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### Research Article

# **Fourth-Order Differential Equation with Deviating Argument**

### M. Bartušek, M. Cecchi, Z. Došlá, and M. Marini 2

Correspondence should be addressed to Z. Došlá, dosla@math.muni.cz

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We consider the fourth-order differential equation with middle-term and deviating argument  $x^{(4)}(t) + q(t)x^{(2)}(t) + r(t)f(x(\varphi(t))) = 0$ , in case when the corresponding second-order equation h'' + q(t)h = 0 is oscillatory. Necessary and sufficient conditions for the existence of bounded and unbounded asymptotically linear solutions are given. The roles of the deviating argument and the nonlinearity are explained, too.

#### 1. Introduction

The aim of this paper is to investigate the fourth-order nonlinear differential equation with middle-term and deviating argument

$$x^{(4)}(t) + q(t)x^{(2)}(t) + r(t)f(x(\varphi(t))) = 0.$$
(1.1)

The following assumptions will be made.

(i) q is a continuously differentiable bounded away from zero function, that is,  $q(t) \ge q_0 > 0$  for large t such that

$$\int_0^\infty |q'(t)| dt < \infty. \tag{1.2}$$

(ii)  $r, \varphi$  are continuous functions for  $t \ge 0$ , r is not identically zero for large  $t, \varphi(t) \ge 0$ , and  $\varphi(0) = 0$ ,  $\lim_{t \to \infty} \varphi(t) = \infty$ .

<sup>&</sup>lt;sup>1</sup> Department of Mathematics and Statistics, Masaryk University, 61137 Brno, Czech Republic

<sup>&</sup>lt;sup>2</sup> Department of Electronics and Telecommunications, University of Florence, 50139 Florence, Italy

(iii) f is a continuous function such that f(u)u > 0 for  $u \neq 0$ .

Observe that (i) implies that there exists a positive constant Q such that  $q(t) \leq Q$  and the linear second-order equation

$$h''(t) + q(t)h(t) = 0 (1.3)$$

is oscillatory. Moreover, solutions of (1.3) are bounded together with their derivatives, see for example, [1, Theorem 2].

By a solution of (1.1) we mean a function x defined on  $[T_x, \infty)$ ,  $T_x \ge 0$ , which is differentiable up to the fourth order and satisfies (1.1) on  $[T_x, \infty)$  and  $\sup\{|x(t)| : t \ge T\} > 0$  for  $T \ge T_x$ .

A solution x of (1.1) is said to be asymptotically linear (AL-solution) if either

$$\lim_{t \to \infty} x(t) = c_x \neq 0, \qquad \lim_{t \to \infty} x'(t) = 0, \tag{1.4}$$

or

$$\lim_{t \to \infty} |x(t)| = \infty, \qquad \lim_{t \to \infty} x'(t) = d_x \neq 0, \tag{1.5}$$

for some constants  $c_x$ ,  $d_x$ .

Fourth-order nonlinear differential equations naturally appear in models concerning physical, biological, and chemical phenomena, such as, for instance, problems of elasticity, deformation of structures, or soil settlement, see, for example, [2, 3].

When (1.3) is nonoscillatory and h is its eventually positive solution, it is known that (1.1) can be written as the two-term equation

$$\left(h^{2}(t)\left(\frac{x''(t)}{h(t)}\right)'\right)' + h(t)r(t)f(x(t)) = 0.$$
(1.6)

In this case, the question of oscillation and asymptotics of such class of equations has been investigated with sufficient thoroughness, see, for example, the papers [3–10] or the monographs [11, 12] and references therein.

Nevertheless, as far we known, there are only few results concerning (1.1) when (1.3) is oscillatory. For instance, the equation without deviating argument

$$x^{(n)}(t) + q(t)x^{(n-2)}(t) + r(t)f(x(t)) = 0$$
(1.7)

has been investigated by Kiguradze in [13] in case  $q(t) \equiv 1$  and by the authors in [14, 15] when q satisfies (i). In particular, in [14] the oscillation of (1.1) in the case n=3 is studied. In [15], the existence of positive bounded and unbounded solutions as well as of oscillatory solutions for (1.7) has been considered and the case n=4 has been analyzed in detail. Other results can be found in [16] and references therein, in which the existence and uniqueness of almost periodic solutions for equations of type (1.1) with almost periodic coefficients q, r are studied.

Motivated by [14, 15], here we study the existence of AL-solutions for (1.1). The approach is completely different from the one used in [15], in which an iteration process, jointly with a comparison with the linear equation  $y^{(4)} + q(t)y^{(2)} = 0$ , is employed. Our tools are based on a topological method, certain integral inequalities, and some auxiliary functions. In particular, for proving the continuity in the Fréchet space  $C[t_0, \infty)$  of the fixed point operators here considered, we use a similar argument to that in the Vitali convergence theorem.

