(r_L/2f)^2$ is the solid angle under which the atomic fluorescence is seen by the photodiode (see Fig. 2) and \hbar is the reduced Planck's constant. At $T = 450^\circ\text{C}$, taking into account the contributions of all the strontium isotopes, we estimate a total flux of $\mathcal{F} \simeq 1.0 \times 10^{13} \text{ s}^{-1} \text{cm}^{-2}$, corresponding to a flow rate of $\dot{N} = \mathcal{F}(\pi r^2) \simeq 1.2 \times 10^{12} \text{ s}^{-1}$. Since from Eq. (7), we find that atomic beam divergence ($\theta \simeq 37$ mrad) is determined by the cold aperture ($r/d \simeq 35$ mrad) and not by the geometry of the capillaries ($\theta_t \equiv 1.68 a/L \simeq 21$ mrad¹⁵), we can define an atomic beam intensity peak (atoms s^{-1} steradian $^{-1}$) given by

$$J \equiv \frac{\dot{N}}{\Phi}, \quad (12)$$

where $\Phi = \pi(r/d)^2$ is the solid angle selection performed by the cold aperture (see Fig. 2). In order to establish the

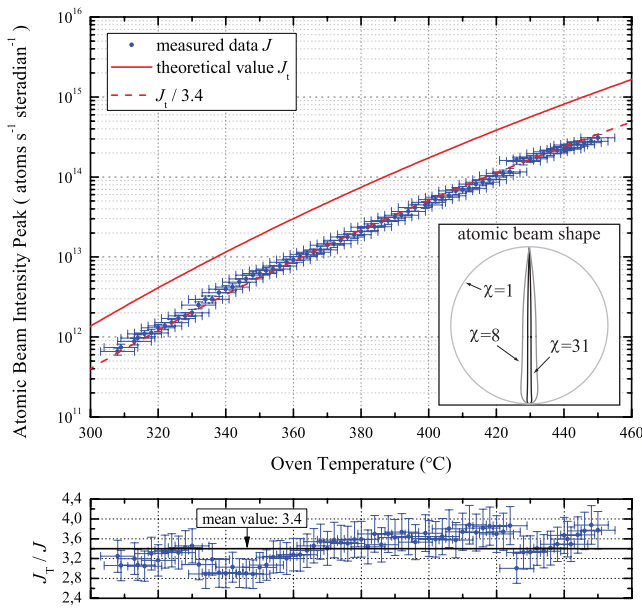


FIG. 6. Semilog plot of the measured atomic beam peak intensity (dots) as a function of the oven temperature, in comparison with the theoretical curve (solid curve) for a collision-free source. Experimental and theoretical values are in agreement except for a proportional factor of ~ 3.4 (see the text). The dashed line represents the theoretical value scaled by this factor. Horizontal error bars of $\pm 5^\circ\text{C}$ are given by the temperature uncertainty of the thermocouple, vertical error bars of 10% take into account the total uncertainty on the atomic beam intensity measurement (due to the calibration of the photodiode, temperature, and geometrical factors). The inset shows the polar plot of the atomic beam intensity as a function of the emission angle¹⁶ for different values of the peaking factor χ .

effusion regime, the atomic beam intensity peak is measured as a function of the oven temperature and is compared with the theoretical calculation in the limit of no collisions occurring in the capillaries (Knudsen regime),^{15,16}

$$J_t = \frac{(\pi a^2) \bar{v} n}{4\pi} N_{\text{cap}}, \quad (13)$$

where $\bar{v} = (8k_B T / \pi M)^{1/2}$ and $n = P / k_B T$ are, respectively, the mean atomic speed and the atomic density in the reservoir, with P vapor pressure of strontium given by¹⁷

$$\log(P(\text{Pa})) = 14.232 - \frac{8572}{T(\text{K})} - 1.1926 \log(T(\text{K})). \quad (14)$$

