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THE SEMI-INFINITE COHOMOLOGY OF WEYL MODULES WITH TWO SINGULAR POINTS

GIORGIA FORTUNA, DAVIDE LOMBARDO, ANDREA MAFFEI, VALERIO MELANI

A Claudio Procesi, con ammirazione.

"a volte due punti sono più vicini di quanto non sembri, ma per unirli ci vuole un'idea. Lui era la persona giusta." Nonmaterial lifeform, di F. B. Amadou, Urania.

ABSTRACT. In their study of spherical representations of an affine Lie algebra at the critical level and of unramified opers, Frenkel and Gaitsgory introduced what they called the Weyl module \mathbb{V}^{λ} corresponding to a dominant weight λ . This object plays an important role in the theory. In [4], we introduced a possible analogue $\mathbb{V}_2^{\lambda,\mu}$ of the Weyl module in the setting of opers with two singular points, and in the case of $\mathfrak{sl}(2)$ we proved that it has the 'correct' endomorphism ring. In this paper, we compute the semi-infinite cohomology of $\mathbb{V}_2^{\lambda,\mu}$ and we show that it does not share some of the properties of the semi-infinite cohomology of the Weyl module of Frenkel and Gaitsgory. For this reason, we introduce a new module $\mathbb{V}_2^{\lambda,\mu}$ which, in the case of $\mathfrak{sl}(2)$, enjoys all the expected properties of a Weyl module.

1. Introduction

Let $\mathfrak g$ be a complex simple Lie algebra and let $\hat{\mathfrak g}$ be its affinization. Choose a Borel subalgebra and a maximal toral subalgebra, and let G be a simply connected algebraic group with Lie algebra equal to $\mathfrak g$. As a particular case of a more general conjecture, Frenkel and Gaitsgory proved in [6] that the semi-infinite cohomology gives an isomorphism between the category $\hat{\mathfrak g}_{crit}$ -mod JG of spherical representations of $\hat{\mathfrak g}$ at the critical level (that is, representations of $\hat{\mathfrak g}$ at the critical level with a compatible action of $JG = G(\mathbb C[[t]])$) and the category of quasi-coherent sheaves on the space of unramified opers $\operatorname{Op}_1^{\operatorname{unr}}$ over $\mathfrak g^L$, the Langlands dual of $\mathfrak g$. As they explain, the space of unramified opers is the disjoint union of its connected components $\operatorname{Op}_1^{\lambda,\operatorname{unr}}$, and the category of spherical representations is the product of certain subcategories $\hat{\mathfrak g}_{crit}$ -mod $^{JG,\lambda}$, where in both cases λ ranges over all dominant weights of G. The equivalence given by semi-infinite cohomology specialises to an equivalence between $\hat{\mathfrak g}_{crit}$ -mod $^{JG,\lambda}$ and the category of quasi-coherent sheaves over $\operatorname{Op}_1^{\lambda,\operatorname{unr}}$. The space $\operatorname{Op}_1^{\lambda,\operatorname{unr}}$ is a non-reduced indscheme, and its reduced version, denoted by $\operatorname{Op}_1^{\lambda}$, is an affine scheme. In this paper we will denote by Z_1^{λ} its coordinate ring.

In this theory, an important role is played by the Weyl module \mathbb{V}_1^{λ} . This module enjoys the following fundamental properties:

$$\operatorname{End}_{\hat{\mathfrak{g}}}(\mathbb{V}_1^{\lambda}) \simeq Z_1^{\lambda} \quad \text{ and } \quad \Psi^0(\mathbb{V}_1^{\lambda}) \simeq Z_1^{\lambda},$$

where Ψ^n is the *n*-th semi-infinite cohomology group. Moreover the semi-infinite cohomology groups $\Psi^n(\mathbb{V}^{\lambda}_1)$ are trivial for $n \neq 0$.

Dennis Gaitsgory suggested to Giorgia Fortuna to study the space of unramified opers and spherical representations in a more general context, see [3]; in fact, the

definition of unramified opers as well as the definition of spherical representations can be generalized in the presence of more than one singularity, raising the question on whether or not certain statements remain true and what happens when these singularities collide.

In [4] we took some steps in this direction, by studying the case of $\mathfrak{sl}(2)$. In particular, we introduced a version of the Weyl module $\mathbb{V}_{2}^{\lambda,\mu}$ of critical level of the affine Lie algebra with two singularities $\hat{\mathfrak{g}}_{2}$. Thinking of t as a coordinate near the first singularity and s as a coordinate near the second singularity, this is the version of the affine Lie algebra over the ring $A = \mathbb{C}[[a]]$, where a = (t-s). As an A module is equal to $K_2 \otimes_{\mathbb{C}} \oplus AC_2$ where $K_2 = \mathbb{C}[[a,t]][1/t(t-a)]$ and C_2 is a central element (see [4], Section 3.3 for the complete definition).

We also introduced reduced scheme over A of unramified opers $\operatorname{Op}_2^{\lambda,\mu}$ which generalize the schemes Op_1^{λ} . Both objects depend on two integral dominant weights λ , μ of G, and we proved that

$$\operatorname{End}_{\hat{\mathfrak{g}}_2}(\mathbb{V}_2^{\lambda,\mu}) \simeq Z_2^{\lambda,\mu},$$

where $Z_2^{\lambda,\mu}$ is the coordinate ring of $\operatorname{Op}_2^{\lambda,\mu}$. In this article we study the semi-infinite cohomology of $\mathbb{V}_2^{\lambda,\mu}$ and its relation with the ring $Z_2^{\lambda,\mu}$ in order to understand how the equivalence $\Psi^{\bar{0}}(\mathbb{V}_1^{\lambda}) \simeq Z_1^{\lambda}$ generalizes. This is done in Section 4, where we compute the cohomology of $\mathbb{V}_2^{\lambda,\mu}$; in Section 5 we study the action of Z_2 , the center of a completion \hat{U}_2 of the enveloping algebra of $\hat{\mathfrak{g}}_2$ at the critical level on this module (see Section 2.2).

In particular, we prove that the specialisation at a = 0 and the localization at $a \neq 0$ of the semi-infinite cohomology of $\mathbb{V}_2^{\lambda,\mu}$ are isomorphic to the specialisation and localization of $Z_2^{\lambda,\mu}$, respectively. However, in contrast to our intuition, we also show the following result which says that $\Psi^0(\mathbb{V}_2^{\lambda,\mu})$ doesn't exactly generalize the equivalence $\Psi^0(\mathbb{V}_1^{\lambda}) \simeq Z_1^{\lambda}$ as expected:

Theorem A (Theorem 4.9 and Proposition 5.3). We have $\Psi^n(\mathbb{V}_2^{\lambda,\mu}) = 0$ for $n \neq 0$. Moreover, $\Psi^0(\mathbb{V}_2^{\lambda,\mu})$ is not isomorphic to $Z_2^{\lambda,\mu}$ as a Z_2 -module.

For this computation, we rely on the formalism introduced by Casarin in [1], which makes it possible to use vertex algebras also in the context of opers with two singularities. Once this formalism is in place, for the computation of the semiinfinite cohomology we can follow closely the approach taken by Frenkel and Ben Zvi in [5, Chapter 15] for the case of one singularity.

In the last section, we restrict our attention to the Lie algebra $\mathfrak{sl}(2)$ and introduce a submodule $\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}$ of $\mathbb{V}_{2}^{\lambda,\mu}$, which is generated by the highest weight vector. We prove that this module is the correct one to consider, in the sense that it has the expected cohomology groups and endomorphism ring, as the following result shows.

Theorem B (Proposition 6.3, Theorem 6.5 and Proposition 6.6). If $\mathfrak{g} = \mathfrak{sl}(2)$ then we have $\Psi^n(\widetilde{\mathbb{Y}}_2^{\lambda,\mu}) = 0$ for $n \neq 0$. Moreover, we have

$$\operatorname{End}_{\widehat{\mathfrak{g}}_2}(\widetilde{\mathbb{V}}_2^{\lambda,\mu}) \simeq Z_2^{\lambda,\mu} \quad \text{ and } \quad \Psi^0(\widetilde{\mathbb{V}}_2^{\lambda,\mu}) \simeq Z_2^{\lambda,\mu}.$$

We now briefly explain the connection between these results and Conjecture 3.6.1 in Fortuna's Thesis [3]. As a particular case the conjecture predicts an equivalence between quasi-coherent sheaves over the space of unramified opers with two singularities and the category of spherical representations over $\hat{\mathfrak{g}}_2$: that is the space of smooth representations of $\hat{\mathfrak{g}}_2$ with a compatible action of $J_2G = G(\mathbb{C}[[a,t]])$.

The conjecture stated in [3] predicts an equivalence of similar categories not only in the presence of two singularities but in the presence of n-possible singularities. In particular for any finite set with n elements I we can define the space of opers on the formal disc with n-singularities Op_I and the subspace of unramified opers $\operatorname{Op}_I^{\operatorname{unr}}$ (see Section 3.5 in [3]). These are spaces over the product of n-copies of the formal disc. These are easily seen to be factorization spaces, which means that this spaces specialise nicely when restricted along or outside the diagonals of this product (see Section 3.1.5 in [3]). There are not substantial differences between the treatment we do here or in [4] of Op_2 and the general case. The only minor difference is that we fix a singularity to be 0. These spaces are indschemes, and so we can define the categories $\operatorname{QCoh}(\operatorname{Op}_I)$, and $\operatorname{QCoh}(\operatorname{Op}_I^{\operatorname{unr}})$ of quasi-coherent sheaves on Op_I and $\operatorname{Op}_I^{\operatorname{ur}}$ (see Section 3.5.3 in [3] for the actual definition), and the nice factorization properties which make them factorization categories (see Section 3.1.2 in [3]).

Similarly, for a finite set I we can define a Lie algebra $\hat{\mathfrak{g}}_I$ and study its smooth representations at the critical level. The objects constructed in this way live also on the product of n copies of the formal disc, and they also have nice factorization properties, in particular the collection of (completions of the) enveloping algebras specialized at the critical level \hat{U}_I of the algebras $\hat{\mathfrak{g}}_I$, is what is called a factorization algebra (see Section 3.1.3 in [3]). As a consequence the collection of the categories of smooth representations at the critical of the Lie algebras $\hat{\mathfrak{g}}_I$, denoted by $\hat{\mathfrak{g}}_{I,crit}$ -mod and their subcategories of spherical representations $\hat{\mathfrak{g}}_{I,crit}$ -mod JG can be organized also in a factorization category. The semi-infinite cohomology can be defined also in this generality and defines a functor

$$\Psi_I : \hat{\mathfrak{g}}_{I,crit}\text{-mod} \longrightarrow D(\operatorname{QCoh}(\operatorname{Op}_I))$$

compatible with the factorization properties. While in Fortuna's thesis all these constructions are obtained somehow for free using the language of chiral algebras (see Section 3.1.6 in [3]), in this paper we use the language of vertex algebras and the formalism introduced by Casarin [1]. Let us notice that, from this point of view, there are no differences in treating the case with two singular points and the case with an arbitrary finite number of singular points. For example, the proof of Theorem A above can be repeated verbatim in the case of n singular points. More generally we believe that all the technical difficulties in the study of this problem already appear in the case of two singularities.

It is easy to see from the factorization properties and the analogous statement for the case of one singularty by Frenkel and Gaitsgory (see [7]) that the semi-infinite cohomology of a $\hat{\mathfrak{g}}_I$ -spherical module is supported on $\operatorname{Op}_I^{\operatorname{unr}}$. Hence semi-infinite cohomology restricts to a functor $\Psi_I: \hat{\mathfrak{g}}_{I,crit}$ -mod $^{JG} \longrightarrow D(\operatorname{QCoh}(\operatorname{Op}_I^{\operatorname{unr}}))$. Conjecture 3.6.1 in [3] states that this functor is exact and that

$$\Psi_I^0: \hat{\mathfrak{g}}_{I.crit}\text{-mod}^{JG} \longrightarrow \operatorname{QCoh}(\operatorname{Op}_I^{\operatorname{unr}})$$

is an equivalence of categories. In fact, it can be seen that the first part of Theorem A implies that Ψ_I is exact. Moreover, in the case of $\mathfrak{g} = \mathfrak{sl}(2)$, Theorem B yields that the restriction of Ψ_I^0 to modules with reduced support is an equivalence. The details will be given in a forthcoming paper.

The paper is organized as follows. In the first section we recall some definitions from [4]. In Section 3 we recall the formalism introduced by Casarin [1] and we use it to define semi-infinite cohomology and prove some of its basic properties. In Sections 3 and 4 we compute the semi-infinite cohomology of $\mathbb{V}_2^{\lambda,\mu}$ and in Section 5 we compute the semi-infinite cohomology of $\widetilde{\mathbb{V}}_2^{\lambda,\mu}$.

We thank Luca Casarin for many useful discussions and in particular for explaining to us the formalism introduced in [1]. It seems to us that Casarin's approach provides a natural framework to treat questions concerning opers with several singularities, making the theory much more transparent than it was in [4]. In particular,

the results of [1] allowed us to streamline several arguments and calculations which would have been quite hard to carry out using the direct approach of [4].

2. Basic constructions

In this section we recall some basic constructions from [4], to which we refer for further details, and we introduce the notion of semi-infinite cohomology in the context of affine Lie algebras with more than one singular point.

2.1. Rings. We follow [4, Section 1], to which the reader is referred for more details. We introduce the rings

$$A = \mathbb{C}[[a]], \qquad Q = \mathbb{C}((a)), \qquad R_2 = \mathbb{C}[[t, s]], \qquad K_2 = \mathbb{C}[[t, s]][1/ts],$$

where a = t - s. Recall that we have expansion maps (given by suitable natural inclusions) and a specialisation map (which sends a to 0 and t, s to t, see Section 1.1 in [4]

$$E_t: K_2[a^{-1}] \longrightarrow Q((t)), \qquad E_s: K_2[a^{-1}] \longrightarrow Q((s)), \qquad \operatorname{Sp}: K_2 \longrightarrow \mathbb{C}((t)).$$

We also write $E = E_t \times E_s : K_2[a^{-1}] \longrightarrow Q((t)) \times Q((s))$. Recall from [4, Section 1.1] that Sp induces an isomorphism $K_2/(a) \simeq \mathbb{C}((t))$. These rings have natural topologies: with respect to these, the image of E is dense, and $E(R_2[a^{-1}])$ is dense in $Q((t)) \times Q((s))$.

These rings are also equipped with residue maps

$$\operatorname{Res}_2: K_2 \to A \quad \operatorname{Res}_1: \mathbb{C}((t)) \to \mathbb{C}, \quad \operatorname{Res}_t: Q((t)) \to Q, \quad \operatorname{Res}_s: Q((s)) \to Q,$$

which behave nicely with respect to specialisation and expansion (see [4, Section 1.2). Finally, we recall Lemma 1.10 in [4].

Lemma 2.1 ([4], Lemma 1.10). Let M, N be two A-modules and $\varphi: M \longrightarrow N$ be a morphism of A-modules. Then

- a) if M is flat and $\varphi_a: M[a^{-1}] \longrightarrow N[a^{-1}]$ is injective, then φ is injective. b) if N is flat, $\varphi_a: M[a^{-1}] \longrightarrow N[a^{-1}]$ is surjective, and $\overline{\varphi}: M/aM \longrightarrow N/aN$ is injective, then φ is surjective.