Our results extend to the case with deviating argument analogues ones stated in [15] for (1.7) when n = 4. We obtain sharper conditions for the existence of unbounded AL-solutions of (1.1), and, in addition, we show that under additional assumptions on q, r, these conditions become also necessary for the existence of AL-solutions, in both the bounded and unbounded cases. In the final part, we consider the particular case

$$f(u) = |u|^{\lambda} \operatorname{sgn} u \quad (\lambda > 0)$$
(1.8)

and we study the possible coexistence of bounded and unbounded AL-solutions. The role of deviating argument and the one of the growth of the nonlinearity are also discussed and illustrated by some examples.

#### 2. Unbounded Solutions

Here we study the existence of unbounded AL-solutions of (1.1). Our first main result is the following.

**Theorem 2.1.** For any c,  $0 < c < \infty$ , there exists an unbounded solution x of (1.1) such that

$$\lim_{t \to \infty} x'(t) = c, \qquad \lim_{t \to \infty} x^{(i)}(t) = 0, \quad i = 2, 3, \tag{2.1}$$

provided

$$\int_{0}^{\infty} |r(t)| F(\varphi(t)) dt < \infty, \tag{2.2}$$

where for u > 0

$$F(u) = \max \left\{ f(v) : |v - u| \le \frac{1}{2}u \right\}. \tag{2.3}$$

*Proof.* Without loss of generality, we prove the existence of solutions of (1.1) satisfying (2.1) for c = 1.

Let u and v be two linearly independent solutions of (1.3) with Wronskian d=1. Denote

$$w(s,t) = u(s)v(t) - u(t)v(s), z(s,t) = \frac{\partial}{\partial t}w(s,t). (2.4)$$

As claimed by the assumptions on q, all solutions of (1.3) and their derivatives are bounded. Thus, put

$$M = \sup\{|w(s,t)| + |z(s,t)| : s \ge 0, \ t \ge 0\}, \qquad L = \frac{2(2M+1)}{q_0}. \tag{2.5}$$

Let  $\bar{t} \ge t_0$  be such that  $\varphi(t) \ge t_0$  for  $t \ge \bar{t}$ . Define

$$\overline{\varphi}(t) = \begin{cases} \varphi(t) & \text{if } t \ge \overline{t}, \\ \varphi(\overline{t}) & \text{if } t_0 \le t \le \overline{t}, \end{cases}$$
(2.6)

and choose  $t_0 \ge 0$  large so that

$$\int_{t_0}^{\infty} |r(s)| F(\overline{\varphi}(s)) ds \le \frac{1}{2L}, \qquad \frac{1}{q_0} \int_{t_0}^{\infty} |q'(t)| dt \le \frac{1}{2}. \tag{2.7}$$

Denote by  $C[t_0, \infty)$  the Fréchet space of all continuous functions on  $[t_0, \infty)$ , endowed with the topology of uniform convergence on compact subintervals of  $[t_0, \infty)$ , and consider the set  $\Omega \subset C[t_0, \infty)$  given by

$$\Omega = \left\{ x \in C[t_0, \infty) : \frac{t}{2} \le x(t) \le \frac{3t}{2} \right\}. \tag{2.8}$$

Let  $T > t_0$  and define on  $[t_0, T]$  the function

$$g(t) = \gamma''(t) + q(t)\gamma(t), \tag{2.9}$$

where

$$\gamma(t) = -\int_{t}^{T} \int_{-\tau}^{\infty} r(s) f\left(x(\overline{\varphi}(s))\right) w(s,\tau) ds d\tau \tag{2.10}$$

and  $x \in \Omega$ . Then,

$$\gamma'(t) = \int_{t}^{\infty} r(s) f(x(\overline{\varphi}(s))) w(s, t) ds, \tag{2.11}$$

$$\gamma''(t) = \int_{t}^{\infty} r(s) f(x(\overline{\varphi}(s))) z(s, t) ds,$$

$$\gamma'''(t) = -r(t) f(x(\overline{\varphi}(t))) - q(t) \gamma'(t).$$
(2.12)

Moreover,  $g(T) = \gamma''(T)$ , and it holds for  $t \in [t_0, T]$  that

$$g'(t) = \gamma'''(t) + q(t)\gamma'(t) + q'(t)\gamma(t) = -r(t)f(x(\overline{\varphi}(t))) + q'(t)\gamma(t). \tag{2.13}$$

Integrating, we obtain

$$g(t) = g(T) - \int_{t}^{T} g'(s)ds = \gamma''(T) + \int_{t}^{T} r(s)f\left(x(\overline{\varphi}(s))\right)ds - \int_{t}^{T} q'(s)\gamma(s)ds. \tag{2.14}$$

From here and (1.3), we get

$$\gamma(t) = \frac{1}{q(t)} \left( \gamma''(T) - \gamma''(t) + \int_{t}^{T} r(s) f(x(\overline{\varphi}(s))) ds - \int_{t}^{T} q'(s) \gamma(s) ds \right). \tag{2.15}$$

