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Original Citation:

A bridging mechanism in the homogenization of brittle composites with soft inclusions / Barchiesi, Marco; Lazzaroni, Giuliano; Zeppieri, Caterina Ida. - In: SIAM JOURNAL ON MATHEMATICAL ANALYSIS. - ISSN 0036-1410. - STAMPA. - 48:(2016), pp. 1178-1209. [10.1137/15M1007343]

Availability:

This version is available at: 2158/1149058 since: 2021-03-18T12:41:43Z

Published version:

DOI: 10.1137/15M1007343

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A BRIDGING MECHANISM IN THE HOMOGENISATION OF BRITTLE COMPOSITES WITH SOFT INCLUSIONS

MARCO BARCHIESI, GIULIANO LAZZARONI, AND CATERINA IDA ZEPPIERI

Abstract. We provide a homogenisation result for the energy-functional associated with a purely brittle composite whose microstructure is characterised by soft periodic inclusions embedded in a stiffer matrix. We show that the two constituents as above can be suitably arranged on a microscopic scale ε to obtain, in the limit as ε tends to zero, a homogeneous macroscopic energy-functional explicitly depending on the opening of the crack.

Keywords: Γ -convergence, multiscale analysis, homogenisation, free-discontinuity problems, fracture mechanics.

2010 Mathematics Subject Class. 49J45, 74E30, 74G65, 74Q05, 74R99.

1. INTRODUCTION

Composites are materials that show heterogeneities on length scales that are much larger than the atomic scale but which are essentially homogeneous at macroscopic length scales. Engineered composites are widely used to improve structural performances; indeed with an appropriate choice of the microstructure they may efficiently combine the attributes of their constituents. For this reason, in recent decades there has been an ever increasing interest in the homogenisation theory of composites (and the consequent derivation of effective models) both in the mathematical and in the engineering communities.

The object of the present paper is the homogenisation of a brittle composite whose microstructure consists of soft (or compliant) inclusions periodically arranged in a stiffer matrix.

In the setting of linearised elasticity and antiplane shear, according to Griffith's theory, the microscopic energy corresponding to a displacement $u: \Omega \rightarrow \mathbb{R}$ (where $\Omega \subset \mathbb{R}^2$ is open, bounded, and represents the cross section of a cylindrical body) is given by

$$F_\varepsilon(u) = \int_{\Omega \cap \varepsilon P} |\nabla u|^2 dx + \delta_\varepsilon \int_{\Omega \setminus \varepsilon P} |\nabla u|^2 dx + \mathcal{H}^1(S_u),$$

where $\varepsilon > 0$ is the length-scale of the microstructure, εP is the ε -scaled copy of a connected, Q -periodic, open set $P \subset \mathbb{R}^2$, with $Q = (-1/2, 1/2)^2$, and $\delta_\varepsilon \rightarrow 0^+$ is the elastic modulus of the soft material. Thus, $\Omega \cap \varepsilon P$ and $\Omega \setminus \varepsilon P$ represent, respectively, the stiff and the compliant constituent of the brittle composite Ω (see Figure 1); moreover these two constituents are characterised by the same toughness, which here is normalised to one. The energy F_ε is given by the sum of three terms: two bulk energy contributions relative to the two constituents of Ω , and a surface energy contribution accounting for the energy needed to open the crack. The latter is identified with the one-dimensional discontinuity set S_u of a displacement u assumed to belong to $SBV^2(\Omega)$, the space of special functions of bounded variation for which the microscopic energy F_ε is finite for every fixed ε .