In particular, if M and N are flat, $\varphi_a: M[a^{-1}] \longrightarrow N[a^{-1}]$ is an isomorphism, and $\overline{\varphi}: M/aM \longrightarrow N/aN$ is injective, then φ is an isomorphism.

2.2. Affine Lie algebras and completion of the enveloping algebra. We follow [4, Section 3]. Let g be a finite-dimensional Lie algebra over the complex numbers and denote by κ the Killing form of \mathfrak{g} . Recall from [4, Sections 3.1 and 3.3] that for each of the rings of the previous section we introduce an affine Lie algebra: $\hat{\mathfrak{g}}_1$ is the usual affine Lie algebra (we take for convenience the version defined by Laurent polynomial and not Laurent series), $\hat{\mathfrak{g}}_t$ and $\hat{\mathfrak{g}}_s$ are also versions of the usual affine Lie algebra, while $\hat{\mathfrak{g}}_2$ is an A-Lie algebra having as underlying A-module the space

$$\hat{\mathfrak{g}}_2 = \mathbb{C}[t,s][1/ts] \otimes_{\mathbb{C}} \mathfrak{g} \oplus AC_2.$$

We also introduce the Lie algebra $\hat{\mathfrak{g}}_{t,s} = \hat{\mathfrak{g}}_t \oplus \hat{\mathfrak{g}}_s/(C_t - C_s)$ (see [4, Section 3.3]).

For each of these Lie algebras, we introduce the corresponding universal enveloping algebra, which we suitably complete and then specialize at the critical level by imposing that the central element acts as -1/2 (see Sections 3.1 and 3.3 in [4]). In particular

$$\hat{U}_2 = \varprojlim_n \frac{U(\hat{\mathfrak{g}}_2)}{(C_2 = -1/2, \ t^n s^n \mathbb{C}[t, s] \otimes \mathfrak{g})_{\text{left.id.}}}$$

Recall from [4, Section 3.4] that the expansion maps and the specialisation maps induce morphisms at the level of Lie algebras. In particular, the specialisation map Sp: $\hat{U}_2 \longrightarrow \hat{U}_1$ induces an isomorphism between $\hat{U}_2/a\hat{U}_2$ and \hat{U}_1 , while the expansion map induces a morphism $E: \hat{U}_2[a^{-1}] \longrightarrow \hat{U}_{t,s}$ which is injective and has dense image.

Moreover, the natural inclusions $\hat{\mathfrak{g}}_t \hookrightarrow \hat{\mathfrak{g}}_{t,s}$ and $\hat{\mathfrak{g}}_s \hookrightarrow \hat{\mathfrak{g}}_{t,s}$ induce a morphism

$$\hat{U}_t \otimes \hat{U}_s \longrightarrow \hat{U}_t$$

which is also injective and with dense image (see [4, Section 3.3]).

2.3. Weyl modules. We follow [4, Section 6]. We choose a Borel subalgebra and a maximal toral subalgebra of \mathfrak{g} , which we denote by \mathfrak{b} and \mathfrak{t} respectively. This data induces a choice of weights, integral weights and dominant weights. For every integral dominant weight λ , [7] introduced the Weyl module \mathbb{V}_1^{λ} over the affine Lie algebra $\hat{\mathfrak{g}}_1$. The representation $\mathbb{V} = \mathbb{V}_1^0$, which has a structure of vertex algebra, will play a particularly important role for us. This vertex algebra enjoys the following universal property.

Lemma 2.2. Let U be a vertex algebra such that there exists a linear map $x \mapsto u_x$ from \mathfrak{g} to U such that

$$(u_x)_{(0)}(u_y) = u_{[x,y]}$$
 $(u_x)_{(1)}(u_y) = -\frac{1}{2}\kappa(x,y)|0\rangle_U$ $(u_x)_{(n)}(u_y) = 0$

for all $n \ge 2$. There exists a unique morphism of vertex algebras $\alpha : \mathbb{V} \to U$ such that $\alpha(xt^{-1}|0\rangle_{\mathbb{V}}) = u_x$ for all $x \in \mathfrak{g}$.

Weyl modules \mathbb{V}_t^{λ} and \mathbb{V}_s^{λ} can also be defined for the Lie algebras $\hat{\mathfrak{g}}_t$ and $\hat{\mathfrak{g}}_s$, without any significant change from [7]. In [4], we introduced a generalization of these modules. Given two dominant weights λ, μ , we consider the irreducible representations V^{λ} and V^{μ} of the Lie algebra \mathfrak{g} having highest weights λ, μ , respectively. In [4, Definition 6.2], given two dominant integral weights λ, μ we introduced the module

$$\mathbb{V}_{2}^{\lambda,\mu} = \operatorname{Ind}_{\hat{\mathfrak{g}}_{2}^{+}}^{\hat{\mathfrak{g}}_{2}} \left(A \otimes_{\mathbb{C}} V^{\lambda} \otimes_{\mathbb{C}} V^{\mu} \right),$$

where $\hat{\mathfrak{g}}_2^+ = \mathbb{C}[t,s] \otimes \mathfrak{g} \oplus A C_2$ acts on $A \otimes_{\mathbb{C}} V^{\lambda} \otimes_{\mathbb{C}} V^{\mu}$ as

$$f(t,s)x \cdot (p(a) \otimes u \otimes v) = f(0,-a)p(a) \otimes xu \otimes v + f(a,0)p(a) \otimes u \otimes xv,$$

while C_2 acts as -1/2. In [4] we called this object the Weyl module of weights (λ, μ) , although, as we will see, it does not have the same properties as its 1-singularity analogue.

We also define

$$\mathbb{W}_{1}^{\lambda,\mu} = \operatorname{Ind}_{\hat{\mathfrak{g}}_{1}^{+}}^{\hat{\mathfrak{g}}_{1}} \left(V^{\lambda} \otimes_{\mathbb{C}} V^{\mu} \right),$$

where $\hat{\mathfrak{g}}_1^+ = \mathbb{C}[t] \otimes \mathfrak{g} \oplus \mathbb{C} C_1$ acts on $V^{\lambda} \otimes_{\mathbb{C}} V^{\mu}$ as $f(t)x \cdot (u \otimes v) = f(0)x \cdot (u \otimes v)$ and C_1 acts as -1/2.

The specialisation and expansion maps are defined also for Weyl modules, and induce the following isomorphisms [4, Lemma 6.3]:

$$\frac{\mathbb{V}_{2}^{\lambda,\mu}}{a\mathbb{V}_{2}^{\lambda,\mu}} \simeq \mathbb{W}_{1}^{\lambda,\mu}, \qquad \mathbb{V}_{2}^{\lambda,\mu}[a^{-1}] \simeq \mathbb{V}_{t}^{\lambda} \otimes_{Q} \mathbb{V}_{s}^{\mu}. \tag{2.1}$$

2.4. Clifford algebra. We now define the Clifford algebra with two singularities, generalizing the construction of the classical case (see for example [5, Chapter 15]). Let \mathfrak{n}_+ be the nilpotent radical of \mathfrak{b} and set

$$X_2 = K_2 \otimes_{\mathbb{C}} \mathfrak{n}_+ \oplus K_2 \otimes_{\mathbb{C}} \mathfrak{n}_+^*.$$

We equip X_2 with the unique A-bilinear form such that $K_2 \otimes_{\mathbb{C}} \mathfrak{n}_+$ and $K_2 \otimes_{\mathbb{C}} \mathfrak{n}_+^*$ are isotropic subspaces and

$$(f \otimes x; q \otimes \varphi) = \operatorname{Res}_2(fq) \varphi(x)$$

for all $f, g \in K_2$, $x \in \mathfrak{n}_+$ and $\varphi \in \mathfrak{n}_+^*$. We denote by $\mathcal{C}\ell_2$ the associated Clifford algebra over A.

There are obvious variants of the same construction where we replace K_2 with the ring $\mathbb{C}[t^{\pm 1}]$ or one of the rings $Q[t^{\pm 1}]$, $Q[s^{\pm 1}]$, $Q[t^{\pm 1}] \times Q[s^{\pm 1}]$. We obtain Clifford algebras that we denote by $\mathcal{C}\ell_1, \mathcal{C}\ell_t, \mathcal{C}\ell_s, \mathcal{C}\ell_{t,s}$. The algebra $\mathcal{C}\ell_U$ in [5, Section 15.1.1] is a completion of $\mathcal{C}\ell_1$.

These Clifford algebras have a natural grading called the *charge* and denoted by ch. It can be defined as follows: the elements of the base ring have charge 0, while for $\psi \in \mathfrak{n}$ and $\psi^* \in \mathfrak{n}^*$ we have

$$\operatorname{ch} \psi = -1, \qquad \operatorname{ch} \psi^* = 1. \tag{2.2}$$

The relations defining each Clifford algebra are homogeneous, hence the charge induces a well-defined grading on the Clifford algebra.

We now introduce completions of the tensor product $\hat{U}_2 \otimes_A \mathcal{C}\ell_2$. We define

$$\hat{U}_2 \hat{\otimes}_A \mathcal{C} \ell_2 = \varprojlim_n \frac{\hat{U}_2 \otimes_A \mathcal{C} \ell_2}{\left((ts)^n R_2 \mathfrak{g} \otimes 1, \ 1 \otimes (ts)^n R_2 \mathfrak{n}_+, \ 1 \otimes (ts)^n R_2 \mathfrak{n}_+^* \right)_{\text{left ideal}}}$$

and we notice that, as in the case of the algebra \hat{U}_2 , this A-module has a natural structure of A-algebra. We introduce the completed Clifford algebras $\hat{U}_1 \hat{\otimes} \mathcal{C} \ell_1$, $\hat{U}_t \hat{\otimes}_Q \mathcal{C} \ell_t$, $\hat{U}_s \hat{\otimes}_Q \mathcal{C} \ell_s$, and $\hat{U}_{t,s} \hat{\otimes}_Q \mathcal{C} \ell_{t,s}$. The specialisation and expansion map determine morphisms

$$\operatorname{Sp}: \hat{U}_2 \hat{\otimes}_A \mathcal{C}\ell_2 \longrightarrow \hat{U}_1 \hat{\otimes} \mathcal{C}\ell_1 \qquad \text{and} \qquad E: (\hat{U}_2 \hat{\otimes}_A \mathcal{C}\ell_2)[a^{-1}] \longrightarrow \hat{U}_{t,s} \hat{\otimes}_Q \mathcal{C}\ell_{t,s}.$$

Arguing exacly as in [4, Lemmas 3.7 and 3.9] we see that E is injective with dense image, while the specialisation map induces an isomorphism $\hat{U}_2\hat{\otimes}_A\mathcal{C}\ell_2/a(\hat{U}_2\hat{\otimes}_A\mathcal{C}\ell_2)\simeq \hat{U}_1\hat{\otimes}\mathcal{C}\ell_1$. Finally, we have an injective map $I:\hat{U}_t\hat{\otimes}_Q\mathcal{C}\ell_t\to\hat{U}_{t,s}\hat{\otimes}_Q\mathcal{C}\ell_{t,s}$ induced by the natural inclusion $K_t\to K_{t,s}=K_t\times K_s$ given by $f\mapsto (f,0)$. Similarly, we have an injective map $J:\hat{U}_s\hat{\otimes}_Q\mathcal{C}\ell_s\to\hat{U}_{t,s}\hat{\otimes}_Q\mathcal{C}\ell_{t,s}$. As in Section 3.3 of [4], the product of these maps $I\otimes J:(\hat{U}_t\hat{\otimes}_Q\mathcal{C}\ell_t)\otimes_Q(\hat{U}_s\hat{\otimes}_Q\mathcal{C}\ell_s)\to\hat{U}_{t,s}\hat{\otimes}_Q\mathcal{C}\ell_{t,s}$ is injective with dense image.

2.5. Fock module. We now describe the "fermionic" Fock spaces corresponding to the Clifford algebras defined in the previous section. As above, for the construction in the case of one singularity we refer to [5, Section 15.1.4]: here we mimic this definition in the case of two singularities. We define $\mathcal{C}\ell_2^+$ as the A-subalgebra of $\mathcal{C}\ell_2$ generated by $R_2 \otimes \mathfrak{n}_+$ and $R_2 \otimes \mathfrak{n}_+^*$ and we define the Fock module

$$\Lambda_2 = \mathcal{C}\ell_2 \otimes_{\mathcal{C}\ell_2^+} A |0\rangle_{\Lambda_2}$$

where $R_2 \otimes \mathfrak{n}_+$ and $R_2 \otimes \mathfrak{n}_+^*$ acts trivially on $|0\rangle_{\Lambda_2}$. The charge (see equation (2.2)) induces a grading on the Fock space by setting

$$\operatorname{ch}|0\rangle_{\Lambda_2}=0.$$

We denote by Λ_2^n the subspace of homogeneous elements of charge equal to n. Similar constructions can be given for all the other Clifford algebras $\mathcal{C}\ell_1, \mathcal{C}\ell_t, \mathcal{C}\ell_s$, and $\mathcal{C}\ell_{t,s}$, giving Fock modules Λ_1 , Λ_t , Λ_s , and $\Lambda_{t,s}$.

Specialisation and expansion, induce maps also at the level of the Fock spaces. Arguing as in [4, Section 6] (where we considered the module $\mathbb{V}_2^{\lambda,\mu}$), it is easy to prove the following Lemma:

Lemma 2.3.

- a) The specialisation map $\operatorname{Sp}: \Lambda_2^{\bullet} \longrightarrow \Lambda_1^{\bullet}$ is homogeneous of degree zero and induces an isomorphism $\Lambda_2^{\bullet}/a\Lambda_2^{\bullet} \simeq \Lambda_1^{\bullet}$.
- b) We have a homogeneous isomorphism of degree zero $\Lambda_{t,s}^{\bullet} \simeq \Lambda_t^{\bullet} \otimes_Q \Lambda_s^{\bullet}$.

c) The expansion map $E: \Lambda_2^{\bullet}[a^{-1}] \longrightarrow \Lambda_t^{\bullet} \otimes_Q \Lambda_s^{\bullet}$ is a homogeneous isomorphism of degree zero.

Recall also that the Fock space $\Lambda = \Lambda_1$ has a natural structure of vertex superalgebra with the following universal property.

Lemma 2.4. Let U be a vertex superalgebra such that there exists a linear map $x \mapsto u_x$ from $\mathfrak{n}_+^* \oplus \mathfrak{n}_+^*$ to the space of odd elements of U such that

(1) for all $\varphi, \psi \in \mathfrak{n}$ and for all $\varphi^*, \psi^* \in \mathfrak{n}_+^*$

$$(u_{\psi})_{(n)}(u_{\varphi}) = (u_{\psi^*})_{(n)}(u_{\varphi^*}) = (u_{\psi})_{(m)}(u_{\psi^*}) = (u_{\psi^*})_{(m)}(u_{\psi}) = 0$$

for all $n \ge 0$ and for all $m \ge 1$;

(2) $(u_{\psi})_{(0)}(u_{\psi^*}) = (u_{\psi^*})_{(0)}(u_{\psi}) = \langle \psi, \psi^* \rangle |0\rangle_U \text{ for all } \psi \in \mathfrak{n} \text{ and } \psi^* \in \mathfrak{n}_+^*.$

Then there exists a unique morhism of vertex superalgebras $\alpha: \Lambda \to U$ such that $\alpha(\psi t^{-1}|0\rangle_{\Lambda}) = u_{\psi}$ and $\alpha(\psi^* t^{-1}|0\rangle_{\Lambda}) = u_{\psi^*}$.