From the comparison, shown in Fig. 6, we find that the behavior of the experimental atomic intensity as a function of the oven temperature is in good agreement with the theoretical curve, except for a proportional factor of ~ 3.4 . We conclude that the atomic source is operating in the collision-free regime, as also confirmed by the calculation of the mean free path l of strontium atoms inside the reservoir (at $T = 450^\circ\text{C}$),

$$l = (\sqrt{2} \pi n \sigma^2)^{-1} \simeq 13 \text{ cm} \gg a, L, \quad (15)$$

with $\sigma = 4.3 \text{ \AA}$ the atomic diameter of strontium. The disagreement between the experimental and theoretical absolute values of the atomic beam intensity, together with the evidence that the measured atomic beam divergence is larger

than the theoretical Knudsen regime value θ_t^{15} can be both explained by a slightly imperfect mutual alignment of the capillaries. To understand more quantitatively if the reduced intensity is compliant with the measured increased divergence we can assume that for a small misalignment of the capillaries the number of atoms effused in the half-intensity solid angle $J\pi\theta^2$ is the same of ideal case $J_t\pi\theta_t^2$, leading to $J_t/J \sim (\theta/\theta_t)^2 \simeq 3$, which is very close to the measured value.

In our collision-free regime, the total flow rate of atoms emitted in the half solid angle by the oven ($T = 450^\circ\text{C}$) is given by¹⁵

$$\dot{N}_{\text{tot}} = \frac{2\pi n \bar{v} a^3}{3L} N_{\text{cap}} \simeq 1.2 \times 10^{14} \text{ s}^{-1}, \quad (16)$$

corresponding to a continuous operation lifetime (with 6 g of strontium) of about 10 years. The system started continuous operation about 2 years ago, it is presently running without appreciable degradation of the atomic flux and it is regularly used for loading $\sim 10^8$ atoms in a magneto-optical trap operating on the $5s^2 1S_0 \leftrightarrow 5s5p 1P_1$ transition.

Considering that the measured value of the peak intensity at 450°C is $J \simeq 3.1 \times 10^{14} \text{ s}^{-1} \text{ steradian}^{-1}$, we can evaluate the peaking factor¹⁶ of our source,

$$\chi \equiv \pi J / \dot{N}_{\text{tot}} \simeq 8, \quad (17)$$

where π is added so that $\chi = 1$ for an effusive cosine emitter source (ideal thin-walled orifice). The measured χ has to be compared with the theoretical value $\chi_t \equiv 1/W = 31$ (see inset of Fig. 6), where $W = (8a/3L)/(1 + 8a/3L)$ is the transmission probability (or Clausing factor) of the capillaries. We can conclude that because of the difficulty in aligning 120 capillaries with $L/a = 80$ the theoretical peaking factor $\chi_t = 31$ is not reached. The measured value $\chi \simeq 8$ can be realized by employing a significantly simplified capillaries nozzle with $L/a \simeq 20$.

IV. CONCLUSION

In summary, we have presented a compact (length 70 mm, diameter 37 mm) and efficient strontium oven, capable of producing a 37 mrad collimated atomic beam, with a flux of $1.0 \times 10^{13} \text{ s}^{-1} \text{ cm}^{-2}$ at the oven temperature of 450°C , reached with an electrical power consumption of 36 W. The latter value can be further reduced by employing a lower emissivity thermal shield ($\varepsilon < 0.2$). Reduced dimension and high efficiency simplify the heat management of the vacuum system and keep the contribution of BBR frequency shift due to the oven on the strontium clock transition below the 10^{-18} fractional level. Laser-spectroscopy on the strontium transition $5s^2 1S_0 \leftrightarrow 5s5p 1P_1$ at $\lambda = 461 \text{ nm}$ was performed to characterize the atomic beam. We have shown that the collimation of the atomic beam allows frequency stabilization of a 461 nm laser on the atomic transition with a negligible systematic shift, in a simple linear absorption regime spectroscopy with a single probe beam. Atomic beam intensity was measured as a function of the oven temperature and was compared with the theoretical model. From the comparison, we have concluded that the oven is operating in the collision-free regime. The simultaneous evidence that atomic

beam intensity and divergence are, respectively, slightly below and beyond the theoretical values can be explained by an imperfect mutual alignment of the capillaries. From the theoretical model, we have estimated a continuous operation lifetime of 10 years, experimentally we have not observed appreciable degradation of the atomic flux on a time scale of at least 2 years. The measured peaking factor of the oven ($\chi \simeq 8$) is compliant with a significantly simplified capillaries nozzle with $L/\lambda \simeq 20$.

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