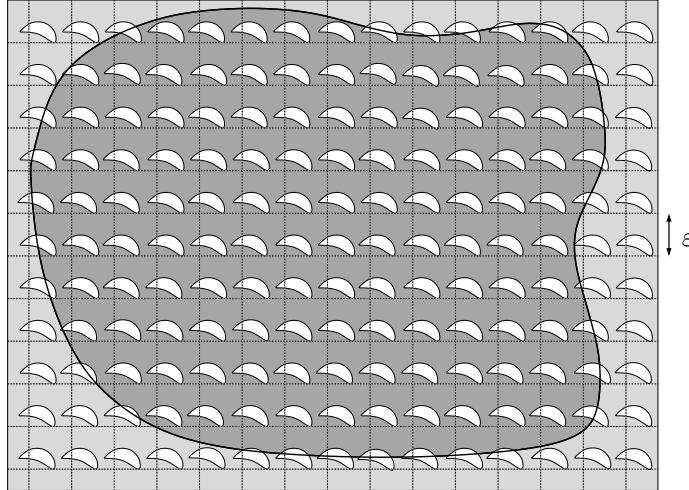


FIGURE 1. In dark grey the stiff matrix $\Omega \cap \varepsilon P$.

Using the tool of Γ -convergence [9, 17], in this paper we show that the limit behaviour of the functionals F_ε is not, in general, a simple superposition of the corresponding limit behaviour of the bulk and surface energy terms. More precisely, we exhibit an elementary micro-geometry and an elastic modulus δ_ε for which optimal sequences of displacements may depend on the interplay between bulk and surface energy, thus favouring the presence of high gradients or discontinuities of u inside the soft inclusions $\Omega \setminus \varepsilon P$. We show on a concrete example that even though the surface term in F_ε does not depend on the opening of the crack, this dependence explicitly appears in the limit as ε tends to zero.

Specifically, appealing to the localisation method of Γ -convergence and to the integral representation in SBV [7], we prove that (up to subsequences) the functionals F_ε Γ -converge to a homogenised functional F of the form

$$F(u) = \int_{\Omega} f_0(\nabla u) dx + \int_{S_u} g_0([u], \nu_u) d\mathcal{H}^1 \quad \text{for all } u \in SBV^2(\Omega), \quad (1.1)$$

where $[u]$ denotes the opening of the crack and ν_u its orientation, while the energy densities f_0 and g_0 satisfy

$$c_1 |\xi|^2 \leq f_0(\xi) \leq |\xi|^2 \quad \text{and} \quad c_2 \leq g_0(t, \nu) \leq 1,$$

for every $\xi \in \mathbb{R}^2$, $t \neq 0$, and $\nu \in \mathbb{S}^1$, and for some positive constants c_1, c_2 . Moreover, f_0 depends only on P (hence in particular not on δ_ε) and it can be shown that $t \mapsto g_0(t, \nu)$ is nondecreasing and left-continuous for $t > 0$, and satisfies the symmetry condition $g_0(-t, -\nu) = g_0(t, \nu)$.

The above Γ -convergence analysis strongly relies on some recent results for free-discontinuity problems in perforated domains [22, 14, 5]. In fact, the asymptotic analysis carried out in [22, 14, 5] allows us to bound from below the functional Γ -lim inf F_ε with the Mumford-Shah functional, thus obtaining the equi-coerciveness of F_ε . (We now overlook the fact that the domain of F is larger than $SBV^2(\Omega)$ and refer to Theorem 1 for a precise statement.)

The main result of this paper is Theorem 2 in which we specialise the functionals F_ε choosing the two-dimensional micro-geometry given by

$$P = \mathbb{R}^2 \setminus \bigcup_{i \in \mathbb{Z}^2} \left(\frac{1}{4} \overline{Q} + i \right)$$

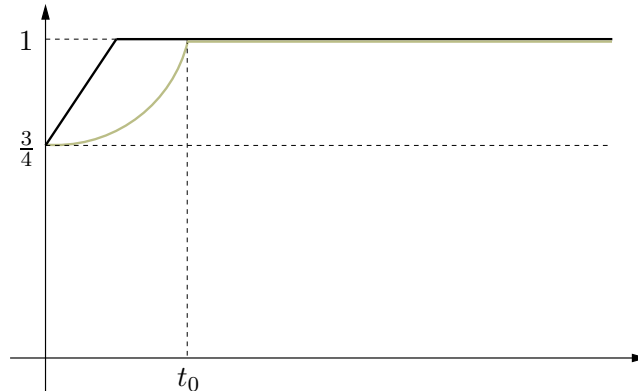


FIGURE 2. The surface energy density $g_0(\cdot, e_i)$ is bounded from above by $\min\{3/4 + \sqrt{2}t, 1\}$ and from below by $\min\{3/4 + c_3t^2, 1\}$. In particular, $g_0(0^+, e_i) = 3/4$ while $g_0(t, e_i) = 1$ for t larger than a threshold t_0 .