2.6. **Bases.** For each of the objects introduced above – base rings, enveloping algebras, Clifford algebras, and Fock spaces – it is not hard to construct explicit bases (or topological bases). We give the details in the case of two singularities. The construction of a basis depends on the choice of a basis of $\mathbb{C}[t,s][1/ts]$ as an A-module. Following [4], Section 1.1 and Equation (4.1) we introduce the following bases, indexed by $\frac{1}{2}\mathbb{Z}$: for $n \in \mathbb{Z}$ we define

$$\begin{cases} z_n = t^n s^n \\ z_{n+\frac{1}{2}} = t^{n+1} s^n \end{cases} \qquad \begin{cases} w_n = t^n s^n \\ w_{n+\frac{1}{2}} = t^n s^{n+1} \end{cases}$$

The elements z_m for $m \in \frac{1}{2}\mathbb{Z}$ form a basis of $\mathbb{C}[t, s][1/ts]$ as an A-module, and the elements w_n are the dual basis with respect to the residue bilinear form: more precisely, one has

$$\operatorname{Res}_2(z_n w_{-m-\frac{1}{2}}) = \delta_{n,m}.$$

This specific choice of basis is not particularly important, and several others would be possible. However, some properties need to be satisfied for our approach to work. In particular with our choice, the elements z_m (or w_m) with $m \ge 0$ form an A-basis of $\mathbb{C}[t,s]$.

Since K_2 is an A-free module, we deduce that the enveloping algebras of \mathfrak{g}_2 and $\mathcal{C}\ell_2$ are A-free modules. Moreover, as R_2 is a direct summand of K_2 , we also deduce that $\mathbb{V}_2^{\lambda,\mu}$ and Λ_2 are also A-free modules. Explicit bases of these modules, as well as an explicit topological basis of the algebra $\hat{U}_2\hat{\otimes}_A\mathcal{C}\ell_2$, can be obtained using the Poincaré-Birkhoff-Witt theorem and its analogue for Clifford algebras.

3. Vertex algebras and semi-infinite cohomology

In this section, we recall some results obtained by Casarin [1] which allow us to use the formalism of vertex algebras also in the context of several singularities. In particular, using this formalism we develop a notion of semi-infinite cohomology for \hat{U}_2 -modules.

3.1. Distributions and vertex algebra morphisms. Let \mathcal{R} be a complete topological associative A-algebra. Following [1, Definition 3.0.4], we denote by $\mathcal{F}_A(K_2, \mathcal{R})$ the space of continuous A-linear morphisms from K_2 to \mathcal{R} and call it the space of 2-fields. We refer to [1] for the definitions of mutually local 2-fields (Definition 3.1.1), of the n-products $X_{(n)}Y$ of two 2-fields (Definitions 3.1.2 and 3.1.7) and of the derivative $\partial(X)$ of a 2-field (before Definition 3.0.3). The definition in [1] applies also to the other rings we are considering: $K_1, K_t, K_s, K_{t,s}$.

In particular to define n products it is necessary to choose what in [1], Definition 2.3.8, is called a global coordinate. We choose always t as a global coordinate. More explicitly for the rings K_2, K_1, K_t and K_s we choose t = s + a as a global coordinate, and for the ring $K_{t,s} = K_t \times K_s$ we choose (t,t) = (t,s+a).

We also use some foundational results proved in this context in [1]. In particular, the following result will be crucial for us.

Theorem 3.1 ([1], Theorem 3.2.3). Let \mathcal{F} be a \mathbb{C} -linear subspace of $\mathcal{F}_A(K_2, \mathcal{R})$ of mutually local 2-fields closed under derivation and n-products. Let $\mathbb{1}$ be a field such that $\mathbb{1}(f)$ is central for every $f \in K_2$, that $\partial \mathbb{1} = 0$ and such that $\mathbb{1}_{(n)}X = \delta_{n,-1}X$ for all $X \in \mathcal{F}$. Then the vector space $\mathcal{F} + \mathbb{C}\mathbb{1}$, endowed with n-products and derivation $T = \partial$, is a \mathbb{C} -vertex algebra with $\mathbb{1}$ as vacuum vector.

It is straightforward to generalize the constructions and results in [1] to the case of superalgebras \mathcal{R} .

We are interested in the case where \mathcal{R} is the superalgebra $\hat{U}_2 \hat{\otimes}_A \mathcal{C} \ell_2$. For $x \in \mathfrak{g}$, $\psi \in \mathfrak{n}_+$ and $\psi^* \in \mathfrak{n}_+^*$ we define the 2-fields

$$x^{(2)}(g) = (x \otimes g) \otimes 1_{\mathcal{C}\ell_2}, \quad \psi^{[2]}(g) = 1_{\hat{U}_2} \otimes (\psi \otimes g), \quad (\psi^*)^{[2]}(g) = 1_{\hat{U}_2} \otimes (\psi^* \otimes g) \quad (3.1)$$

for all $g \in K_2$. The first of these fields has even parity with respect to the superalgebra structure, while the second and third ones are odd. These fields are mutually local. We consider the minimal \mathbb{C} -linear subspace $\mathcal{F}^{(2)}$ of $\hat{U}_2 \hat{\otimes}_A \mathcal{C} \ell_2$ closed under n-products and derivation and containing the fields (3.1). Moreover, we define

$$\mathbb{1}_2(f) = \operatorname{Res}_2(f) \big(1_{\hat{U}_2} \otimes 1_{\mathcal{C}\ell_2} \big).$$

It is easy to check that this data satisfies the hypothesis of Theorem 3.1. Therefore, $\mathcal{V}^{(2)} = \mathcal{F}^{(2)} + \mathbb{C}\mathbb{1}_2$ has a structure of vertex superalgebra, and by the universal properties of the vertex algebra \mathbb{V} (Lemma 2.2) and of the vertex superalgebra Λ^{\bullet} (Lemma 2.4) it follows that there exists a morphism of vertex superalgebras

$$\Phi^{(2)}: \mathbb{V} \otimes_{\mathbb{C}} \Lambda^{\bullet} \longrightarrow \mathcal{V}^{(2)}. \tag{3.2}$$

This homomorphism will allow us to easily introduce many elements in $\mathcal{V}^{(2)}$, hence also in $\hat{U}_2 \hat{\otimes}_A \mathcal{C} \ell_2$.

Similar constructions apply if the algebra $\hat{U}_2 \hat{\otimes}_A \mathcal{C} \ell_2$ is replaced by the algebras $\hat{U}_1 \hat{\otimes} \mathcal{C} \ell_1$, $\hat{U}_t \hat{\otimes}_Q \mathcal{C} \ell_t$, etc. Hence, we construct the fields $x^{(1)}$, $\psi^{[1]}$, $x^{(t)}$, $\psi^{[t]}$, the vertex superalgebras $\mathcal{V}^{(1)}$, $\mathcal{V}^{(t)}$, and homomorphisms of vertex algebras $\Phi^{(1)}: \mathbb{V} \otimes_{\mathbb{C}} \Lambda^{\bullet} \longrightarrow \mathcal{V}^{(1)}$, $\Phi^{(t)}: \mathbb{V} \otimes_{\mathbb{C}} \Lambda^{\bullet} \longrightarrow \mathcal{V}^{(t)}$, etc.

Notice that we have a specialisation morphism $\operatorname{Sp}_{\mathcal{F}}: \mathcal{F}_A(K_2, \hat{U}_2 \hat{\otimes} \mathcal{C}\ell_2) \longrightarrow \mathcal{F}_{\mathbb{C}}(K_1, \hat{U}_2 \hat{\otimes} \mathcal{C}\ell_1)$ and an expansion map $E_{\mathcal{F}}: \mathcal{F}_A(K_2, \hat{U}_2 \hat{\otimes}_A \mathcal{C}\ell_2) \longrightarrow \mathcal{F}_Q(K_{t,s}, \hat{U}_2 \hat{\otimes}_Q \mathcal{C}\ell_{t,s})$, determined by the conditions

$$(\operatorname{Sp}_{\mathcal{F}}(X))(\operatorname{Sp}(f)) = \operatorname{Sp}(X(f))$$
 and $(E_{\mathcal{F}}(X))(E(f)) = E(X(f)).$

These maps commute with *n*-products and derivations and satisfy $\operatorname{Sp}_{\mathcal{F}}(\mathbb{1}_2) = \mathbb{1}_1$ and $E_{\mathcal{F}}(\mathbb{1}_2) = \mathbb{1}_{t,s}$. Moreover, by construction they satisfy

$$\operatorname{Sp}_{\mathcal{F}}(x^{(2)}) = x^{(1)}$$
 and $E_{\mathcal{F}}(x^{(2)}) = x^{(t,s)}$

for $x \in \mathfrak{g}$. Similar relations hold for $\psi^{[2]}$ and $(\psi^*)^{[2]}$. This implies in particular that the homomorphisms $\operatorname{Sp}_{\mathcal{F}}$ and $E_{\mathcal{F}}$ restrict to homomorphisms of vertex algebras $\operatorname{Sp}: \mathcal{V}^{(2)} \longrightarrow \mathcal{V}^{(1)}$ and $E: \mathcal{V}^{(2)} \longrightarrow \mathcal{V}^{(t,s)}$ such that

$$\operatorname{Sp} \circ \Phi^{(2)} = \Phi^{(1)}$$
 $E \circ \Phi^{(2)} = \Phi^{(t,s)}$.

We can also describe the morphism $\Phi^{(2)}$ through the morphisms $\Phi^{(t)}$ and $\Phi^{(s)}$. Recall from the end of Section 2.4 the maps I, J from $\hat{U}_t \hat{\otimes}_Q \mathcal{C}\ell_t$ and $\hat{U}_s \hat{\otimes}_Q \mathcal{C}\ell_s$ to $\hat{U}_{t,s} \hat{\otimes}_Q \mathcal{C}\ell_{t,s}$. These maps induce maps at the level of fields $I_{\mathcal{F}} : \mathcal{F}_Q(K_t, \hat{U}_t \hat{\otimes}_Q \mathcal{C}\ell_t) \to 0$ $\mathcal{F}_Q(K_{t,s}, \hat{U}_{t,s} \hat{\otimes}_Q \mathcal{C}\ell_{t,s})$ and $J_{\mathcal{F}}: \mathcal{F}_Q(K_s, \hat{U}_s \hat{\otimes}_Q \mathcal{C}\ell_s) \to \mathcal{F}_Q(K_{t,s}, \hat{U}_{t,s} \hat{\otimes}_Q \mathcal{C}\ell_{t,s})$, given by

$$I_{\mathcal{F}}(X)(f,g) = I(X(f))$$
 and $J_{\mathcal{F}}(X)(f,g) = J(X(g))$

for all $(f,g) \in K_t \times K_s = K_{t,s}$. The maps $I_{\mathcal{F}}$ and $J_{\mathcal{F}}$ preserve *n*-products, commute with derivations, and satisfy $I_{\mathcal{F}}(\mathbbm{1}_t) + J_{\mathcal{F}}(\mathbbm{1}_s) = \mathbbm{1}_{t,s}$. Moreover we notice that I(u) and J(v) commute for all $u \in \hat{U}_t \hat{\otimes}_Q \mathcal{C}\ell_t$ and $v \in \hat{U}_s \hat{\otimes}_Q \mathcal{C}\ell_s$. By the discussion in [1, Section 7.2], this implies

$$I_{\mathcal{F}} \circ \Phi^{(t)} + J_{\mathcal{F}} \circ \Phi^{(s)} = \Phi^{(t,s)}.$$

This is the only statement where it is relevant the choice of the global coordinate we have done in Section 3.1.

3.2. Semi-infinite cohomology. We now define a notion of semi-infinite cohomology for \hat{U}_2 -modules, in analogy with the analogous notion for \hat{U}_1 -modules described for example in [5, Chapter 15]. To this end, we introduce some notation for elements in the vertex superalgebra $\mathbb{V} \otimes \Lambda^{\bullet}$. As in the case of \hat{U}_1 , to describe these elements we choose a basis J^a of \mathfrak{g} compatible with the decomposition $\mathfrak{g} = \mathfrak{n}_- \oplus \mathfrak{t} \oplus \mathfrak{n}_+$, where \mathfrak{n}_+ is the nilpotent radical of \mathfrak{b} and \mathfrak{n}_- is the radical of the opposite nilpotent borel subalgebra. We denote by $e^{b,d}_e$ the structure coefficients of the Lie bracket with respect to this basis. We denote by $\Phi \sqcup \Gamma$ the indexing set of the basis J^a , so that, if $\alpha \in \Phi$, then $J^\alpha = e_\alpha = f_{-\alpha}$ is a root vector of weight α and, if $\alpha \in \Gamma$, then $J^\alpha \in \mathfrak{t}$. We also denote by ψ^*_α for $\alpha \in \Phi^+$ the basis of \mathfrak{n}^*_+ dual to the basis e_α of \mathfrak{n}_+ .

With each element in $\mathfrak{n}_+ \otimes \cdots \otimes \mathfrak{n}_+ \otimes \mathfrak{n}_+^* \otimes \cdots \otimes \mathfrak{n}_+^*$ we associate an element in the vertex superalgebra Λ as follows:

$$N(\psi_1 \otimes \dots \psi_\ell \otimes \psi_1^* \otimes \dots \otimes \psi_m^*) = (\psi_1 t^{-1}) \cdots (\psi_\ell t^{-1}) \cdot (\psi_1^* t^{-1}) \cdots (\psi_m^* t^{-1}) \cdot |0\rangle_{\Lambda}.$$

Similarly, with an element in $\mathfrak{g} \otimes \mathfrak{n}_+^*$ we associate an element in the vertex superalgebra $\mathbb{V} \otimes \Lambda^*$ by setting

$$M(x \otimes \psi^*) = (xt^{-1}) \cdot |0\rangle_{\mathbb{V}} \otimes (\psi^*t^{-1}) \cdot |0\rangle_{\Lambda}.$$

Following [5, Chapter 15] we define

$$q = M(I) - \frac{1}{2} |0\rangle_{\mathbb{V}} \otimes N(B) = \sum_{\alpha \in \Phi^{+}} (e_{\alpha}t^{-1}) \cdot |0\rangle_{\mathbb{V}} \otimes (\psi_{\alpha}^{*}t^{-1}) \cdot |0\rangle_{\Lambda}$$
$$- \frac{1}{2} \sum_{\alpha,\beta \in \Phi^{+}} c_{\alpha+\beta}^{\alpha,\beta} |0\rangle_{\mathbb{V}} \otimes (e_{\alpha+\beta}t^{-1}) \cdot (\psi_{\alpha}^{*}t^{-1}) \cdot (\psi_{\beta}^{*}t^{-1}) \cdot |0\rangle_{\Lambda},$$

where $I \in \mathfrak{g} \otimes \mathfrak{n}_+^*$ represents the inclusion of \mathfrak{n}_+ in \mathfrak{g} and $B \in \mathfrak{n}_+ \otimes \mathfrak{n}_+^* \otimes \mathfrak{n}_+^*$ is the Lie bracket. We now define the boundary operator $d_{\mathrm{std}}^{(2)} \in \hat{U}_2 \hat{\otimes}_A \mathcal{C} \ell_2$ as follows:

$$d_{\text{std}}^{(2)} := (\Phi^{(2)}(q))(1).$$

The boundary operator that we will use to define the semi-infinite cohomology is a deformation of $d_{\mathrm{std}}^{(2)}$. Let $\psi_{\mathrm{pr}}^* = \sum_{\alpha \, \mathrm{simple}} \psi_{\alpha}^* \in \mathfrak{n}_+^*$, and define

$$\chi^{(2)} = 1_{\hat{U}_2} \otimes \psi_{\text{pr}}^* = \Phi^{(2)}(N(\psi_{\text{pr}}^*))(1) \in \hat{U}_2 \otimes_A \mathcal{C}\ell_2.$$

Similar constructions yield $\chi^{(s)}$, $\chi^{(t)}$, $\chi^{(s)}$, and $\chi^{(s,t)}$. Finally set

$$d^{(2)} = d_{\text{std}}^{(2)} + \chi^{(2)}.$$

As we will check in Section 3.3, this is an element that squares to zero, and therefore, it can be used to define the semi-infinite cohomology of a \hat{U}_2 -module.