Thus,

$$\left|\gamma(t)\right| \leq \frac{1+2M}{q_0} \int_t^{\infty} |r(s)| F(\overline{\varphi}(s)) ds + \frac{1}{q_0} \max_{t \leq s \leq T} \left|\gamma(s)\right| \int_{t_0}^{\infty} |q'(s)| ds, \tag{2.16}$$

and so

$$\left(1 - \frac{1}{q_0} \int_{t_0}^{\infty} \left| q'(s) \right| ds \right) \max_{t \le s \le T} \left| \gamma(s) \right| \le \frac{1 + 2M}{q_0} \int_{t}^{\infty} |r(s)| F(\overline{\varphi}(s)) ds, \tag{2.17}$$

or, in view of (2.7),

$$|\gamma(t)| \le L \int_{t}^{\infty} |r(s)| F(\overline{\varphi}(s)) ds.$$
 (2.18)

Thus, from (2.10), as  $T \to \infty$ , we get

$$\left| \int_{t}^{\infty} \int_{\tau}^{\infty} r(s) f\left(x(\overline{\varphi}(s))\right) w(s,\tau) ds d\tau \right| \le L \int_{t}^{\infty} |r(s)| F(\overline{\varphi}(s)) ds. \tag{2.19}$$

Hence, the operator  $T: \Omega \to \Omega$  given by

$$\mathcal{T}(x)(t) = t - \int_{t_0}^t \int_{\sigma}^{\infty} \int_{\tau}^{\infty} r(s) f\left(x(\overline{\varphi}(s))\right) w(s, \tau) ds d\tau d\sigma \tag{2.20}$$

is well defined for any  $x \in \Omega$ . Moreover, in view of (2.19), we have

$$\left| \mathcal{T}'(x)(t) - 1 \right| \le L \int_{t}^{\infty} |r(s)| F(\overline{\varphi}(s)) ds. \tag{2.21}$$

From here, in virtue of (2.7) we get

$$|\mathcal{T}(x)(t) - t| \le Lt \int_{t_0}^{\infty} |r(s)| F(\overline{\varphi}(s)) ds \le \frac{1}{2}t. \tag{2.22}$$

Hence,  $\mathcal{T}(\Omega) \subset \Omega$ . From (2.5) and (2.11), we have

$$\left| \mathcal{T}''(x)(t) \right| = \left| \gamma'(t) \right| \le M \int_{t}^{\infty} |r(s)| F(\overline{\varphi}(s)) ds, \tag{2.23}$$

and so  $\lim_{t\to\infty} T''(x)(t) = 0$ . Similarly,

$$\left|\mathcal{T}'''(x)(t)\right| = \left|\gamma''(t)\right| \le M \int_{t}^{\infty} |r(s)| F(\overline{\varphi}(s)) ds, \tag{2.24}$$

and thus,  $\lim_{t\to\infty} T'''(x)(t) = 0$ , too. In addition,

$$C^{(4)}(x)(t) = \gamma'''(t) = -q(t)C''(x)(t) - r(t)f(x(\overline{\varphi}(t))).$$
(2.25)

Hence, any fixed point of T is a solution of (1.1) for large t.

Let us show that  $\mathcal{T}(\Omega)$  is relatively compact, that is,  $\mathcal{T}(\Omega)$  consists of functions equibounded and equicontinuous on every compact interval of  $[t_0, \infty)$ . Because  $\mathcal{T}(\Omega) \subset \Omega$ , the equiboundedness follows. Moreover, in view of (2.7),  $\mathcal{T}'(u)(t)$  is bounded for any  $u \in \Omega$ , which yields the equicontinuity of the elements in  $\mathcal{T}(\Omega)$ .

Now we prove the continuity of  $\mathcal{T}$  in  $\Omega$ . Let  $\{x_n\}$ ,  $n \in \mathbb{N}$ , be a sequence in  $\Omega$ , which uniformly converges to  $\overline{x} \in \Omega$  on every compact interval of  $[t_0, \infty)$ . Fixing  $T > t_0$ , in virtue of (2.23), the dominated convergence Lebesgue theorem gives

$$\lim_{n\to\infty} \int_{\sigma}^{T} \int_{\tau}^{\infty} r(s) \left( f\left(x_n(\overline{\varphi}(s))\right) \right) w(s,\tau) ds d\tau = \int_{\sigma}^{T} \int_{\tau}^{\infty} r(s) \left( f\left(x(\overline{\varphi}(s))\right) \right) w(s,\tau) ds d\tau. \quad (2.26)$$

Moreover,

$$\left| \int_{\sigma}^{\infty} \int_{\tau}^{\infty} r(s) \left( f\left(x_{n}(\overline{\varphi}(s))\right) - f\left(\overline{x}(\overline{\varphi}(s))\right) \right) w(s,\tau) ds d\tau \right|$$

$$\leq \left| \int_{\sigma}^{T} \int_{\tau}^{\infty} r(s) \left( f\left(x_{n}(\overline{\varphi}(s))\right) - f\left(\overline{x}(\overline{\varphi}(s))\right) \right) w(s,\tau) ds d\tau \right|$$

$$+ \int_{T}^{\infty} \int_{\tau}^{\infty} |r(s)| \left( f\left(x_{n}(\overline{\varphi}(s))\right) + f\left(\overline{x}(\overline{\varphi}(s))\right) \right) |w(s,\tau)| ds d\tau.$$

$$(2.27)$$

In view of (2.19), we have

$$\int_{T}^{\infty} \int_{\tau}^{\infty} |r(s)| \left( f\left(x_{n}(\overline{\varphi}(s))\right) + f\left(\overline{x}(\overline{\varphi}(s))\right) \right) |w(s,\tau)| ds d\tau \leq 2M \int_{T}^{\infty} |r(s)| F\left(\overline{\varphi}(s)\right) ds. \tag{2.28}$$