and setting $\delta_\varepsilon = \varepsilon$. We prove that for such a choice the following estimate holds true:

$$\min\left\{\frac{3}{4} + c_3t^2, 1\right\} \leq g_0(t, e_i) \leq \min\left\{\frac{3}{4} + \sqrt{2}t, 1\right\} \quad \text{for } t > 0 \text{ and } i = 1, 2, \quad (1.2)$$

where g_0 is as in (1.1). We mention here that the results which for modelling reasons have been described so far only in dimension two, can be suitably extended to the case of n -dimensional micro-geometries (and scalar displacement).

The estimate from below in (1.2) allows us to deduce that the surface energy density g_0 depends on $[u]$ in a nontrivial way. Moreover, the combination of the lower and upper bounds in (1.2) implies that $g_0(\cdot, e_i)$ is constant for every t larger than a positive threshold t_0 (see Figure 2). For these reasons, our result can be also interpreted as a possible mesoscopic justification of cohesive zone energies; the latter being characterised by a surface energy density whose dependence on the crack-opening is nonconstant and nondecreasing for small openings, and bounded for large ones. Cohesive zone models have been introduced by Barenblatt in [6] and are widely employed in fracture mechanics since they provide a more accurate description of the process of crack growth if compared with those based on Griffith's criterion. Indeed, explicitly depending on the crack-opening $[u]$, they take into account that fracture is a gradual process due to the fact that atomic bonds stretch before breaking. In a discrete-to-continuum setting, cohesive surface energies have been derived by means of Γ -convergence in [10] starting from one-dimensional discrete systems and in [12] by mixing quadratic and defected springs. Moreover, cohesive-type models have been obtained in e.g. [15, 18] via Ambrosio-Tortorelli approximation and in [20] as limits of coupled elastoplastic-damage models. Further, in [23] it has been proved that cohesive energies converge to the Mumford-Shah functional if the size of the body tends to infinity.

The analysis carried out in the present paper is quite different: we show that a nontrivial dependence of the limit surface energy on the opening of the crack may arise from the interplay, on a mesoscopic scale, between two different brittle constituents, even though the micro-geometry is rather simple. We shall also mention that our limit model has a positive activation threshold $g_0(0^+, e_i)$. From a mechanical point of view it would be desirable to derive a cohesive model without activation threshold (or equivalently a limit surface density with $g_0(0^+, \nu) = 0$ for every $\nu \in \mathbb{S}^1$) starting from an energy of brittle type; however, this seems to be a challenging question at the moment.

The presence of an activation threshold can be also found in a problem investigated in [4, 21], where a limit energy depending on the crack-opening is obtained homogenising a composite made

of a brittle constituent reinforced by an unbreakable structure; *i.e.*, by a structure with infinite toughness. It is worth to point out that in [4, 21] the limit surface energy density grows linearly in the crack-opening, thus showing an unphysical behaviour. Motivated by this observation, in the last section of this paper we briefly discuss a microscopic model similar to that analysed in [4, 21]. We show that a cohesive model with bounded surface energy density can be obtained if the unbreakable structure is replaced by periodically distributed brittle fibres with toughness tending to infinity as the microscopic scale tends to zero.

