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Comment on ``Statistical efficiency of methods for computing free energy of hydration'' [J. Chem. Phys. 149, 144111 (2018)]

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Comment on “Statistical efficiency of methods for computing free energy of hydration” [J. Chem. Phys. 149, 144111 (2018)]

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In Ref. 1, Yildirim *et al.* compared the statistical efficiency for the computation of the hydration free energy (HFE) of a set of small organic molecules using nonequilibrium work (NEW) methods and the standard equilibrium free energy perturbation (FEP) approach with stratification (MBAR). Based on the analysis of the 34 rigid molecules set, the authors conclude that “the nonequilibrium methods tested here for the prediction of HFE have lower computational efficiency than the MBAR method.”

This conclusion is based on the comparison of a so-called “statistical efficiency” $\epsilon_i = 1/T_i\sigma_i^2$, where T_i and σ_i^2 are the total simulation time and the mean variance of method i , respectively. The authors state that the simulation time, T , “is used as a proxy for the amount of information in the simulation.” Their definition of ϵ is not, by any means, a statistical efficiency and depends on their particular choice of the simulation protocol. The statistical efficiency of an unbiased estimator for a given parameter, say, ΔG , is rigorously defined as the ratio of the inverse of the Fischer information (or minimum variance) and the sample variance for ΔG (e.g., computed with bootstrap with resampling), i.e.,

$$\epsilon = \frac{1/I(\Delta G)}{\sigma^2}, \quad (1)$$

where

$$I(\Delta G) = -E\left[\frac{\partial^2 \log(f(X, \Delta G))}{\partial \Delta G^2} \middle| \Delta G\right] \quad (2)$$

is the Fischer information and $f(X, \Delta G)$ is the distribution of the estimator ΔG . The statistical efficiency is hence a number such that $0 \leq \epsilon \leq 1$. The inverse of the Fischer information for a fixed number of simulations, in the case of BAR, is known and is given by

Eq. (10) of Ref. 2. If the authors had used the correct Fischer information instead of their “proxy,” they would have found a statistical efficiency close to one. This is so since the sample variance, σ^2 , obtained by computing many time ΔG using resampling with replacement, must be close to and *larger* than the theoretical variance $1/I(\Delta G)$ for a fixed number of studies given in Ref. 2. Hence, the true ϵ says nothing about the computational efficiency, except for the known fact³ that BAR is a statistically efficient estimator. Nonetheless, ϵ could be regarded as a legitimate measure of *computational efficiency*. However, in order to test the efficiency of the NEW method, it would have been interesting to assess, at fixed total simulation time T , the effect on ΔG and its variance of the duration of the nonequilibrium trajectories and of the switching protocol during the decoupling/recoupling of the solute. By the same token, the optimal protocol⁴ in terms of ϵ could have been determined for the FEP/MBAR simulation (number and schedule of λ windows). Unfortunately, this kind of analysis is missing in the Yildirim paper.

Concerning the FEP-based equilibrium simulations, MBAR⁵ is known to “perform similarly to BAR when the spacing between intermediate states is moderate and therefore only neighboring states have significant phase space overlap.” This sentence is taken *verbatim* from Ref. 7⁶ of the Yildirim paper. Besides, in Ref. 7, I have shown that the free energies computed with BAR or MBAR are virtually indistinguishable in a challenging equilibrium alchemical application, provided that the overlap between neighboring distribution is significant, which makes MBAR somewhat redundant in well-designed equilibrium alchemical simulations. Consequently, the variance in MBAR must be similar to that of BAR. The latter is basically given by the sum over all inner strata of terms such as those given by Eq. (10) of Ref. 2. However, in FEP with stratification, BAR

or MBAR can give a reliable estimate if and only if the sampling is adequate in each stratum of the alchemical coordinate. In Refs. 6, 8, and 9, it has been authoritatively pointed out that “sampling (in alchemical calculations with stratification) remains a critical issue as the solute size and flexibility grow and as the solvent dynamics or environment becomes heterogeneous, for example, for solvation free energies in octanol.”⁶

In Ref. 1, the sampling issue has been analyzed only for the end states, which allows us to compute reliably the variance only for the NE approach, where the inner states are crossed at fast speed and just the final (end states) NE work distributions matter for producing the estimate ΔG . The authors appear to be somehow aware of this fact, since they did perform “three replicates of the same simulation” to get a “more realistic” uncertainty for the equilibrium MBAR alchemical computation. Even given for granted that three trials are enough to get a “realistic” value of the modified signed-digit (MSD) for each of the 34 solutes in the set, it remains unclear why the authors failed to include the cost of these replicates (100 ns each) in the total simulation time for MBAR, lowering their efficiency measure to 3.9, i.e., in line with the unidirectional NEW Jarzynski approaches and less than that of NEW/BAR.

Finally, I do believe that, for the sound reasons explained above, while the Yildirim *et al.* paper is certainly technically valid and useful, their conclusions are not sufficiently motivated and that a much more thorough and challenging analysis is needed in order to rigorously compare the efficiency of NEW and equilibrium stratification techniques, especially with regard to the issue of the adequacy of sampling. In this respect, it should be taken into account that in NEW, adequate sampling is required only at the end states and that uncertainty on ΔG is essentially a function of the variance (or dissipation) of the work distributions.⁹ In FEP with stratification,

adequate sampling must be checked on each λ stratum and the convergence rate is not guaranteed, by any means, to be the same on each of the strata.¹⁰

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