Thus, choosing T sufficiently large, we get from (2.27)

$$\lim_{n \to \infty} \int_{\sigma}^{\infty} \int_{\tau}^{\infty} r(s) f(x_n(\overline{\varphi}(s))) w(s,\tau) ds d\tau = \int_{\sigma}^{\infty} \int_{\tau}^{\infty} r(s) f(\overline{x}(\overline{\varphi}(s))) w(s,\tau) ds d\tau, \quad (2.29)$$

and so the continuity of  $\mathcal{T}$  in  $\Omega$  follows. By the Tychonoff fixed point theorem, the operator  $\mathcal{T}$  has a fixed point x, which is an unbounded solution of (1.1) satisfying (2.1).

*Remark* 2.2. With minor modifications, Theorem 2.1 gives also the existence of eventually negative unbounded AL-solutions. The details are omitted.

Remark 2.3. When  $\varphi(t) \equiv t$ , Theorem 2.1 is related with Theorem 1 in [15], from which the existence of unbounded AL-solutions of (1.1) can be obtained under stronger assumptions. A comparison between Theorem 1 in [15] and Theorem 2.1 is given in Section 4.

Our next result gives a necessary condition for the existence of unbounded solutions x of (1.1) satisfying for large t and some  $\alpha$  and  $\beta$ 

$$0 < \alpha \le x'(t) \le \beta. \tag{2.30}$$

**Theorem 2.4.** *Assume either*  $r(t) \ge 0$  *or*  $r(t) \le 0$ .

Equation (1.1) does not have eventually positive solutions x satisfying (2.30) for large t and some  $\alpha$  and  $\beta$  provided

$$\int_{0}^{\infty} |r(t)| \overline{F}(\varphi(t)) dt = \infty, \tag{2.31}$$

where for u > 0

$$\overline{F}(u) = \min \left\{ f(v) : \frac{\alpha}{2} u \le v \le 2\beta u \right\}. \tag{2.32}$$

*Proof.* Assume  $r(t) \ge 0$ , and let x be an eventually positive solution of (1.1) satisfying (2.30). Then, there exists  $\tau$  such that

$$\frac{\alpha}{2}t \le x(t) \le 2t\beta \quad \text{for } t \ge \tau.$$
 (2.33)

Consequently, in view of (2.31), we have

$$\lim_{t \to \infty} \int_{\tau}^{t} r(s) f(x(\varphi(s))) ds = \infty.$$
 (2.34)

Thus, integrating (1.1), we get

$$\lim_{t \to \infty} \left( x'''(t) + \int_{\tau}^{t} q(s)x''(s)ds \right) = -\infty. \tag{2.35}$$

Furthermore,

$$\left| \int_{\tau}^{t} q(s)x''(s)ds \right| = \left| q(t)x'(t) - q(\tau)x'(\tau) - \int_{\tau}^{t} q'(s)x'(s)ds \right|$$

$$\leq 2\beta Q + \beta \int_{\tau}^{\infty} |q'(s)|ds < \infty,$$
(2.36)

where  $Q = \sup_{s \ge 0} q(s)$ . Hence  $\lim_{t \to \infty} x'''(t) = -\infty$ , which gives a contradiction with the boundedness of x'. Finally, if  $r(t) \le 0$ , the argument is similar and the details are left to the reader.

#### 3. Bounded Solutions

In this section we study the existence of bounded AL-solutions of (1.1). The following holds.

Theorem 3.1. If

$$\int_0^\infty |r(t)|t\,dt < \infty,\tag{3.1}$$

then, for any  $c \in \mathbb{R} \setminus \{0\}$ , there exists a solution x of (1.1) satisfying

$$\lim_{t \to \infty} x(t) = c, \qquad \lim_{t \to \infty} x^{(i)}(t) = 0, \quad i = 1, 2.$$
 (3.2)

*Proof.* Without loss of generality, we prove the existence of solutions of (1.1) satisfying (3.2) for c = 1.

We proceed by a similar way to that in the proof of Theorem 2.1, and we sketch the proof.