Both in the presence of soft inclusions and of tough fibres, the general idea to prove the bound from below on the limit surface energy density is to approximate a straight crack, or in other words a limit displacement of type $u_t(x) = t \chi_{\{x: x_2 > 0\}}(x)$, with functions u_ε suitably combining jumps and stretchings with high gradients in the different constituents of the composite (see Figures 5 and 10). Loosely speaking, at a microscopic level we observe (soft or tough) regions that stretch without breaking, thus acting as “bridges” between the two opposite sides of the macroscopic crack. This microscopic phenomenon is known in the mechanical literature as “bridging mechanism” and is experimentally observed *e.g.* in the fracture of fibre-reinforced plastics or ceramics and in the crazing of polymers [3, Chapter 6]: the crack propagation is preceded by the nucleation of micro-cracks and micro-voids, with some “bridging elements” that contribute in transferring stresses between the crack’s faces, thus adding resistance to large crack-opening and to further growth.

2. SETTING OF THE PROBLEM AND STATEMENT OF THE MAIN RESULT

In this section we recall some definitions and introduce a few notation we employ in the paper. For the sake of generality we now work in dimension $n \geq 2$.

For the general theory of *special functions of bounded variation* we refer to [2] (see also [8]).

Let U be an open bounded subset of \mathbb{R}^n . The space of special functions of bounded variation on U is denoted by $SBV(U)$. For every $u \in SBV(U)$, ∇u denotes the *approximate gradient* of u , S_u the *approximate discontinuity set* of u , and ν_u the *generalised normal* to S_u , which is defined up to the sign. If u^+ and u^- are the traces of u on the sides of S_u determined by ν_u and $-\nu_u$, respectively, the difference $u^+ - u^-$ is called the *jump* of u , and is denoted by $[u]$. Note that, with this convention, if we reverse the orientation of ν_u , we change the sign of $[u]$. It turns out that $[u] \in L^1(S_u; \mathcal{H}^{n-1})$.

We consider the vector subspace of $SBV(U)$

$$SBV^2(U) := \{u \in SBV(U) : \nabla u \in L^2(U; \mathbb{R}^n) \text{ and } \mathcal{H}^{n-1}(S_u) < +\infty\}.$$

We consider also the larger space of *generalised special functions of bounded variation* on U , $GSBV(U)$, which is made of all functions $u \in L^1(U)$ whose truncations $u^m := (u \wedge m) \vee (-m)$ belong to $SBV(U)$ for every $m \in \mathbb{N}$. By analogy with the case of SBV functions, we say that $u \in GSBV^2(U)$ if $u \in GSBV(U)$, $\nabla u \in L^2(U; \mathbb{R}^n)$ and $\mathcal{H}^{n-1}(S_u) < +\infty$.

For $r > 0$ we denote by Q_r the n -dimensional cube with side-length r , centred at the origin; *i.e.*, $Q_r := (-r/2, r/2)^n$; while we simply write Q in place of Q_1 .

The canonical basis in \mathbb{R}^n is denoted by $\{e_1, \dots, e_n\}$.

Unless otherwise stated, in what follows the Γ -convergence of functionals is always understood with respect to the strong L^1 -topology,

Let Ω be an open bounded subset of \mathbb{R}^n and let $\delta_\varepsilon > 0$ be such that $\delta_\varepsilon \rightarrow 0$ as $\varepsilon \rightarrow 0$. We introduce the functionals $F_\varepsilon : L^1(\Omega) \rightarrow [0, +\infty]$ defined as

$$F_\varepsilon(u) := \begin{cases} \int_{\Omega \cap \varepsilon P} |\nabla u|^2 dx + \delta_\varepsilon \int_{\Omega \setminus \varepsilon P} |\nabla u|^2 dx + \mathcal{H}^{n-1}(S_u) & \text{if } u \in SBV^2(\Omega), \\ +\infty & \text{otherwise in } L^1(\Omega), \end{cases}$$

where $P \subset \mathbb{R}^n$ is open, connected, and Q -periodic. We recall that in our model Ω represents the reference configuration of a periodic brittle composite made of two constituents having different elastic properties. More precisely, the elastic modulus of the constituent located in $\Omega \setminus \varepsilon P$ is represented by the vanishing sequence δ_ε . For this reason, in what follows, $\Omega \setminus \varepsilon P$ is referred to as the *soft inclusions*. In order to keep our analysis as simple as possible, we assume that the two constituents have the same toughness, here normalised to one. However, straightforward computations show that analogous results also hold when the two constituents have different toughnesses independent of ε . For the case where the toughness of the soft inclusions scales as ε^β with $\beta > 0$, we refer to Remark 4 below.