Similarly we can define $d_{\text{std}}^{(1)}$, $\chi^{(1)}$, $d_{\text{std}}^{(1)}$, $\chi^{(t)}$, $d_{\text{std}}^{(t)}$, $\chi^{(t)}$, and so on, as elements of the corresponding superalgebras. By the discussion at the end of Section 3.1 we have

$$\operatorname{Sp}(d^{(2)}) = d^{(1)}, \qquad E(d^{(2)}) = d^{(t,s)}, \quad \text{and} \quad I(d^{(t)}) + J(d^{(s)}) = d^{(t,s)}.$$

Definition 3.2. Let M be an \hat{U}_2 module. Consider the $\hat{U}_2 \hat{\otimes}_A \mathcal{C} \ell_2$ -graded module $M \otimes_A \Lambda_2^{\bullet}$, where the grading is given by charge on Λ_2^{\bullet} . The element $d^{(2)}$ acts on this module as a boundary operator of degree one. Define $\Psi^n(M)$ as the corresponding cohomology of degree n.

Similar constructions apply to modules over the algebras \hat{U}_1 , \hat{U}_t , \hat{U}_s or $\hat{U}_{t,s}$.

Let Z_2 be the center of the algebra \hat{U}_2 , and similarly introduce the center Z_1 of \hat{U}_1 and the centers Z_t and Z_s of \hat{U}_t and \hat{U}_s . If M is an \hat{U}_2 -module, the action of Z_2 on $M \otimes_A \Lambda_2^{\bullet}$ commutes with the differential $d^{(2)}$ and preserves the charge, hence induces an action of Z_2 on the semi-infinite cohomology groups of M. A similar action is defined in the case of \hat{U}_1 -modules or \hat{U}_t -modules.

Recall that a module M over a topological algebra \mathcal{R} is said to be *smooth* if the action of \mathcal{R} on M is continuous with respect to the discrete topology on M. Notice that, if M is a smooth \hat{U}_2 -module, then, since the map E has dense image, the action of \hat{U}_2 on M extends to a smooth action of $\hat{U}_{t,s}$ on $M[a^{-1}]$. Similarly, if M_t is a smooth \hat{U}_t -module and M_s is a smooth \hat{U}_s -module, then there is an induced action of $\hat{U}_{t,s}$ on $M_t \otimes_Q M_s$. In the next section we will use the following properties of the semi-infinite cohomology.

Lemma 3.3.

- a) Given a short exact sequence of \hat{U}_2 -modules, there is an induced long exact sequence in semi-infinite cohomology.
- b) Let M be an \hat{U}_1 -module. The semi-infinite cohomology of M as an \hat{U}_1 -module is isomorphic to the semi-infinite cohomology of M considered as an \hat{U}_2 -module through the map Sp.
- c) Let M be an $\hat{U}_{t,s}$ -module. The semi-infinite cohomology of M as an $\hat{U}_{t,s}$ -module is isomorphic to the semi-infinite cohomology of M considered as an \hat{U}_2 -module through the map E. In particular, this applies to the case where $M = N[a^{-1}]$ is the localization of a smooth \hat{U}_2 -module N.
- d) Let M_t be a smooth \hat{U}_t -module, M_s be a smooth \hat{U}_s -module, and let $M:=M_t\otimes_Q M_s$, regarded as a $\hat{U}_{t,s}$ -module. The complex computing the semi-infinite cohomology of M is the total complex associated with the double complex given by the tensor product of the complex computing the semi-infinite cohomology of M_t and that of M_s . In particular, being the base ring Q a field, if M_t and M_s have non zero semi-infinite cohomology only in degree zero, then M considered as an $\hat{U}_{t,s}$ -module has semi-infinite cohomology only in degree zero and the cohomology in degree zero is isomorphic to the product of the tensor product of $\Psi^0(M_t)$ and $\Psi^0(M_s)$.

Proof. Part a) follows from the fact that Λ_2 is a free module over A.

Part b) follows from the fact that, since $a \in A$ acts trivially on M, by Lemma 2.3 a) we have

$$M \otimes_A \Lambda_2^{\bullet} \simeq M \otimes_{\mathbb{C}} \frac{\Lambda_2^{\bullet}}{a \Lambda_2^{\bullet}} \simeq M \otimes_{\mathbb{C}} \Lambda_1^{\bullet}$$

and moreover, by construction, $d^{(1)} = \operatorname{Sp}(d^{(2)})$.

Part c) follows from the fact that, since the action of a on M is invertible, by Lemma 2.3 c) we have

$$M \otimes_A \Lambda_2^{\bullet} = M \otimes_A \Lambda_2^{\bullet}[a^{-1}] = M \otimes_A \Lambda_{t,s}^{\bullet}$$

and, moreover, by construction, $d^{(t,s)} = E(d^{(2)})$.

Finally, from Lemma 2.3 c) we have

$$(M_t \otimes_Q \Lambda_t^{\bullet}) \otimes_Q (M_s \otimes_Q \Lambda_s^{\bullet}) \simeq M \otimes_Q \Lambda_{t,s}^{\bullet}.$$

Part d) then follows from the equality $d^{(t,s)} = I(d^{(t)}) + J(d^{(s)})$.

3.3. Commutation relations. For their computation of the semi-infinite cohomology of \mathbb{V} , Frenkel and Ben Zvi (see [5] Chapter 15) relied on the choice of a clever basis of $\mathbb{V} \otimes \Lambda$. For all $x \in \mathfrak{g}$, they define

$$\hat{x} = xt^{-1} \cdot |0\rangle_{\mathbb{V}} \otimes |0\rangle_{\Lambda} + N(\alpha_x),$$

where $\alpha_x \in \mathfrak{n}_+ \otimes \mathfrak{n}_+^*$ represents the linear map $\mathfrak{n}_+ \to \mathfrak{n}_+$ obtained as the composition of $\mathrm{ad}_x : \mathfrak{n}_+ \longrightarrow \mathfrak{n}_+$, the natural projection $\pi : \mathfrak{g} \longrightarrow \mathfrak{g}/\mathfrak{b}_-$, and the inverse of the isomorphism $\mathfrak{n}_+ \cong \mathfrak{g}/\mathfrak{b}_-$ induced by π . Using the map $\Phi^{(2)}$ from Equation (3.2) we define

$$\hat{x}^{(2)} = \Phi^{(2)}(\hat{x}).$$

To compute the semi-infinite cohomology of $\mathbb{V}_2^{\lambda,\mu}$ we will need some information about the commutation relations among the elements $\hat{x}^{(2)}$, $\psi^{[2]}$, and $(\psi^*)^{[2]}$, and the boundary operators. These are easy to compute because all these objects are constructed through the map $\Phi^{(2)}$. Let us make this remark precise. Given an element x in $\mathbb{V}\otimes\Lambda$, denote by x(z) the corresponding field in the vertex superalgebra and by $x^{(2)}: K_2 \longrightarrow \hat{U}_2 \hat{\otimes} \mathcal{C}\ell_2$ the 2-field $\Phi^{(2)}(x)$. For any choice of elements $x, y \in \mathbb{V} \otimes \Lambda$, the commutator of the corresponding fields is given by

$$[x(z), y(w)] = \sum_{n \ge 0} \frac{1}{n!} (x_{(n)}y)(w) \partial_w^n \delta(z - w).$$

We have a similar Operator Product Expansion formula for 2-fields (see [1], Proposition 3.1.3)

$$[x^{(2)}(f), y^{(2)}(g)] = \sum_{n>0} \frac{1}{n!} \left((x^{(2)})_{(n)} (y^{(2)}) \right) (g \, \partial^n f),$$

where the product $(x^{(2)})_{(n)}(y^{(2))}$ is the product of 2-fields defined in [1]. However, since $\Phi^{(2)}$ is a map of vertex algebras we get $(x^{(2)})_{(n)}(y^{(2)}) = (x_{(n)}y)^{(2)}$. Hence, if we know the commutator of x(z), y(w), we immediately deduce that of $x^{(2)}$ and $y^{(2)}$.

Similar considerations apply when we want to compute $[x^{(2)}(1), y^{(2)}(g)]$ assuming we know the commutator of $x_{(0)}$ and y(w). In this case, the usual OPE formula gives $[x_{(0)}, y(w)] = (x_{(0)}y)(w)$, while the OPE formula for 2-fields gives

$$[x^{(2)}(1), y^{(2)}(g)] = ((x^{(2)})_{(0)}(y^{(2)}))(g).$$

Using again the fact that $\Phi^{(2)}$ is a map of vertex algebras, we get

$$[x^{(2)}(1), y^{(2)}] = \Phi^{(2)} \left(\left([x_{(0)}, y(w)](|0\rangle_{\mathbb{V}} \otimes |0\rangle_{\Lambda}) \right)|_{w=0} \right).$$

These formulas are enough to determine all commutation relations among the elements $\hat{x}^{(2)}$, $\psi^{[2]}$, $(\psi^*)^{[2]}$ and the boundary operators from those obtained by Frenkel and Ben Zvi in [5, Chapter 15], without the need of any further computation. We summarise these results in Proposition 3.4 below, which (in light of the above) follows from Sections 15.2.4 and 15.2.9 of [5]. In the statement, we denote by $e_{\rm pr}, h_{\rm pr}, f_{\rm pr}$ the $\mathfrak{sl}(2)$ -triple such that $f_{\rm pr} = \sum_{\alpha \text{ simple}} \lambda_{\alpha} f_{\alpha}$, $\kappa(f_{pr}, e_{\alpha}) = 1$ for all simple root α and $h_{pr} \in \mathfrak{t}$.

Proposition 3.4. for all $x \in \mathfrak{g}$, $y \in \mathfrak{b}$, $z \in \mathfrak{n}_+$, $w \in \mathfrak{b}_-$, $\psi \in \mathfrak{n}_+$ and $\psi^* \in \mathfrak{n}_+^*$ we have:

$$\begin{array}{ll} a) & (d_{\rm std}^{(2)})^2 = 0, & [d_{\rm std}^{(2)}, \chi^{(2)}]_+ = 0, \\ b) & (\chi^{(2)})^2 = 0, & (d^{(2)})^2 = 0, \\ c) & [\chi^{(2)}, \psi^{[2]}]_+ = \langle \psi_{\rm pr}^*, \psi \rangle \, \mathbbm{1}, & [\chi^{(2)}, (\psi^*)^{[2]}]_+ = 0, \\ d) & [\chi^{(2)}, \hat{z}^{(2)}] = 0, & [\chi^{(2)}, \hat{w}^{(2)}] = \sum_{\alpha \in \Phi^+} \kappa([f_{\rm pr}, z], e_\alpha) \psi_\alpha^*, \\ e) & [d_{\rm std}^{(2)}, \psi^{[2]}]_+ = \hat{\psi}^{(2)}, & [d_{\rm std}^{(2)}, (\psi^*)^{[2]}]_+ = -\frac{1}{2} \Phi \big(1_{\hat{U}_2} \otimes N(\psi^* \circ B) \big), \\ f) & [d_{\rm std}^{(2)}, \hat{y}^{(2)}]_+ = 0 & \end{array}$$

where in the second formula of e) the element $\psi^* \circ B \in \mathfrak{n}_+^* \otimes \mathfrak{n}_+^*$ represents the composition of the bracket with the map ψ^* . Moreover, if we choose a basis J^a as at the beginnin of Section 3.2, for all $\gamma \in \Phi^+$ we have

$$\begin{split} [d_{\mathrm{std}}^{(2)}, \hat{f}_{\gamma}^{(2)}]_{+} &= \sum_{\alpha \in \Phi^{+}, a \in \Phi^{-} \sqcup \Gamma} c_{a}^{\alpha, -\gamma} (\hat{J}^{a})^{(2)}{}_{(-1)} (\psi_{\alpha}^{*})^{[2]} \\ &- \frac{1}{2} \, \kappa(e_{-\gamma}, f_{\gamma}) \, \partial(\psi_{-\gamma}^{*})^{[2]} - \sum_{\alpha, \beta \in \Phi^{+}, \, a \in \Phi \sqcup \Gamma} c_{\beta}^{\alpha, a} c_{a}^{\beta, -\gamma} \partial(\psi_{\alpha}^{*})^{[2]} \end{split}$$

By specialisation and localization we obtain that similar formulas hold also in the case of our various other superalgebras $\hat{U}_t \hat{\otimes} \mathcal{C}\ell_t$, $\hat{U}_{t,s} \otimes \mathcal{C}\ell_{t,s}$, ...

4. The semi-infinite cohomology of $\mathbb{V}_2^{\lambda,\mu}$

In this section we compute the semi-infinite cohomology of $\mathbb{V}_2^{\lambda,\mu}$. We denote by $C_2^{\bullet} = C_2^{\bullet}(\lambda,\mu)$ the complex $\mathbb{V}_2^{\lambda,\mu} \otimes_A \Lambda_2^{\bullet}$ and similarly we introduce the complexes $C_t^{\bullet} = C_t^{\bullet}(\lambda) = \mathbb{V}_t^{\lambda} \otimes_Q \Lambda_t^{\bullet}$ and $C_s^{\bullet} = C_s^{\bullet}(\mu) = \mathbb{V}_s^{\mu} \otimes_Q \Lambda_s^{\bullet}$. We further introduce the complexes $C_1^{\bullet}(\nu) = \mathbb{V}_1^{\nu} \otimes_C \Lambda_1^{\bullet}$ and $C_1^{\bullet}(\lambda,\mu) = \mathbb{W}_1^{\lambda,\mu} \otimes_C \Lambda_1^{\bullet}$. Hence, we have $C_1^{\bullet}(\lambda,\mu) \simeq \oplus C_1^{\bullet}(\nu)$, where the sum ranges over the irreducible factors of $V^{\lambda} \otimes V^{\mu}$ counted with multiplicity.

We denote by Op_1 the indscheme of opers on the punctured disc and, for every integral dominant weight ν , we write $\operatorname{Op}_1^{\nu}$ for the associated connected component of the space of unramified opers without monodromy, equipped with its reduced structure (see, for example, [7] for a more complete definition). We also denote by v_{ν} a highest weight vector in the \mathfrak{g} -module V^{ν} . Feigin and Frenkel [2] constructed an isomorphism \mathcal{F}_1 : Funct(Op_1) $\longrightarrow Z_1$ between the space of functions over Op_1 and the center Z_1 of \hat{U}_1 . Recall the following result, which combines Theorem 1, Theorem 2 and the proof of Proposition 1 in [7].