Let M be the constant given in (2.5), and let

$$K = \max\left\{f(u): \frac{1}{2} \le u \le \frac{3}{2}\right\}, \qquad L_1 = \frac{2K(2M+1)}{q_0}. \tag{3.3}$$

Choose  $t_0 \ge 0$  large so that

$$\int_{t_0}^{\infty} t |r(t)| dt \le \frac{1}{2L_1}, \qquad \frac{1}{q_0} \int_{t_0}^{\infty} |q'(s)| ds \le \frac{1}{2}, \tag{3.4}$$

and define  $\overline{\varphi}$  as in (2.6). Denote by  $C[t_0, \infty)$  the Fréchet space of all continuous functions on  $[t_0, \infty)$ , endowed with the topology of uniform convergence on compact subintervals of  $[t_0, \infty)$ , and consider the set  $\Omega \subset C[t_0, \infty)$  given by

$$\Omega = \left\{ x \in C[t_0, \infty) : \frac{1}{2} \le x(t) \le \frac{3}{2} \right\}. \tag{3.5}$$

Let  $T > t_0$ , and, for any  $x \in \Omega$ , consider again the function  $\gamma$  given in (2.10). Reasoning as in the proof of Theorem 2.1, with minor changes, we obtain

$$\left| \int_{t}^{\infty} \int_{\tau}^{\infty} r(s) f(x(\overline{\varphi}(s))) w(s,\tau) ds d\tau \right| \le L_{1} \int_{t}^{\infty} |r(s)| ds. \tag{3.6}$$

Hence, in virtue of (3.1), the operator  $\mathcal{H}: \Omega \to \Omega$  given by

$$\mathcal{H}(x)(t) = 1 + \int_{t}^{\infty} \int_{\sigma}^{\infty} \int_{\tau}^{\infty} r(s) f\left(x(\overline{\varphi}(s))\right) w(s, \tau) ds d\tau d\sigma \tag{3.7}$$

is well defined and  $\lim_{t\to\infty} \mathcal{H}(x)(t) = 1$ . In view of (3.6), we get

$$\left| \mathcal{H}'(x)(t) \right| \le L_1 \int_t^\infty |r(s)| ds. \tag{3.8}$$

A similar estimation holds for  $|\mathcal{H}''(x)|$ . Thus,  $\lim_{t\to\infty}\mathcal{H}^{(i)}(x)(t)=0$ , i=1,2. In view of (3.4), from (3.8), we obtain

$$|\mathcal{H}(x)(t) - 1| \le L_1 \int_t^\infty s|r(s)|ds \le \frac{1}{2},\tag{3.9}$$

that is,  $\mathcal{H}(\Omega) \subset \Omega$ . Moreover, a standard calculation gives

$$\mathcal{H}^{(4)}(x)(t) = -q(t)\mathcal{H}^{(2)}(x)(t) - r(t)f(x(\overline{\varphi}(s))), \tag{3.10}$$

and so any fixed point of  $\mathcal{A}$  is, for large t, a solution of (1.1). Proceeding by a similar way to that in the proof of Theorem 2.1, we obtain that  $\mathcal{A}(\Omega)$  is relatively compact.

Now we prove the continuity of  $\mathcal{H}$  in  $\Omega$ . Let  $\{x_n\}$ ,  $n \in \mathbb{N}$ , be a sequence in  $\Omega$ , which uniformly converges to  $\overline{x} \in \Omega$  on every compact interval of  $[t_0, \infty)$ . Since

$$\left| \int_{\tau}^{\infty} r(s) \left( f\left( x_n(\overline{\varphi}(s)) \right) \right) w(s, \tau) ds \right| \le KM \int_{\tau}^{\infty} |r(s)| ds, \tag{3.11}$$

in virtue of (3.1), the dominated convergence Lebesgue theorem gives

$$\lim_{n\to\infty} \int_{\sigma}^{\infty} \int_{\tau}^{\infty} r(s) \left( f\left(x_n(\overline{\varphi}(s))\right) \right) w(s,\tau) ds d\tau = \int_{\sigma}^{\infty} \int_{\tau}^{\infty} r(s) \left( f\left(x(\overline{\varphi}(s))\right) \right) w(s,\tau) ds d\tau. \quad (3.12)$$

Moreover, fixing  $T > t_0$ , we have

$$\left| \int_{t}^{\infty} \int_{\sigma}^{\infty} \int_{\tau}^{\infty} r(s) \left( f\left(x_{n}(\overline{\varphi}(s))\right) - f\left(\overline{x}(\overline{\varphi}(s))\right) \right) w(s,\tau) ds \, d\tau \, d\sigma \right|$$

$$\leq \left| \int_{t}^{T} \int_{\sigma}^{\infty} \int_{\tau}^{\infty} r(s) \left( f\left(x_{n}(\overline{\varphi}(s))\right) - f\left(\overline{x}(\overline{\varphi}(s))\right) \right) w(s,\tau) ds \, d\tau \, d\sigma \right|$$

$$+ \int_{T}^{\infty} \int_{\sigma}^{\infty} \int_{\tau}^{\infty} |r(s)| \left( f\left(x_{n}(\overline{\varphi}(s))\right) + f\left(\overline{x}(\overline{\varphi}(s))\right) \right) |w(s,\tau)| ds \, d\tau \, d\sigma.$$

$$(3.13)$$

In view of (3.9), we have

$$\int_{T}^{\infty} \int_{\sigma}^{\infty} \int_{\tau}^{\infty} |r(s)| \left( f\left(x_{n}(\overline{\varphi}(s)\right)\right) + f\left(\overline{x}(\overline{\varphi}(s))\right) \right) |w(s,\tau)| ds \, d\tau \, d\sigma \leq 2L_{1} \int_{T}^{\infty} s|r(s)| ds, \quad (3.14)$$

and so, choosing T sufficiently large, from (3.13) we obtain the continuity of  $\mathcal{A}$  in  $\Omega$ . Hence, by the Tychonoff fixed point theorem, the operator  $\mathcal{A}$  has a fixed point x, which is a bounded solution of (1.1) satisfying (3.2).