We also consider the functionals $\hat{F}_\varepsilon : L^1(\Omega) \rightarrow [0, +\infty]$ given by

$$\hat{F}_\varepsilon(u) := \begin{cases} \int_{\Omega \cap \varepsilon P} |\nabla u|^2 dx + \mathcal{H}^{n-1}(S_u \cap \Omega \cap \varepsilon P) & \text{if } u|_{\Omega \cap \varepsilon P} \in SBV^2(\Omega \cap \varepsilon P), \\ +\infty & \text{otherwise in } L^1(\Omega). \end{cases}$$

Notice that in this case $\Omega \setminus \varepsilon P$ represents a so-called *perforation*. The asymptotic behaviour of \hat{F}_ε has been recently studied in [5, Theorem 4] (see also [14, 22]). Specifically, \hat{F}_ε Γ -converges with respect to the strong L^1 -topology to

$$\hat{F}(u) := \begin{cases} \int_{\Omega} \hat{f}(\nabla u) dx + \int_{S_u} \hat{g}(\nu_u) d\mathcal{H}^{n-1} & \text{if } u \in GSBV^2(\Omega), \\ +\infty & \text{otherwise in } L^1(\Omega), \end{cases} \quad (2.1)$$

where $\hat{f} : \mathbb{R}^n \rightarrow [0, +\infty)$ and $\hat{g} : \mathbb{S}^{n-1} \rightarrow [0, +\infty)$ satisfy

$$\begin{aligned} c_1 |\xi|^2 &\leq \hat{f}(\xi) \leq |\xi|^2 && \text{for every } \xi \in \mathbb{R}^n, \\ c_2 &\leq \hat{g}(\nu) \leq 1 && \text{for every } \nu \in \mathbb{S}^{n-1}, \end{aligned} \quad (2.2)$$

for some constants $c_1, c_2 > 0$ only depending on n and P . Moreover, \hat{f} is a quadratic form given by the following homogenisation formula:

$$\hat{f}(\xi) = \inf \left\{ \int_{Q \cap P} |\xi + \nabla w|^2 dx : w \in H_{\text{per}}^1(Q \cap P) \right\}, \quad (2.3)$$

where $H_{\text{per}}^1(Q \cap P) := \{u \in H_{\text{loc}}^1(P) : u \text{ is } Q\text{-periodic and } u|_{Q \cap P} \in H^1(Q \cap P)\}$. If P is Lipschitz, the infimum in (2.3) can be taken over $H_{\text{per}}^1(Q)$.

The convergence result as above is proved in [5] with respect to the strong L^2 -topology; a technical but standard argument (see e.g. [21, Theorem 3.4]) allows us to deduce that the Γ -convergence of \hat{F}_ε to \hat{F} takes place also with respect to the strong L^1 -topology.

The estimates in (2.2) and the fact that the domain of \hat{F} is $GSBV^2(\Omega)$ are direct consequences of the compactness result provided in [5, Theorem 1], while (2.3) follows from [5, formula (40)].

It is also convenient to consider the Mumford-Shah functional

$$MS(u) := \begin{cases} \int_{\Omega} |\nabla u|^2 dx + \mathcal{H}^{n-1}(S_u) & \text{if } u \in GSBV^2(\Omega), \\ +\infty & \text{otherwise in } L^1(\Omega). \end{cases}$$

We notice that by virtue of the compactness result in *GSBV* [2, Theorem 4.36] MS is $L^1(\Omega)$ -lower semicontinuous on $GSBV^2(\Omega)$.