Theorem 4.1 (Frenkel and Gaitsgory [7]). The action of Z_1 on \mathbb{V}_1^{ν} and the Feigin-Frenkel isomorphism induce an isomorphism

$$\mathcal{G}_1: \operatorname{Funct}(\operatorname{Op}_1^{\nu}) \longrightarrow \operatorname{End}_{\hat{\mathfrak{a}}_1}(\mathbb{V}_1^{\nu}).$$

Moreover, the element $v_{\nu} \otimes |0\rangle_{\Lambda}$ is a cocycle in $C_1^{\bullet}(\nu)$ and the map $z \mapsto [z \cdot v_{\nu} \otimes |0\rangle_{\Lambda}]$ from Z_1 to $\Psi^0(\mathbb{V}^{\nu}_1)$ induces isomorphisms of Z_1 -modules

$$\operatorname{Funct}(\operatorname{Op}_1^{\nu}) \simeq \operatorname{End}_{\hat{\mathfrak{a}}_1}(\mathbb{V}_1^{\nu}) \simeq \Psi^0(\mathbb{V}_1^{\nu}).$$

Finally, $\Psi^n(\mathbb{V}_1^{\nu})$ vanishes for all $n \neq 0$.

The result of Frenkel and Gaitsgory generalises easily to the case of the modules \mathbb{V}^{λ}_{t} and \mathbb{V}^{μ}_{s} .

By Lemma 2.1 and Lemma 2.3, as in the proof of Lemma 3.3, by the compatibility of boundary operators we get homomorphisms of complexes Sp : $C_2^{\bullet} \to C_1^{\bullet}(\lambda, \mu)$ and $E: C_2^{\bullet} \to C_t^{\bullet}(\lambda) \otimes_Q C_s^{\bullet}(\mu)$. These induce isomorphisms

$$C_2^{\bullet}[a^{-1}] \simeq C_t^{\bullet}(\lambda) \otimes_Q C_s^{\bullet}(\mu)$$
 and $\frac{C_2^{\bullet}}{aC_2^{\bullet}} \simeq C_1^{\bullet}(\lambda, \mu).$ (4.1)

From these isomorphisms and Theorem 4.1 it follows easily that $\Psi^n(\mathbb{V}_2^{\lambda,\mu})$ is zero for $n \neq 0, 1$, and we could also get information on the cohomology in degrees zero and one.

However, it is easier to compute these cohomology groups directly by adapting the strategy employed by Frenkel and Ben Zvi in [5, Chapter 15]. In order to do this, we now introduce certain subcomplexes of C_2^{\bullet} . We denote by $1_{\mathbb{V}_2^{0,0}}$ the element $1 \in A \otimes_{\mathbb{C}} \mathbb{C} \otimes_{\mathbb{C}} \mathbb{C} \subset \mathbb{V}_2^{0,0}$.

Definition 4.2. We denote by E_2^{\bullet} the subcomplex of $C_2^{\bullet}(0,0)$ spanned by elements of the form

$$\hat{x}_1^{(2)}(g_1)\cdots\hat{x}_a^{(2)}(g_a)\cdot 1_{\mathbb{V}_0^{0,0}}\otimes\psi_1^{(2)}(\ell_1)\cdots\psi_b^{(2)}(\ell_b)\cdot |0\rangle_{\Lambda_2}$$
(4.2)

where $x_i, \psi_i \in \mathfrak{n}_+$ and $g_1, \dots, g_a, \ell_1, \dots \ell_b \in K_2$. By the commutation relations of Section 3.3 we see that E_2^{\bullet} is a subcomplex of $C_2^{\bullet}(0,0)$.

We define also analogous complexes E_t^{\bullet} , E_s^{\bullet} and E_1^{\bullet} . These complexes were denoted by C' in [5] and by C_0 in [7]. By construction, these subcomplexes are compatible with specialisation and localization, and there are isomorphisms $E_2^{\bullet}/aE_2^{\bullet} \simeq E_1^{\bullet}$ and $E_2^{\bullet}[a^{-1}] \simeq E_t^{\bullet} \otimes_Q E_s^{\bullet}$.

Definition 4.3. We denote by $D_2^{\bullet} = D_2^{\bullet}(\lambda, \mu)$ the subcomplex of $C_2^{\bullet}(\lambda, \mu)$ spanned by elements of the form

$$\hat{y}_1^{(2)}(h_1)\cdots\hat{y}_c^{(2)}(h_c)\cdot w\otimes(\psi_1^*)^{(2)}(k_1)\cdots(\psi_d^*)^{(2)}(k_d)\cdot|0\rangle_{\Lambda_2}$$
(4.3)

where $w \in V^{\lambda} \otimes V^{\mu}$, $y_i \in \mathfrak{b}_- = \mathfrak{n}_- + \mathfrak{t}$, $\psi_i^* \in \mathfrak{n}_+^*$ and $h_1, \ldots, h_c, k_1, \ldots, k_d \in K_2$. By the commutation relations of Section 3.3 we see that E_2^{\bullet} is a subcomplex of $C_2^{\bullet}(\lambda, \mu)$.

We define also analogous complexes $D_t^{\bullet}(\lambda)$, $D_s^{\bullet}(\mu)$ and $D_1^{\bullet}(\nu)$. These complexes were denoted by C_0 in [5] and by C' in [7]. Finally, we denote by $D_1^{\bullet}(\lambda,\mu)$ the analogous subcomplex of $C_1^{\bullet}(\lambda,\mu)$. By construction, these subcomplexes are compatible with specialisation and localization, and there are isomorphisms $D_2^{\bullet}/aD_2^{\bullet} \simeq D_1^{\bullet}(\lambda,\mu)$ and $D_2^{\bullet}[a^{-1}] \simeq D_t^{\bullet}(\lambda) \otimes_Q D_s^{\bullet}(\mu)$.

There is an isomorphism of complexes $E_2^{\bullet} \otimes D_2^{\bullet} \longrightarrow C_2^{\bullet}$ defined by

$$\left(\underline{x}\cdot 1_{\mathbb{V}_2^{0,0}}\otimes \underline{\psi}\cdot |0\rangle_{\Lambda_2}\right)\otimes \left(\underline{y}\cdot w\otimes \underline{\psi}^*\cdot |0\rangle_{\Lambda_2}\right)\longmapsto \underline{x}\cdot \underline{y}\cdot w\otimes \underline{\psi}\cdot \underline{\psi}^*\cdot |0\rangle_{\Lambda_2},$$

where $\underline{x} = \hat{x}_1^{(2)}(g_1) \cdots \hat{x}_a^{(2)}(g_a)$ and $\underline{\psi} = \psi_1^{(2)}(\ell_1) \cdots \psi_b^{(2)}(\ell_b)$ are as in Equation (4.2), $\underline{y} = \hat{y}_1^{(2)}(h_1) \cdots \hat{y}_c^{(2)}(h_c)$ and $\underline{\psi}^* = (\psi^*)^{(2)}(k_1) \cdots (\psi^*)^{(2)}(k_d)$ are as in Equation (4.3), and w is an element of $V^{\lambda} \otimes V^{\mu}$.

We now compute the cohomology of the complex E_2^{\bullet} . We will need the following result by Frenkel and Ben Zvi.

Lemma 4.4 ([5, Section 15.2.6]). $H^n(E_1^{\bullet}) = 0$ for $n \neq 0$ and $\Psi^0(E_1^{\bullet}) = \mathbb{C}[|0\rangle_{\mathbb{V}} \otimes |0\rangle_{\Lambda}].$

This result generalizes easily to the case of E_t^{\bullet} and E_s^{\bullet} . Localizing and specializing, we deduce the following lemma.

Lemma 4.5.
$$H^n(E_2^{\bullet}) = 0$$
 for $n \neq 0$ and $H^0(E_2^{\bullet}) = A[1_{\mathbb{V}_2^{0,0}} \otimes |0\rangle_{\Lambda_2}].$

Proof. By definition, the complex E_2^{\bullet} is concentrated in non-positive degrees. Hence, the long exact sequence induced by

$$0 \longrightarrow E_2^{\bullet} \xrightarrow{a \cdot} E_2^{\bullet} \longrightarrow E_1^{\bullet} \longrightarrow 0$$

implies that $H^n(E_2^{\bullet})$ is torsion free for every n, and that the specialisation of $H^0(E_2^{\bullet})$ is isomorphic to $H^0(E_1^{\bullet})$. Since semi-infinite cohomology commutes with localization (Lemma 3.3), using Lemma 2.1 and Lemma 4.4 we get the desired result.

We now compute the cohomology of D_2^{\bullet} . The strategy is similar, but the argument is less straightforward since we do not have an explicit representative for $H^0(D_1^{\bullet})$. Following the strategy in [5], we introduce the following bigraded structure on D_2^{\bullet} . Recall that the *height* $\operatorname{ht}(\alpha)$ of a root α is equal to the sum of the coefficients of α when written as a sum of simple roots. Let also $e_{\operatorname{pr}}, h_{\operatorname{pr}}, f_{\operatorname{pr}}$ be an $\mathfrak{sl}(2)$ -triple such that $f_{\operatorname{pr}} = \sum_{\alpha \text{ simple}} f_{\alpha}$ and h_{pr} belongs to \mathfrak{t} .

Definition 4.6. We define a bidegree, with values in $\frac{1}{2}\mathbb{Z} \times \frac{1}{2}\mathbb{Z}$ and denoted by bideg, as follows. On elements of $\hat{\mathfrak{g}}_2$, we set

$$bideg(x \otimes g) = (-n, n)$$

if $x \in \mathfrak{g}$ is such that $[h_{\mathrm{pr}}, x] = 2 n x$ and $g \in K_2$. We set also the bidegree of the central element $C_2 \in \hat{\mathfrak{g}}_2$ to be (0,0). This induces a bidegree on $U(\hat{\mathfrak{g}}_2)$. On the space $X_2 = K_2 \otimes \mathfrak{n}_+ \oplus K_2 \otimes \mathfrak{n}_+^*$ (see Section 2.4) we define

$$\operatorname{bideg} e_{\alpha} \otimes g = (-\operatorname{ht}(\alpha), -1 + \operatorname{ht}(\alpha))$$

bideg
$$\psi_{\alpha}^* \otimes g = (\operatorname{ht}(\alpha), 1 - \operatorname{ht}(\alpha))$$

for α a positive root and g any element of K_2 . This induces a bidegree on the Clifford algebra $\mathcal{C}\ell_2$. Moreover, if W is any finite-dimensional representation of \mathfrak{g} , then we set

$$bideg w = (-n, n)$$

if $w \in W$ is such that $h_{\text{pr}} \cdot w = 2 n w$. These choices induces a bidegree on the module $C_2^{\bullet}(\lambda, \mu)$, and the element $\hat{x}^{(2)}(g)$ is homogeneous of bidegree (-n, n) if $[h_{\text{pr}}, x] = 2 n x$. Finally, notice that if an element has bidegree (p, q), then it has charge p+q. In particular, we introduce the submodule $D_2^{p,q}$ of elements of D_2^{p+q} of bidegree (p, q).

We notice also that bideg $d_{\mathrm{std}}^{(2)}=(0,1)$ and that bideg $\chi^{(2)}=(1,0)$. In particular, $D_2^{\bullet,\bullet}$ is a double complex and D_2^{\bullet} is the associated total complex. Following Frenkel and Ben Zvi [5, Chapter 15], the cohomology of the rows of this double complex is easy to describe. Let $\mathfrak a$ be the centralizer of f_{pr} in $\mathfrak g$. Recall from [5, Lemma 15.1.3 and Section 15.2.9] that the space spanned by monomials of the form $(\hat p_1)_{n_1}\cdots(\hat p_k)_{n_k}\cdot|0\rangle_{\mathbb V}\otimes|0\rangle_{\Lambda}$ with $p_i\in\mathfrak a$ generates a commutative vertex subalgebra F_1 of $\mathbb V\otimes\Lambda^{\bullet}$ isomorphic to $S^{\bullet}(\mathfrak a\otimes t^{-1}\mathbb C[t^{-1}])$. As in Section 3.3, it follows that for $x,y\in\mathfrak a$ the fields $\hat x^{(2)}$ and $\hat y^{(2)}$ commute.

We define $F_2(\lambda, \mu)$ as the span of elements of the form

$$\hat{x}_1^{(2)}(g_1)\cdots\hat{x}_k^{(2)}(g_k)\cdot(v\otimes|0\rangle_{\Lambda_2})\in\mathbb{V}_2^{\lambda,\mu}\otimes_A\Lambda_2^{\bullet}$$

with $x_1, \ldots, x_k \in \mathfrak{a}$ and $v \in V^{\lambda} \otimes V^{\mu}$. Notice that all these elements have charge equal to zero, and that the space $F_2(\lambda, \mu)$ splits as a direct sum $F_2(\lambda, \mu) = \bigoplus_q F_2^{-q,q}(\lambda, \mu)$ according to the bidegree introduced above. Moreover, by Proposition 3.4 d), these elements are annihilated by the action of $\chi^{(2)}$.

Similarly we construct subspaces $F_1^{-q,q}(\nu) \subset \mathbb{V}_1^{\nu} \otimes_{\mathbb{C}} \Lambda_1^{\bullet}, \ F_t^{-q,q}(\lambda) \subset \mathbb{V}_t^{\lambda} \otimes_{Q} \Lambda_t^{\bullet}, F_s^{-q,q}(\lambda) \subset \mathbb{V}_t^{\lambda} \otimes_{Q} \Lambda_t^{\bullet},$ and $F_1^{-q,q}(\lambda,\mu) \subset \mathbb{W}_1^{\lambda,\mu} \otimes_{\mathbb{C}} \Lambda_1^{\bullet},$ In particular, $F_1^{-q,q}(\lambda,\mu) = \mathbb{E}[X_t^{-q,q}(\lambda,\mu)]$

 $\bigoplus_{\nu} F_1^{-q,q}(\nu)$ where the sum is over all irreducible factors of $V^{\lambda} \otimes_{\mathbb{C}} V^{\mu}$. By construction, the specialisation and localization maps induce isomorphisms

$$\frac{F_2^{-q,q}(\lambda,\mu)}{aF_2^{-q,q}(\lambda,\mu)} \simeq F_1^{-q,q}(\lambda,\mu) \quad \text{and} \quad F_2^{-q,q}(\lambda,\mu)[a^{-1}] \simeq \bigoplus_{b+c=q} F_t^{-b,b}(\lambda) \otimes_Q F_s^{-c,c}(\mu).$$

Recall the following result on the cohomology of $D_1^{\bullet,q}$ with respect to the boundary $\chi^{(1)}$.

Lemma 4.7 ([5, Lemma 15.2.10] and [7]). Let $2p_{\nu} = \langle \nu, h_{\text{pr}} \rangle$.