Remark 3.2. When n = 4, Theorem 3.1 extends to equations with deviating argument of a similar result stated in [15] for (1.7). Observe that our approach used here is completely different from that in [15].

The next result shows that, under additional assumptions, condition (3.1) can be also necessary for the existence of bounded AL-solutions of (1.1).

#### **Theorem 3.3.** Assume either

$$r(t) \ge 0$$
,  $q''(t) \ge 0$  for large  $t$  (3.15)

or

$$r(t) \le 0$$
,  $q''(t) \le 0$  for large  $t$ . (3.16)

If

$$\int_0^\infty |r(t)|t\,dt = \infty,\tag{3.17}$$

then (1.1) does not have solutions x satisfying

$$0 < \alpha \le x(t) \le \beta, \tag{3.18}$$

for large t and some  $\alpha$  and  $\beta$ . Consequently, every bounded solution x of (1.1) satisfies

$$\lim_{t \to \infty} \inf |x(t)| = 0.$$
(3.19)

The following lemmas are needed for proving Theorem 3.3.

**Lemma 3.4.** Assume  $q''(t) \ge 0$  for  $t \ge T \ge 0$ , and let x be a solution of (1.1) satisfying (3.18) for  $t \ge T$ . Then, there exist two constants  $M_1$ ,  $M_2$  such that for  $t \ge T$ 

$$-\int_{T}^{t} sq(s)x''(s)ds < tq'(t)x(t) - tq(t)x'(t) + M_{1},$$
(3.20)

$$\int_{T}^{t} (s - T)q'(s)x'(s)ds < M_{2}. \tag{3.21}$$

If  $q''(t) \le 0$  for  $t \ge T \ge 0$ , inequalities (3.20), (3.21) hold in the opposite order.

*Proof.* Suppose  $q''(t) \ge 0$  on  $[T, \infty)$ . We have

$$\int_{T}^{t} sq(s)x''(s)ds = tq(t)x'(t) - Tq(T)x'(T) - \int_{T}^{t} q(s)x'(s)ds - \int_{T}^{t} sq'(s)x'(s)ds.$$
 (3.22)

Since

$$\int_{T}^{t} sq'(s)x'(s)ds = tq'(t)x(t) - Tq'(T)x(T) - \int_{T}^{t} q'(s)x(s)ds - \int_{T}^{t} sq''(s)x(s)ds,$$

$$\int_{T}^{t} q(s)x'(s)ds = q(t)x(t) - q(T)x(T) - \int_{T}^{t} q'(s)x(s)ds,$$
(3.23)

from (3.22), we get

$$-\int_{T}^{t} sq(s)x''(s)ds = tq'(t)x(t) - tq(t)x'(t) + q(t)x(t)$$

$$-2\int_{T}^{t} q'(s)x(s)ds - \int_{T}^{t} sq''(s)x(s)ds + K_{1},$$
(3.24)

where  $K_1$  is a suitable constant. Since q, x are bounded,  $q''(t) \ge 0$ , in view of (1.1), inequality (3.20) follows.

Moreover, q' is nondecreasing for  $t \ge T$ . Because q is a positive bounded function, then  $q'(t) \le 0$  on  $[T, \infty)$ . Thus, inequality (3.21) follows integrating by parts and using (1.1). Finally, if  $q''(t) \le 0$  on  $[T, \infty)$ , the argument is similar.

**Lemma 3.5.** Let x be a solution of (1.1) satisfying (3.18) for large t. If

$$\int_{0}^{\infty} |r(t)|dt < \infty, \tag{3.25}$$

then x" is bounded. If, in addition,  $r(t) \ge 0$ ,  $q''(t) \ge 0$  for  $t \ge T \ge 0$  and (3.17) holds, then for large t

$$x'''(t) + q(t)x'(t) < q'(t)x(t).$$
(3.26)

If  $r(t) \le 0$ ,  $q''(t) \le 0$  for  $t \ge T \ge 0$ , inequality (3.26) holds in the opposite order.

*Proof.* Since  $\lim_{t\to\infty} \varphi(t) = \infty$ , there exists  $\tau$  such that for  $t \ge \tau$ 

$$0 < \alpha \le x(\varphi(t)) \le \beta. \tag{3.27}$$

Without loss of generality, let  $\tau = T$ . Thus,  $\inf_{t \ge T} f(x(\varphi(t))) > 0$ .

Let u and v be two linearly independent solutions of (1.3) with Wronskian d=1. By assumptions on q, all solutions of (1.3) and their derivatives are bounded. Thus, by the variation constant formula, there exist constants  $c_1$  and  $c_2$  such that

$$x''(t) = c_1 u(t) + c_2 v(t) - \int_T^t (u(s)v(t) - u(t)v(s))r(s)f(x(\varphi(s)))ds,$$
 (3.28)

and, in view of (3.25), x'' is bounded.