- a) $D_1^{p,q}(\nu) = 0$ for $q > p_{\nu}$ and for p < -q. In particular, $D_1^{p,q} = 0$ for $q > p_{\lambda+\mu}$ and for p < -q; b) $H^n(D_1^{\bullet,q}(\nu)) = 0$ for $n \neq -q$. In particular, $H^n(D_1^{\bullet,q}(\lambda,\mu)) = 0$ for $n \neq -q$; c) The map $v \mapsto [v]$ from $F_1^{-q,q}(\nu)$ to $H^{-q}(D_1^{\bullet,q}(\nu))$ is an isomorphism.

Finally, it follows from c) that the map $v \mapsto [v]$ from $F_1^{-q,q}(\lambda,\mu)$ to $H^{-q}(D_1^{\bullet,q}(\lambda,\mu))$ is also an isomorphism.

Similar results hold for the complexes $D_t^{\bullet,q}(\lambda)$ and $D_s^{\bullet,q}(\mu)$. From this result we deduce the cohomology of the complex $D_q^{\bullet,q}$ with respect to the boundary operator

Lemma 4.8. Let $2p_0 = \langle \lambda + \mu, h_{pr} \rangle$ as above.

- a) $D_2^{p,q}=0$ for $q>p_0$ and for p<-q; b) $H^n(D_2^{\bullet,q})=0$ for $n\neq -q$; c) The map $v\mapsto [v]$ from $F_2^{-q,q}(\lambda,\mu)$ to $H^{-q}(D_2^{\bullet,q}(\lambda,\mu))$ is an isomorphism of

Proof. Part a) is clear for the definition of $D_2^{p,q} = 0$. For parts b) and c), we start by studying the localization of the cohomology groups of $D_2^{\bullet,q}$. Equivalently, we aim to compute the cohomology of the localization of the row $D_2^{\bullet,q}$. This localization can be rewritten as

$$\bigoplus_{b+c=q} D_t^{\bullet,b}(\lambda) \otimes D_s^{\bullet,c}(\mu).$$

In particular, it follows from Lemma 4.7 that its cohomology is concentrated in degree -q, and that its cohomology in this degree is given by

$$\bigoplus_{b+c=q} F_t^{-b,b}(\lambda) \otimes F_s^{-c,c}(\mu),$$

which is the localization of $F_2^{-q,q}(\lambda,\mu)$. Since specialisation is compatible with bideg, we have an isomorphism $D_2^{\bullet,q}/aD_2^{\bullet,q}\simeq D_1^{\bullet,q}(\lambda,\mu)$. Using Lemma 4.7, the associated long exact sequence shows that $H^n(D_2^{\bullet,q})$ is torsion-free for $n\neq -q+1$, and that the map

$$\iota: H^{-q}(D_2^{\bullet,q})/aH^{-q}(D_2^{\bullet,q}) \to H^{-q}(D_1^{\bullet,q}(\lambda,\mu))$$

We now prove c). Notice that both $F_2^{-q,q}(\lambda,\mu)$ and $H^{-q}(D_2^{\bullet,q}(\lambda,\mu))$ are torsionfree. We have already shown that the localization of the natural maps between them is an isomorphism. To study its specialisation, we compose it with the injection ι . This composition is the isomorphism of the last remark of Lemma 4.7. We conclude by applying Lemma 2.1.

In order to prove b), it is enough to notice that from the above discussion we know that, for $n \neq -q$, the module $H^n(D_2^{\bullet,q}) = 0$ is torsion-free, and that its localization is trivial.

Let now be $\varphi_i^{(q)}$ be an A-basis of $F_2^{-q,q}(\lambda,\mu)$. Since the cohomology in degree -q of the complex $D_2^{\bullet,q+1}$ is zero, there exists an element $\varphi_{i,1}^{(q)} \in D_2^{-q-1,q+1}$ such that $\chi^{(2)}(\varphi_{i,1}^{(q)}) = -d_{\mathrm{std}}^{(2)}(\varphi_i^{(q)})$. By induction, we can construct elements $\varphi_{i,0}^{(q)} = \varphi_i^{(q)}$ and $\varphi_{i,\ell}^{(q)} \in D_2^{-q-\ell,q+\ell}$ such that their sum

$$\tilde{\varphi}_i^{(q)} = \sum_{\ell=0}^{p_0 - q} \varphi_{i,\ell}^{(q)}$$

satisfies $d^{(2)}(\tilde{\varphi}_i^{(q)}) = 0$. We now prove the main result of this section.

Theorem 4.9. The following hold.

- a) $\Psi^n(\mathbb{V}_2^{\lambda,\mu}) = 0$ for $n \neq 0$.
- b) We have an isomorphism

$$\frac{\Psi^0(\mathbb{V}_2^{\lambda,\mu})}{a\Psi^0(\mathbb{V}_2^{\lambda,\mu})} \simeq \Psi^0(\mathbb{W}_1^{\lambda,\mu}) \simeq \bigoplus_{\nu} \Psi^0(\mathbb{V}_1^{\nu})$$

where the sum ranges over all irreducible components V^{ν} of $V^{\lambda} \otimes V^{\mu}$, counted with multiplicity.

c) The elements $\left[\tilde{\varphi}_{i}^{(q)}\right]$ are an A-basis of $\Psi^{0}(\mathbb{V}^{\lambda,\mu})$.

Proof. From Lemma 4.8 we deduce that the classes of the elements $\tilde{\varphi}_i^{(q)}$ form an A-basis of $H^0(D_2^{\bullet})$, and that $H^n(D_2^{\bullet}) = 0$ for $n \neq 0$. As the complex D_2^{\bullet} is concentrated in non-negative degrees, by a standard homological argument we deduce that $H^n(\mathbb{V}_2^{\lambda,\mu})$ is isomorphic to the n-th cohomology of the complex $H^0(D_2^{\bullet}) \otimes_A E_2^{\bullet}$. Using Lemma 4.5, we immediately obtain parts a) and c).

The second isomorphism appearing in part b) is clear, while the first follows from a) and the long exact sequence associated with the isomorphism

$$\frac{C_2^{\bullet}}{aC_2^{\bullet}} \simeq C_1^{\bullet}(\lambda, \mu).$$

We will use the following Corollary in the next Section.

Corollary 4.10. The element $[v_{\lambda} \otimes v_{\mu}] \in \Psi^{0}(\mathbb{V}^{\lambda,\mu})$ is indivisible.

Proof. By the previous theorem we can choose $[v_{\lambda} \otimes v_{\mu}]$ as an element of a basis of the free A module $\Psi^{0}(\mathbb{V}^{\lambda,\mu})$.

5. The action of the center

In this section we study the action of the center Z_2 on the semi-infinite cohomology of the module $\mathbb{V}_2^{\lambda,\mu}$.

In this section we show that $\mathbb{V}_2^{\lambda,\mu}$ is not a perfect analogue of the Weyl module \mathbb{V}_1^{ν} . Indeed, we show that, as a Z_2 -module, the semi-infinite cohomology of $\mathbb{V}_2^{\lambda,\mu}$ is not isomorphic to $\operatorname{End}_{\hat{\mathfrak{g}}_2}(\mathbb{V}_2^{\lambda,\mu})$ or to $\operatorname{Funct}(\operatorname{Op}_2^{\lambda,\mu})$.

We begin by observing that the module $\Psi^0(\mathbb{V}_1^{\nu})$ has no non-trivial Z_1 -equivariant automorphisms.

First we notice, that by construction, the action of \mathbb{Z}_2 commutes with localization and specialisation, as introduced before Equation (4.1). Concretely, we have:

$$E_t(z \cdot x) = E_t(z) \cdot E_t(x), \quad E_s(z \cdot x) = E_s(z) \cdot E_s(x), \quad \operatorname{Sp}(z \cdot x) = \operatorname{Sp}(z) \cdot \operatorname{Sp}(x)$$
 for all $z \in \mathbb{Z}_2$ and for all $x \in \Psi^0(\mathbb{V}_2^{\lambda,\mu})$.

Lemma 5.1. If $\mathcal{K}: \operatorname{End}_{\hat{\mathfrak{g}}_t}(\mathbb{V}_t^{\lambda}) \otimes_Q \operatorname{End}_{\hat{\mathfrak{g}}_s}(\mathbb{V}_s^{\mu}) \longrightarrow \Psi^0(\mathbb{V}_t^{\lambda}) \otimes_Q \Psi^0(\mathbb{V}_s^{\mu})$ is a $(Z_t \otimes Z_s)$ -equivariant isomorphism, then $\mathcal{K}(\operatorname{Id}_{\mathbb{V}_t^{\lambda}} \otimes \operatorname{Id}_{\mathbb{V}_s^{\mu}}) = q[v_{\lambda}] \otimes [v_{\mu}]$ for some $q \in Q \setminus \{0\}$.

Proof. It follows from Theorem 4.1 that $\operatorname{End}_{\hat{\mathfrak{g}}_t \times \hat{\mathfrak{g}}_s}(\mathbb{V}^{\lambda_t} \otimes_Q \mathbb{V}^{\mu_s})$ is isomorphic to $\operatorname{Funct}(\operatorname{Op}_t^{\lambda} \times_{\operatorname{Spec} Q} \operatorname{Op}_s^{\mu}) = \operatorname{Funct}(\operatorname{Op}_t^{\lambda}) \otimes_Q \operatorname{Funct}(\operatorname{Op}_s^{\mu})$ and this is a polynomial ring in infinitely many variables over the field Q. In particular, its only invertible elements are the non-zero scalars in Q.

Moreover, Theorem 4.1 also implies that Funct($\operatorname{Op}_t^{\lambda}$) is isomorphic as a Z_t -module to $\Psi^0(\mathbb{V}_t^{\lambda})$, with an isomorphism given by $z \longrightarrow \mathcal{G}_t(z) \cdot [v_{\lambda}]$. The claim follows.

Before proving that $\mathbb{V}_2^{\lambda,\mu}$ does not have the "right" semi-infinite cohomology we recall some properties of the modules \mathbb{V}_1^{ν} that will be needed also in the next section.

Remark 5.2. We denote by Z_1^{ν} the coordinate ring of the scheme $\operatorname{Op}_1^{\nu}$. Recall that the schemes $\operatorname{Op}_1^{\nu}$ for different values of ν are disjoint, so that the map $Z_1 \longrightarrow Z_1^{\nu_1} \times \cdots \times Z_1^{\nu_k}$ is surjective if the weights ν_i are distinct. Recall also that the ring Z_1^{ν} is a polynomial ring in infinitely many variables. This implies that

- (1) There are no nontrivial $\hat{\mathfrak{g}}_1$ -morphisms between the \hat{U}_1 -modules \mathbb{V}_1^{ν} and $\mathbb{V}_1^{\nu'}$ if $\nu \neq \nu'$.
- (2) There are no nontrivial extensions between the \hat{U}_1 -modules \mathbb{V}_1^{ν} and $\mathbb{V}_1^{\nu'}$ if $\nu \neq \nu'$.
- (3) Assume that $\alpha: \prod Z^{\nu_i} \longrightarrow \prod Z^{\nu_i}$ is a map of Z-modules and that the weights ν_i are distinct. If 1 is in the image of α then α is an isomorphism and $\alpha(Z_1^{\nu_i}) = Z_1^{\nu_i}$.

By the Feigin-Frenkel Theorem (see [4] Theorem 5.2) the ring Funct(Op₂) is isomorphic to Z_2 . In the sequel we will identify these rings through this isomorphism. In particular the ring Funct(Op₂^{λ,μ}) is a quotient of Z_2 . We will denote Funct(Op₂^{λ,μ}) by $Z_2^{\lambda,\mu}$.

We now prove that $Z_2^{\lambda,\mu}$ and $\Psi^0(\mathbb{V}_2^{\lambda,\mu})$ are not isomorphic.

Proposition 5.3. Assume that $V^{\lambda} \otimes V^{\mu}$ is not irreducible. Then the two Z_2 -modules $\operatorname{End}_{\mathfrak{g}_2}(\mathbb{V}_2^{\lambda,\mu})$ and $\Psi^0(\mathbb{V}_2^{\lambda,\mu})$ are not isomorphic. Similarly the two Z_2 -modules $Z_2^{\lambda,\mu}$ and $\Psi^0(\mathbb{V}_2^{\lambda,\mu})$ are not isomorphic.

Proof. Suppose $\mathcal{H}: \operatorname{End}_{\hat{\mathfrak{g}}_2}(\mathbb{V}_2^{\lambda,\mu}) \longrightarrow \Psi^0(\mathbb{V}_2^{\lambda,\mu})$ is a Z_2 -equivariant isomorphism. Recall from Lemma 4.28 in [4] that $Z_2[1/a]$ is dense in $Z_{t,s}$, and therefore the localization of \mathcal{H} is a $(Z_t \otimes_Q Z_s)$ -equivariant isomorphism

$$\operatorname{End}_{\hat{\mathfrak{g}}_t}(\mathbb{V}_t^{\lambda}) \otimes_Q \operatorname{End}_{\hat{\mathfrak{g}}_s}(\mathbb{V}_s^{\mu}) \longrightarrow \Psi^0(\mathbb{V}_t^{\lambda}) \otimes_Q \Psi^0(\mathbb{V}_s^{\mu}),$$

where we used the identification of the localization of $\Psi^0(\mathbb{V}_2^{\lambda,\mu})$ with $\Psi^0(\mathbb{V}_t^{\lambda}) \otimes_Q \Psi^0(\mathbb{V}_t^{\mu})$.

From Lemma 5.1 and 4.10 we deduce that $\mathcal{H}(\mathrm{Id}_{\mathbb{V}_{2}^{\lambda,\mu}}) = [q \, v_{\lambda} \otimes v_{\mu}]$, where $q \in A$ and $qv_{\lambda} \otimes v_{\mu} \in \mathbb{V}_{2}^{\lambda,\mu}$. We set $w = qv_{\lambda} \otimes v_{\mu} \in \mathbb{V}_{2}^{\lambda,\mu}$.

By specialisation, \mathcal{H} gives a Z_1 -equivariant isomorphism

$$\overline{\mathcal{H}}: \frac{\operatorname{End}_{\hat{\mathfrak{g}}_2}(\mathbb{V}_2^{\lambda,\mu})}{a\operatorname{End}_{\hat{\mathfrak{g}}_2}(\mathbb{V}_2^{\lambda,\mu})} \longrightarrow \frac{\Psi^0(\mathbb{V}_2^{\lambda,\mu})}{a\Psi^0(\mathbb{V}_2^{\lambda,\mu})}. \tag{5.1}$$

This isomorphism sends $\overline{\operatorname{Id}_{\mathbb{V}_{2}^{\lambda,\mu}}}$ to \overline{w} . Now consider the decomposition $V^{\lambda} \otimes V^{\mu} = \bigoplus V^{\nu}$ as \mathfrak{g} -modules. By Theorem 4.9, the target of the map $\overline{\mathcal{H}}$ in (5.1) decomposes as $\bigoplus \Psi^{0}(\mathbb{V}_{1}^{\lambda})$. The element w is a multiple of $v_{\lambda} \otimes v_{\mu}$ hence its class belongs to $\Psi^{0}(\mathbb{V}_{1}^{\lambda+\mu})$. As $\overline{\mathcal{H}}$ is Z_{1} -equivariant and $\mathbb{V}_{1}^{\lambda+\mu}$ is stable by the action of $\hat{\mathfrak{g}}_{1}$, we get that the image of $\overline{\mathcal{H}}$ is contained in the direct summand $\Psi^{0}(\mathbb{V}_{1}^{\lambda+\mu})$. In particular, if $V^{\lambda} \otimes V^{\mu}$ is not irreducible, the map $\overline{\mathcal{H}}$ cannot be surjective. This proves the

first claim. The second claim follows since the map from $Z_2^{\lambda,\mu}$ to $\Psi^0(\mathbb{V}_2^{\lambda,\mu})$ factors through $\operatorname{End}_{\hat{\mathfrak{g}}_2}(\mathbb{V}_2^{\lambda,\mu})$.

6. A Weyl module for $\mathfrak{sl}(2)$

In this Section, we propose an alternative Weyl module in the context of opers with two singularities, in the case of $\mathfrak{g}=\mathfrak{sl}(2)$. We fix the following notation: e,h,f is an $\mathfrak{sl}(2)$ -triple such that $h\in\mathfrak{t}$ and $e\in\mathfrak{n}_+$, while $\psi^*\in\mathfrak{n}_+^*$ is the dual of e. We identify dominant weights with natural numbers and we assume from now on that $\lambda\geqslant\mu$. In this case, the differential of the complex computing semi-infinite cohomology takes the simpler form $d^{(2)}=\psi^*+\sum ew_n\otimes\psi^*z_{-n-1/2}$.

Let $\widetilde{\mathbb{V}}_2^{\lambda,\mu}$ be the \hat{U}_2 -submodule of $\mathbb{V}_2^{\lambda,\mu}$ generated by the highest weight vector $1 \otimes v_\lambda \otimes v_\mu \in A \otimes V^\lambda \otimes V^\mu$. We will prove that this module has the "correct" semi-infinite cohomology and the "correct" endomorphism ring.

We start by giving a more explicit description of the module $\widetilde{\mathbb{V}}_2^{\lambda,\mu}$. If X is a subspace of $U(\mathfrak{g})$ and Y is a subspace of a \mathfrak{g} -module Z we denote by $X \cdot Y$ the subspace of Z generated by the products $x \cdot y$ with $x \in X$ and $y \in Y$. We define an increasing filtration F^i of $\widetilde{\mathbb{V}}_2^{\lambda,\mu}$ as follows

$$F^{i} = U(\mathfrak{g}) \cdot (\mathbb{C} \operatorname{Id} \otimes \operatorname{Id} + \operatorname{Id} \otimes \mathfrak{g})^{i} \cdot (v_{\lambda} \otimes v_{\mu}).$$

This is an increasing filtration of $V^{\lambda} \otimes V^{\mu}$ by \mathfrak{g} -modules and for i large enough we have $F^i = V^{\lambda} \otimes V^{\mu}$. Choose a \mathfrak{g} -stable complement G^{i+1} of F^i in F^{i+1} and set $G^0 = F^0$, so that $F^i = \bigoplus_{j=0}^i G^j$. If we set $F^i(V^{\mu}) = (\mathbb{C}\mathrm{Id} + \mathfrak{n}_-)^i v_{\mu}$, it is easy to check by induction on i that

$$F^{i} = U(\mathfrak{g}) \cdot (\operatorname{Id} \otimes \operatorname{Id} + \operatorname{Id} \otimes \mathfrak{n}_{-})^{i} (v_{\lambda} \otimes v_{\mu}) = U(\mathfrak{g}) \cdot (\mathbb{C}v_{\lambda} \otimes F^{i}(V^{\mu})).$$

In the case of $\mathfrak{g} = \mathfrak{sl}(2)$ we have $G^i \simeq V^{\lambda + \mu - 2i}$ and $F^{\mu} = V^{\lambda} \otimes V^{\mu}$.

Let $U_2^- \subset U(\hat{\mathfrak{g}}_2)$ be the A-span of Poincaré-Birkhoff-Witt monomials of the form $(x_1w_{a_1})\cdots(x_kw_{a_k})$ with $x_i\in\mathfrak{g}$ and $a_i<0$. This is a complement of $U(\hat{\mathfrak{g}}_2^+)$ in $U(\hat{\mathfrak{g}}_2)$, so that in particular we have

$$\mathbb{V}_{2}^{\lambda,\mu} = U_{2}^{-} \otimes_{\mathbb{C}} (V^{\lambda} \otimes V^{\mu}).$$

Lemma 6.1. *If* $\lambda \geqslant \mu$ *then*

$$\widetilde{\mathbb{V}}_{2}^{\lambda,\mu} = \sum_{i=0}^{\mu} a^{i} U_{2}^{-} \otimes_{\mathbb{C}} F^{i} = \bigoplus_{i=0}^{\mu} a^{i} U_{2}^{-} \otimes_{\mathbb{C}} G^{i}$$

Proof. To understand the module $\widetilde{\mathbb{V}}_2^{\lambda,\mu}$ we need to compute the $\hat{\mathfrak{g}}_2^+$ -submodule of $A \otimes_{\mathbb{C}} V^{\lambda} \otimes_{\mathbb{C}} V^{\mu}$ generated by $1 \otimes v_{\lambda} \otimes v_{\mu}$. Notice that every element of the form xg, with $x \in \mathfrak{g}$ and $g \in \mathbb{C}[[t,s]]$ divisible by ts, acts trivially on $A \otimes V^{\lambda} \otimes V^{\mu}$. Hence we need to understand the action of elements of the form

$$z = x_1 \cdots x_\ell \cdot (y_1 t) \cdots (y_m t) \cdot (v_\lambda \otimes v_\mu),$$

with $x_i, y_i \in \mathfrak{g}$. Moreover, elements of \mathfrak{g} act in the standard way on the tensor product $V^{\lambda} \otimes V^{\mu}$, while elements of the form xt with $x \in \mathfrak{g}$ act via $-a(\operatorname{Id} \otimes x)$. This implies the lemma.

We now describe the specialisation of the module $\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}$. We introduce the following decreasing filtration of $\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}$:

$$\mathbb{F}_i = \widetilde{\mathbb{V}}_2^{\lambda,\mu} \cap a^i \mathbb{V}_2^{\lambda,\mu}. \tag{6.1}$$

By Lemma 6.1 we have the following description of the terms of this filtration as A-modules:

$$\mathbb{F}_i = a^i U_2^- \otimes_{\mathbb{C}} F^i \oplus \bigoplus_{j=i+1}^{\mu} a^j U_2^- \otimes_{\mathbb{C}} G^j$$

In particular we have $\mathbb{F}_0 = \widetilde{\mathbb{V}}_2^{\lambda,\mu}$, $\mathbb{F}_j = a^j \mathbb{V}_2^{\lambda,\mu}$ for $j \geqslant \mu$.

Lemma 6.2. a) Let $u_i \in G^i$ be the highest weight vector and set $\tilde{w}_i = a^i u_i$. Then $\tilde{w}_i \in \widetilde{\mathbb{V}}_2^{\lambda,\mu}$ and $a^{i-1}u_i \notin \widetilde{\mathbb{V}}_2^{\lambda,\mu}$.

b) There is an isomorphism of \hat{U}_1 -modules

$$\frac{\mathbb{F}_i + a\widetilde{\mathbb{V}}_2^{\lambda,\mu}}{a\widetilde{\mathbb{V}}_2^{\lambda,\mu}} \simeq \bigoplus_{j=i}^{\mu} \mathbb{V}_1^{\lambda+\mu-2j}.$$

The quotient $\frac{\mathbb{F}_i + a\widetilde{\mathbb{V}}_2^{\lambda,\mu}}{a\widetilde{\mathbb{V}}_2^{\lambda,\mu}}$ is generated as a \hat{U}_1 -module by the classes of $\tilde{w}_i, \ldots, \tilde{w}_{\mu}$. In particular $\widetilde{\mathbb{V}}_2^{\lambda,\mu}/a\widetilde{\mathbb{V}}_2^{\lambda,\mu} \simeq \mathbb{W}_1^{\lambda,\mu}$ is generated by $\tilde{w}_0, \ldots, \tilde{w}_{\mu}$.

Proof. The first claim follows from Lemma 6.1.

We prove part b) by decreasing induction on i. By Lemma 6.1, for $i > \mu$ the quotient is zero and the claim is true. For $i \leq \mu$, consider the map

$$U_2^- \otimes G^i \longrightarrow \frac{\mathbb{F}_i + a^{i+1} \mathbb{V}_2^{\lambda,\mu}}{a^{i+1} \mathbb{V}_2^{\lambda,\mu} + \mathbb{F}_i \cap a \widetilde{\mathbb{V}}_2^{\lambda,\mu}} \simeq \frac{(\mathbb{F}_i + a \widetilde{\mathbb{V}}_2^{\lambda,\mu}) / a \widetilde{\mathbb{V}}_2^{\lambda,\mu}}{(\mathbb{F}_{i+1} + a \widetilde{\mathbb{V}}_2^{\lambda,\mu}) / a \widetilde{\mathbb{V}}_2^{\lambda,\mu}}$$

sending an element $u \otimes v$ to the class of $a^i u \otimes v$. This map induces an isomorphism

$$\frac{U_2^-}{aU_2^-} \otimes G^i \simeq \frac{(\mathbb{F}_i + a\widetilde{\mathbb{V}}_2^{\lambda,\mu})/a\widetilde{\mathbb{V}}_2^{\lambda,\mu}}{(\mathbb{F}_{i+1} + a\widetilde{\mathbb{V}}_2^{\lambda,\mu})/a\widetilde{\mathbb{V}}_2^{\lambda,\mu}}.$$
 (6.2)

Moreover, notice that $\frac{U_2^-}{aU_2^-}\otimes G^i\simeq U_1^-\otimes G^i$, where $U_1^-=U(t^{-1}\mathfrak{g}[t^{-1}])\subset U(\hat{\mathfrak{g}}_1)=0$ U_1 , and that $U_1^- \otimes G^i$ has a natural structure of U_1 -module, as it can be identified with $\mathbb{V}_1^{\lambda+\mu-2i}$. With this U_1 -action, the isomorphism 6.2 is U_1 -equivariant. Now the claim follows by the inductive hypothesis, combined with the fact that there are no nontrivial extensions between modules \mathbb{V}_1^{ν} and $\mathbb{V}_1^{\nu'}$ if $\nu \neq \nu'$ and that the highest weight vector of V^{ν} generates the module \mathbb{V}_{1}^{ν} as an U_{1} -module.

Notice that, although the specialisations at a=0 of $\mathbb{V}_2^{\lambda,\mu}$ and $\widetilde{\mathbb{V}}_2^{\lambda,\mu}$ are isomorphic, the specialisation of $\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}$, is generated by $v_{\lambda}\otimes v_{\mu}$ while in the first case this vector generates the submodule $\mathbb{V}_{1}^{\lambda+\mu}$.

As a corollary, we get the following result.

Proposition 6.3. The following hold:

- $\begin{array}{l} a) \ \Psi^n(\widetilde{\mathbb{V}}_2^{\lambda,\mu}) = 0 \ for \ n \neq 0. \\ b) \ The \ inclusion \ of \ \widetilde{\mathbb{V}}_2^{\lambda,\mu} \ \ in \ \mathbb{V}_2^{\lambda,\mu} \ \ induces \ isomorphisms \end{array}$

$$\Psi^0(\widetilde{\mathbb{V}}_2^{\lambda,\mu})[a^{-1}] \simeq \Psi^0(\mathbb{V}_2^{\lambda,\mu})[a^{-1}] \simeq \Psi^0(\mathbb{V}_t^{\lambda}) \otimes_Q \Psi^0(\mathbb{V}_s^{\mu}).$$

c) $\Psi^0(\widetilde{\mathbb{Y}}_2^{\lambda,\mu})$ is torsion-free with respect to the action of A, and the natural projection induces isomorphisms

$$\frac{\Psi^0(\widetilde{\mathbb{V}}_2^{\lambda,\mu})}{a\Psi^0(\widetilde{\mathbb{V}}_2^{\lambda,\mu})} \simeq \Psi^0\left(\frac{\widetilde{\mathbb{V}}_2^{\lambda,\mu}}{a\widetilde{\mathbb{V}}_2^{\lambda,\mu}}\right) \simeq \Psi^0(\mathbb{W}_1^{\lambda,\mu}).$$

Proof. We use the filtration introduced in Equation (6.1). Notice that

$$\frac{\mathbb{F}_i}{\mathbb{F}_{i+1}} = \frac{a^i U_2^- \otimes F^i}{a^{i+1} U_2^- \otimes F^i} \simeq U_1^- \otimes_{\mathbb{C}} F^i \simeq \operatorname{Ind}_{\hat{\mathfrak{g}}_1^+}^{\hat{\mathfrak{g}}_1} F^i,$$

where we consider F^i as a $\hat{\mathfrak{g}}_1^+$ -module on which $t\mathfrak{g}[t]$ acts trivially. Notice that $\operatorname{Ind}_{\hat{\mathfrak{g}}_1^+}^{\hat{\mathfrak{g}}_1} F^i$ is a sum of modules of the form \mathbb{V}_1^{ν} , hence in particular has trivial non-zero cohomology.

Hence, arguing by decreasing induction on i, starting from $i = \mu$, it follows that \mathbb{F}_i has trivial semi-infinite cohomology in degree different from zero. Indeed for $i = \mu$ we have $\mathbb{F}_{\mu} = a^{\mu} \mathbb{V}_{2}^{\lambda,\mu} \simeq \mathbb{V}_{2}^{\lambda,\mu}$ and this is the content of Theorem 4.9. For i = 0 this implies claim a).

Part b) follows from the fact that semi-infinite cohomology commutes with localization (see Lemma 3.3) combined with the isomorphism $\hat{\mathbb{V}}_{2}^{\lambda,\mu}[a^{-1}] = \mathbb{V}_{2}^{\lambda,\mu}[a^{-1}] \simeq \mathbb{V}_{t}^{\lambda} \otimes_{Q} \mathbb{V}_{s}^{\mu}$.

To prove c), consider the exact sequence

$$0 \longrightarrow \widetilde{\mathbb{V}}_{2}^{\lambda,\mu} \stackrel{\cdot a}{\longrightarrow} \widetilde{\mathbb{V}}_{2}^{\lambda,\mu} \longrightarrow \frac{\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}}{a\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}} \longrightarrow 0$$

By Lemma 6.2, the last module in this sequence is isomorphic to $\mathbb{W}_{1}^{\lambda,\mu}$. In particular, the semi-infinite cohomology groups Ψ^{n} of the modules appearing in this sequence are zero for $n \neq 0$, and c) follows.

To prove that the semi-infinite cohomology of $\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}$ is isomorphic to $Z_{2}^{\lambda,\mu}$ we will use the action of a particular central element in Z_{2} . Recall from [4] the definition of the 2-Sugawara operator

$$S_{1/2}^{(2)} = \sum_{n \in \frac{1}{2}\mathbb{Z}, b} : (J^b w_n) (J_b z_{-n}) :$$
 (6.3)

where J^1, J^2, J^3 are the basis elements e, h, f and J_1, J_2, J_3 are the dual basis elements f, h/2, e. As proved in [4], the element $S_{1/2}^{(2)}$ is central. Its specialisation is the Sugawara operator

$$S_1^{(1)} = \sum_{n \in \mathbb{Z}_b} : (J^b t^n) (J_b t^{-n}) : (6.4)$$

which is an element of Z_1 . It is straightforward to check that the action of $S_1^{(1)}$ on the Weyl module \mathbb{V}_1^{ν} is given by multiplication by $\nu(\nu+1)$.