Let us prove (3.26), and suppose  $r(t) \ge 0$ ,  $q''(t) \ge 0$  on  $[T, \infty)$ . Multiplying (1.1) by t and integrating from T to t, we get

$$tx'''(t) - x''(t) + \int_{T}^{t} sq(s)x''(s)ds = Tx'''(T) - x''(T) - \int_{T}^{t} sr(s)f(x(\varphi(s)))ds, \tag{3.29}$$

or, in view of Lemma 3.4,

$$tx'''(t) \le x''(t) + tq'(t)x(t) - tq(t)x'(t) - \int_{T}^{t} sr(s)f(x(\varphi(s)))ds + K_{2}, \tag{3.30}$$

where  $K_2$  is a suitable constant. Since x'' is bounded and

$$\int_{T}^{t} sr(s) f(x(\varphi(s))) ds \ge \inf_{t \ge T} f(x(\varphi(t))) \int_{T}^{t} sr(s) ds, \tag{3.31}$$

from (3.17) and (3.30), we have

$$\lim_{t \to \infty} t \left( x'''(t) - q'(t)x(t) + q(t)x'(t) \right) = -\infty, \tag{3.32}$$

which gives the assertion. The case  $r(t) \le 0$ ,  $q''(t) \le 0$  on  $[T, \infty)$  can be treated in a similar way.

*Proof of Theorem 3.3.* Suppose  $r(t) \ge 0$ ,  $q''(t) \ge 0$  for  $t \ge T \ge 0$ . Without loss of generality, assume also that (3.27) holds for  $t \ge T$ . Define

$$v(t) = x''(t) + q(t)x(t), \tag{3.33}$$

$$z(t) = x'''(t) + q(t)x'(t) - \int_{T}^{t} q'(s)x'(s)ds.$$
 (3.34)

Then,  $z'(t) = -r(t) f(x(\varphi(t))) \le 0$  and

$$z(t) = z(T) - \int_{T}^{t} r(s) f(x(\varphi(s))) ds.$$
(3.35)

Since  $q'(t) \le 0$  for  $t \ge T$ , we have

$$v'(t) \le z(t) + \int_{T}^{t} q'(s)x'(s)ds = z(T) - \int_{T}^{t} r(s)f(x(\varphi(s)))ds + \int_{T}^{t} q'(s)x'(s)ds.$$
 (3.36)

Case I. Assume

$$\int_0^\infty r(t)dt = \infty. \tag{3.37}$$

Since for  $t \ge T$  we have  $q''(t) \ge 0$  and, as claimed,  $q'(t) \le 0$ , we get

$$\int_{T}^{t} q'(s)x'(s)ds = q'(t)x(t) - q'(T)x(T) - \int_{T}^{t} q''(s)x(s)ds \le -q'(T)x(T). \tag{3.38}$$

Thus, from (3.36), we obtain  $\lim_{t\to\infty}v'(t)=-\infty$ , that is, v is unbounded. Hence, in view of (3.33), we obtain a contradiction with the boundedness of x.

Case II. Now assume (3.17) and (3.25). In view of Lemma 3.5, without loss of generality, we can suppose that (3.26) holds for  $t \ge T$ . Then,

$$z(T) = x'''(T) + q(T)x'(T) < q'(T)x(T).$$
(3.39)

Hence, z(T) < 0. Integrating (3.36), we get

$$v(t) \le v(T) + z(T)(t - T) - \int_{T}^{t} (s - T)r(s)f(x(\varphi(s)))ds + \int_{T}^{t} (s - T)q'(s)x'(s)ds, \qquad (3.40)$$

and, in view of Lemma 3.4, we have

$$v(t) \le v(T) + z(T)(t - T) + M_2. \tag{3.41}$$

Thus,  $\lim_{t\to\infty} v(t) = -\infty$ , that is, as before, a contradiction. Finally, the case  $r(t) \le 0$ ,  $q''(t) \le 0$  for large t follows in a similar way.

#### 4. Applications

Here we present some applications of our results to a particular case of (1.1), namely, the equation

$$x^{(4)}(t) + q(t)x''(t) + r(t)|x(\varphi(t))|^{\lambda} \operatorname{sgn} x(\varphi(t)) = 0 \quad (\lambda > 0), \tag{4.1}$$

jointly with some suggestions for future research.

#### 4.1. Coexistence of Both Types of AL-Solutions

Applying Theorems 2.1–3.3 to this equation, we obtain the following.

**Corollary 4.1.** (a) Let  $r(t) \neq 0$  for large t. Equation (4.1) has unbounded AL-solutions if and only if

$$\int_0^\infty |r(t)| \varphi^{\lambda}(t) dt < \infty. \tag{4.2}$$

(b) Assume either (3.15) or (3.16). Equation (4.1) has bounded AL-solutions if and only if (3.1) holds.

Corollary 4.1 shows also that the deviating argument can produce a different situation concerning the unboundedness of solutions with respect to the corresponding equation without delay, as the following example illustrates.

Example 4.2. In view of Corollary 4.1(a), the equation

$$x^{(4)}(t) + q(t)x^{(2)}(t) + \frac{1}{(t+1)^2} \left| x(\sqrt{t}) \right|^{3/2} \operatorname{sgn} x(\sqrt{t}) = 0, \tag{4.3}$$

where q satisfies (i), has unbounded AL-solutions, while the corresponding ordinary equation

$$x^{(4)}(t) + q(t)x^{(2)}(t) + \frac{1}{(t+1)^2}|x(t)|^{3/2}\operatorname{sgn} x(t) = 0,$$
(4.4)

in view of Theorem 2.4, does not have unbounded AL-solutions. Moreover, if in addition q''(t) > 0 for large t, then from Corollary 4.1(b) (4.3) does not have bounded AL-solutions.

The following example shows that the opposite situation to the one described in Example 4.2 can occur.

Example 4.3. Consider the equation

$$x^{(4)}(t) + q(t)x^{(2)}(t) + \frac{1}{(t+1)^3}x(t^2) = 0, (4.5)$$

where q satisfies (i). From Theorem 3.1, (4.5) has bounded AL-solutions and the same occurs for the corresponding ordinary equation. Nevertheless, in view of Corollary 4.1(a), (4.5) has no unbounded AL-solutions.

Examples 4.2 and 4.3 illustrate also that the coexistence of both AL-solutions for (4.1) can fail. Sufficient conditions for the coexistence of these solutions immediately follow from Corollary 4.1.

#### **Corollary 4.4.** *Let* $r(t) \neq 0$ *for large* t.

(a) Assume for large t

$$\varphi(t) \ge t^{1/\lambda}.\tag{4.6}$$

If (4.1) has unbounded AL-solutions, then (4.1) also has AS bounded solutions.

(b) Assume for large t

$$\varphi(t) \le t^{1/\lambda}, \qquad \operatorname{sgn} r(t) = \operatorname{sgn} q''(t).$$
 (4.7)

*If* (4.1) *has bounded AL-solutions, then* (4.1) *also has unbounded AL-solutions. For the equation without deviating argument* 

$$x^{(4)}(t) + q(t)x''(t) + r(t)|x(t)|^{\lambda} \operatorname{sgn} x(t) = 0 \quad (\lambda > 0), \tag{4.8}$$

from Corollary 4.4 we get the following.

**Corollary 4.5.** *Let*  $r(t) \neq 0$  *for large* t.

- (a) Assume  $\lambda \geq 1$ . If (4.8) has unbounded AL-solutions, then (4.8) has also bounded AL-solutions.
- (b) Assume  $0 < \lambda \le 1$  and sgn  $r(t) = \operatorname{sgn} q''(t)$  for large t. If (4.8) has bounded AL-solutions, then (4.8) has also unbounded AL-solutions.

#### 4.2. Comparison with Some Results in [15]

As claimed, the existence of unbounded AL-solutions for (4.8) follows also from Theorem 1 in [15]. For n = 4 this result reads as follows.

Theorem A. If

$$\int_{0}^{\infty} |r(t)| t^{\lambda+1} dt < \infty, \tag{4.9}$$

then there exists a solution x of (4.8) such that

$$x^{(i)}(t) = t^{(i)} + \varepsilon_i(t), \quad i = 0, ..., 3,$$
 (4.10)

where  $\varepsilon_i$  are functions of bounded variation for large t and  $\lim_{t\to\infty} \varepsilon_i(t) = 0$ .

Therefore, when  $\varphi(t) \equiv t$ , Theorem 2.1 ensures the existence of unbounded AL-solutions of (4.8) under a weaker condition than (4.9), namely,

$$\int_0^\infty |r(t)| t^\lambda dt < \infty. \tag{4.11}$$

On the other hand, Theorem A gives an asymptotic formula for such solutions.

#### 4.3. An Open Problem

Equation (1.1) can admit also other types of nonoscillatory solutions, as the following examples show.

Example 4.6. Consider the equation

$$x^{(4)}(t) + x^{(2)}(t) - \frac{2t^2 + 4t + 26}{(t+1)^{7/2}} |x(t)|^{3/2} \operatorname{sgn} x(t) = 0.$$
 (4.12)

In virtue of Corollary 4.1(b), (4.12) has no bounded AL-solutions. Nevertheless, this equation admits nonoscillatory bounded solutions because  $x(t) = (1 + t)^{-1}$  is a solution of (4.12).

Example 4.7. Consider the equation

$$x^{(4)}(t) + x^{(2)}(t) + \frac{t^2 + 4t + 10}{(t+2)^4 (\log(t+2))^3} x^3(t) = 0.$$
(4.13)

Thus, (3.1) holds, while  $\int_0^\infty t^3 r(t) dt = \infty$ . Hence, in virtue of Corollary 4.1, (4.13) has bounded AL-solutions, but no unbounded AL-solutions. Nevertheless, this equation admits nonoscillatory unbounded solutions because  $x(t) = \log(t+2)$  is a solution of (4.13).

The existence of nonoscillatory solutions x satisfying either  $\lim_{t\to\infty} x(t) = 0$  or  $\lim |x(t)| = \infty$ ,  $\lim_{t\to\infty} x'(t) = 0$  will be a subject of our next research.

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