Lemma 6.4. The element $\hat{w}_{\ell} = (et^{-1})^{\ell} \tilde{w}_{\ell}$ belongs to $Z_2 \cdot (v_{\lambda} \otimes v_{\mu}) + a \widetilde{\mathbb{V}}_{2}^{\lambda,\mu}$ for $\ell = 0, \ldots, \mu$,

Proof. We notice first that the element $v_{\lambda} \otimes f^{\ell}v_{\mu}$ belongs to $F^{\ell} \setminus F^{\ell-1}$ and has weight $\lambda + \mu - 2\ell$. Hence, up to a non-zero constant we have $v_{\lambda} \otimes f^{\ell}v_{\mu} = u_{\ell} + u'_{\ell}$, where we recall that u_{ℓ} is the highest weight vector in $G^{\ell} \simeq V^{\lambda + \mu - 2\ell} \subset V^{\lambda} \otimes V^{\mu}$ and $u'_{\ell} \in F^{\ell-1}$. In particular, recall from Lemma 6.2 that $a^{\ell-1}F^{\ell} \subset \widetilde{\mathbb{V}}_{2}^{\lambda,\mu}$, hence

$$a^{\ell} \left(et^{-1}\right)^{\ell} v_{\lambda} \otimes f^{\ell} v_{\mu} = \left(et^{-1}\right)^{\ell} \tilde{w}_{\ell} + \left(et^{-1}\right)^{\ell} \left(a^{\ell} u_{\ell}'\right) \equiv \left(et^{-1}\right)^{\ell} \tilde{w}_{\ell} \mod a \widetilde{\mathbb{V}}_{2}^{\lambda,\mu}.$$

Hence, the lemma is equivalent to the fact that $\hat{w}_\ell = a^\ell \left(et^{-1}\right)^\ell v_\lambda \otimes f^\ell v_\mu$ is in $Z_2 \cdot v_\lambda \otimes v_\mu + a\widetilde{\mathbb{V}}_2^{\lambda,\mu}$. We prove this statement by induction on ℓ . For $\ell = 0$ it is trivially true. Now assume \hat{w}_ℓ is in $Z_2 \cdot v_\lambda \otimes v_\mu + a\widetilde{\mathbb{V}}_2^{\lambda,\mu}$. We compute $S_{1/2}^{(2)}(\hat{w}_\ell)$. In order to do this, we notice that the action of xt^is^j on $\widetilde{\mathbb{V}}_2^{\lambda,\mu}/a\widetilde{\mathbb{V}}_2^{\lambda,\mu}$ is equal to the action of xt^{i+j} on the same module, and that $v_\lambda \otimes ef^\ell v_\mu$ is in $F^{\ell-1}$. We have

$$S_{1/2}^{(2)}\hat{w}_{\ell} = 2\sum_{n>0} et^{-n} \cdot ft^{n} \cdot \hat{w}_{\ell} + 2\sum_{n>0} ft^{-n} \cdot et^{n} \cdot \hat{w}_{\ell} + \sum_{n>0} ht^{-n} \cdot ht^{n} \cdot \hat{w}_{\ell} + e \cdot f \cdot \hat{w}_{\ell} + e \cdot f \cdot \hat{w}_{\ell} + \frac{1}{2}h \cdot h \cdot \hat{w}_{\ell}.$$

In the second infinite sum above, the element et^n commutes with et^{-1} , hence $et^n \cdot \hat{w}_{\ell} \in a\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}$ for all n > 0. The summands of the third series are of the form

$$ht^n \cdot (et^{-1})^{\ell} \cdot \hat{w}_{\ell} = (et^{-1})^{\ell} ht^n \cdot \hat{w}_{\ell} + 2\ell(et^{-1})^{\ell-1} et^{n-1} \cdot \hat{w}_{\ell},$$

hence they vanish for $n \ge 3$, while for n = 1, 2 they are easily checked to be elements of $a\widetilde{\mathbb{V}}_2^{\lambda,\mu}$. The summands of the first series are given by

$$ft^n \cdot (et^{-1})^{\ell} \cdot \hat{w}_{\ell} = (et^{-1})^{\ell} ft^n \cdot \hat{w}_{\ell} - \ell(et^{-1})^{\ell-1} ht^{n-1} \cdot \hat{w}_{\ell} - \ell(\ell-1)(et^{-1})^{\ell-2} et^{n-2} \cdot \hat{w}_{\ell},$$

and all terms are zero or in $a\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}$ but for the case n=1, for which we get

$$(et^{-1}) \cdot (ft) \cdot (et^{-1})^{\ell} \cdot \hat{w}_{\ell} = a^{\ell+1} (et^{-1})^{\ell+1} \cdot (v_{\lambda} \otimes f^{\ell+1} v_{\mu}) - \ell (et^{-1})^{\ell} h \cdot (v_{\lambda} \otimes f^{\ell} v_{\mu}) - \ell (\ell-1) (et^{-1})^{\ell} \cdot \hat{w}_{\ell} = \hat{w}_{\ell+1} + K_{1} \hat{w}_{\ell}$$

for some constant K_1 . Finally, $e \cdot f \cdot \hat{w}_{\ell} + e \cdot f \cdot \hat{w}_{\ell} + \frac{1}{2} h \cdot h \cdot \hat{w}_{\ell}$ belongs to $K_2 \hat{w}_{\ell} + a \widetilde{\mathbb{V}}_2^{\lambda,\mu}$ for some constant K_2 . Hence we get

$$S_{1/2}^{(2)}\hat{w}_{\ell} \equiv \hat{w}_{\ell+1} + K\hat{w}_{\ell} \bmod a\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}$$

for some constant K, proving our claim.

We now prove that the zero-th semi-infinite cohomology of the module $\widetilde{\mathbb{V}}_2^{\lambda,\mu}$ is isomorphic to $Z_2^{\lambda,\mu}$.

Theorem 6.5. For $\mathfrak{g} = \mathfrak{sl}(2)$ the map $\Phi: Z_2^{\lambda,\mu} \longrightarrow \Psi^0(\widetilde{\mathbb{V}}_2^{\lambda,\mu})$ given by $\Phi(z) = z \cdot [v_{\lambda} \otimes v_{\mu}]$ is an isomorphism.

Proof. By [4], Theorem 6.4, the action of Z_2 on $\mathbb{V}_2^{\lambda,\mu}$, hence on $\widetilde{\mathbb{V}}_2^{\lambda,\mu}$, factors through $Z_2^{\lambda,\mu}$. Moreover $v_{\lambda} \otimes v_{\mu}$ is a cycle, so the map Φ is well defined. Since we know that both modules are torsion-free, to prove that Φ is an isomorphism it suffices to prove that the localization Φ_a and the specialisation $\overline{\Phi}$ are isomorphisms.

The fact that Φ_a is an isomorphism is the content of part b) of Proposition 6.3. We need to prove that $\overline{\Phi}$ is an isomorphism. By Lemma 6.2, Proposition 6.3 and [4, Theorem 2.13] we have

$$\frac{Z_2^{\lambda,\mu}}{aZ_2^{\lambda,\mu}} \simeq \prod_{i=0}^{\mu} Z_1^{\lambda+\mu-2i} \qquad \text{and} \qquad \frac{\Psi^0(\widetilde{\mathbb{V}}_2^{\lambda,\mu})}{a\Psi^0(\widetilde{\mathbb{V}}_2^{\lambda,\mu})} \simeq \bigoplus_{i=0}^{\mu} \Psi^0(\mathbb{V}_1^{\lambda+\mu-2i}).$$

In particular, by Theorem 4.1 these two Z_1 -modules are isomorphic, but we need to prove that our specific map $\overline{\Phi}$ provides an isomorphism between them. By Remark 5.2 it is enough to prove that $\overline{\Phi}$ is surjective. We prove that the image of $\overline{\Phi}$ contains $\Psi^0(\mathbb{F}_\ell + a\widetilde{\mathbb{V}}_2^{\lambda,\mu}/a\widetilde{\mathbb{V}}_2^{\lambda,\mu})$ arguing by reverse induction on ℓ . For $\ell = 0$ we get our claim. For $\ell > \mu$ there is nothing to prove. Now assume $\ell \leqslant \mu$. Consider again the exact sequence

$$0 {\longrightarrow} \xrightarrow{\mathbb{F}_{\ell+1} + a\widetilde{\mathbb{V}}_2^{\lambda,\mu}} {\longrightarrow} \xrightarrow{\mathbb{F}_{\ell} + a\widetilde{\mathbb{V}}_2^{\lambda,\mu}} {\longrightarrow} a^{\ell}U_2^- \otimes_{\mathbb{C}} G^{\ell} {\longrightarrow} 0.$$

We know that the last module is isomorphic to

$$a^\ell U_2^- \otimes_{\mathbb{C}} G^\ell \simeq \mathbb{V}_1^{\lambda + \mu - 2\ell} = \operatorname{Ind}_{\hat{\mathfrak{g}}_1^+}^{\hat{\mathfrak{g}}_1}(V^{\lambda + \mu - 2\ell})$$

and that it is generated by the element $\tilde{w}_{\ell} \in a^{\ell}G^{\ell}$. Notice this sequence of Z_1 -modules splits by Remark 5.2. Taking semi-infinite cohomology we get a short exact sequence

$$0 \longrightarrow \Psi^0\left(\frac{\mathbb{F}_{\ell+1} + a\widetilde{\mathbb{V}}_2^{\lambda,\mu}}{a\widetilde{\mathbb{V}}_2^{\lambda,\mu}}\right) \longrightarrow \Psi^0\left(\frac{\mathbb{F}_{\ell} + a\widetilde{\mathbb{V}}_2^{\lambda,\mu}}{a\widetilde{\mathbb{V}}_2^{\lambda,\mu}}\right) \longrightarrow \Psi^0\left(a^\ell U_2^- \otimes_{\mathbb{C}} G^\ell\right) \longrightarrow 0.$$

and we know that the last Z_2 -module is generated by \tilde{w}_{ℓ} . Hence it is enough to prove that this element is in the image of $Z_2^{\lambda,\mu}(v_{\lambda}\otimes v_{\mu})$ in $\Psi^0\left(\widetilde{\mathbb{V}}_2^{\lambda,\mu}/\mathbb{F}_{\ell+1}+a\widetilde{\mathbb{V}}_2^{\lambda,\mu}\right)$.

By Lemma 6.4 we know that \hat{w}_{ℓ} is in this image. Now we prove that \tilde{w}_{ℓ} and \hat{w}_{ℓ} define the same element in the semi-infinite cohomology of $a^{\ell}U_2^-\otimes_{\mathbb{C}}G^{\ell}$. This is a claim about the cohomology of the module \mathbb{V}_1^{ν} for $\nu=\lambda+\mu-2\ell$. For any ν we prove that $(et^{-1})^h v_{\nu} + (et^{-1})^{h-1} v_{\nu}$ is a coboundary. Indeed the boundary operator in the case of $\mathfrak{sl}(2)$ is equal to

$$d^{(1)} = \psi^* + \sum_{n \in \mathbb{Z}} (et^n) \otimes \psi^* t^{-1-n},$$

so a simple computation shows

$$d^{(1)}\left(\left(et^{-1}\right)^{h-1}v_{\nu}\otimes(\psi t^{-1})|0\rangle_{\Lambda}\right)=\left(et^{-1}\right)^{h-1}v_{\nu}\otimes|0\rangle_{\Lambda}+\left(et^{-1}\right)^{h}v_{\nu}\otimes|0\rangle_{\Lambda},$$

which implies our claim.

Recall that in [4] we computed the endomorphism ring of $\mathbb{V}_2^{\lambda,\mu}$, showing that it is isomorphic to $Z_2^{\lambda,\mu}$. We now prove the same result for the module $\widetilde{\mathbb{V}}_2^{\lambda,\mu}$.

Proposition 6.6. The action of the center Z_2 on $\widetilde{\mathbb{V}}_2^{\lambda,\mu}$ induces an isomorphism

$$Z_2^{\lambda,\mu} \simeq \operatorname{End}_{\hat{\mathfrak{g}}_2}(\widetilde{\mathbb{V}}_2^{\lambda,\mu}).$$

Proof. We already recalled at the beginning of the proof of Theorem 6.5 that the action of Z_2 on $\widetilde{\mathbb{V}}_2^{\lambda,\mu}$ factors through $Z_2^{\lambda,\mu}$. We denote by $\alpha:Z_2^{\lambda,\mu}\longrightarrow \operatorname{End}(\widetilde{\mathbb{V}}_2^{\lambda,\mu})$ this action. Since both modules have no A-torsion, in order to prove that α is an isomorphism it suffices to show that its localization and its specialisation are isomorphisms. Moreover, since our modules are finitely generated and have no torsion we have

$$\begin{split} \operatorname{End}_{\widehat{\mathfrak{g}}_2}\left(\widetilde{\mathbb{V}}_2^{\lambda,\mu}\right)[a^{-1}] &\simeq \operatorname{End}_{\widehat{\mathfrak{g}}_2[a^{-1}]}\left(\widetilde{\mathbb{V}}_2^{\lambda,\mu}[a^{-1}]\right) \simeq \operatorname{End}_{\widehat{\mathfrak{g}}_{t,s}}\left(\mathbb{V}^{\lambda} \otimes_Q \mathbb{V}_s^{\mu}\right) \\ &\simeq Z_t^{\lambda} \otimes_Q Z_t^{\mu} \simeq Z_2^{\lambda,\mu}[a^{-1}], \end{split}$$

hence the localization of α is an isomorphism.

Finally, we prove that the specialisation of α is also an isomorphism. We have already recalled that by [4, Theorem 2.13] we have $Z_2^{\lambda,\mu}/aZ_2^{\lambda,\mu} \simeq \prod_{i=0}^{\mu} Z_1^{\lambda+\mu-2i}$. Hence by Theorem 4.1 we have the following abstract isomorphisms of Z_1 -modules:

$$\frac{Z_2^{\lambda,\mu}}{aZ_2^{\lambda,\mu}} \simeq \prod_{i=0}^{\mu} Z_1^{\lambda+\mu-2i} \simeq \prod_{i=0}^{\mu} \operatorname{End}_{\hat{\mathfrak{g}}_1}(\mathbb{V}_1^{\lambda+\mu-2i}).$$

Moreover, since $\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}$ has no nontrivial A-torsion, by Lemma 6.2 and Remark 5.2 part (1) we have the inclusion

$$\frac{\operatorname{End}_{\hat{\mathfrak{g}}_{1}}\left(\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}\right)}{a\operatorname{End}_{\hat{\mathfrak{g}}_{1}}\left(\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}\right)}\subset\operatorname{End}_{\hat{\mathfrak{g}}_{1}}\left(\frac{\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}}{a\widetilde{\mathbb{V}}_{2}^{\lambda,\mu}}\right)\simeq\prod_{i=0}^{\mu}\operatorname{End}_{\hat{\mathfrak{g}}_{1}}(\mathbb{V}_{1}^{\lambda+\mu-2i}).$$

Hence, composing the specialisation of the map α with this inclusion and the isomorphisms above we get a Z_1 -equivariant endomorphism of $\prod_{i=0}^{\mu} Z_1^{\lambda+\mu-2i}$. Moreover, $\alpha(1)=1$, hence we conclude by Remark 5.2 (3) that the specialisation of α is also an isomorphism.

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