We remark that

$$\hat{F}_\varepsilon \leq F_\varepsilon \leq MS \quad \text{on } SBV^2(\Omega), \quad (2.4)$$

while, in view of the bounds from below in (2.2), we get

$$\min\{c_1, c_2\} MS \leq \hat{F} \quad \text{on } L^1(\Omega). \quad (2.5)$$

The bounds (2.4) and (2.5) will be crucial in the proof of the homogenisation result below.

Theorem 1. *For every decreasing sequence of positive numbers converging to zero, there exists a subsequence (ε_k) such that (F_{ε_k}) Γ -converges to a functional $F: L^1(\Omega) \rightarrow [0, +\infty]$ of the form*

$$F(u) := \begin{cases} \int_{\Omega} f_0(\nabla u) \, dx + \int_{S_u} g_0([u], \nu_u) \, d\mathcal{H}^{n-1} & \text{if } u \in GSBV^2(\Omega), \\ +\infty & \text{otherwise in } L^1(\Omega), \end{cases} \quad (2.6)$$

where $f_0: \mathbb{R}^n \rightarrow [0, +\infty)$ coincides with the quadratic form \hat{f} defined in (2.3), and $g_0: \mathbb{R} \times \mathbb{S}^{n-1} \rightarrow [0, +\infty)$ is a Borel function satisfying the following properties:

- (i) for every $t \neq 0$ and $\nu \in \mathbb{S}^{n-1}$, $c_2 \leq g_0(t, \nu) \leq 1$, where $c_2 > 0$ is as in (2.2);
- (ii) for any $\nu \in \mathbb{S}^{n-1}$, $g_0(\cdot, \nu)$ is nondecreasing and left-continuous in $(0, +\infty)$ and satisfies the symmetry condition $g_0(-t, -\nu) = g_0(t, \nu)$.

The proof of Theorem 1 is contained in Section 3.

Remark 1. *The assumption that the matrix P is open and connected is needed to invoke the results in [5]. Specifically, the proof of Theorem 1 strongly relies on the fact that $\hat{F} = \Gamma\text{-lim } \hat{F}_\varepsilon$ as well as on the bound (2.5). The latter, in its turn, together with (2.4), allows us to deduce that the domain of F is $GSBV^2(\Omega)$.*

Remark 2. *The quadratic form f_0 does not depend on either the elastic modulus δ_ε or the surface term in F_ε .*

The homogenisation result Theorem 1 asserts that the Γ -limit F may depend on both $[u]$ and ν_u , which in our modelling represent the opening and the orientation of the crack, respectively. In Theorem 2 below we show that for $\delta_\varepsilon = \varepsilon$ we can find an elementary two-dimensional microgeometry εP (see Figure 3) leading to a homogenised functional F whose surface energy density g_0 actually depends on the crack-opening $[u]$. We notice that the choice $\delta_\varepsilon = \varepsilon$ provides the only meaningful scaling for the microgeometry under consideration, indeed other choices of δ_ε give no cohesive behaviour in the limit, see Remark 4 below.

For the sake of simplicity, we prove the dependence of g_0 on $[u]$ only when the normal to the crack ν is one of the coordinate vectors. Moreover, we state the result in the physically relevant case of dimension two; for higher dimension, see Remark 5 below.

The following theorem is the main result of this paper and is proved in Section 4.

Theorem 2. *Let $n = 2$; set $\delta_\varepsilon := \varepsilon$ and*

$$P := \mathbb{R}^2 \setminus \bigcup_{(i,j) \in \mathbb{Z}^2} \bar{Q}_{\frac{1}{4}} + (i, j). \quad (2.7)$$

Then for every $t > 0$ the energy density g_0 appearing in (2.6) satisfies the following growth conditions

$$\min\left\{\frac{3}{4} + c_3 t^2, 1\right\} \leq g_0(t, e_i) \leq \min\left\{\frac{3}{4} + \sqrt{2}t, 1\right\} \quad \text{for } i = 1, 2, \quad (2.8)$$

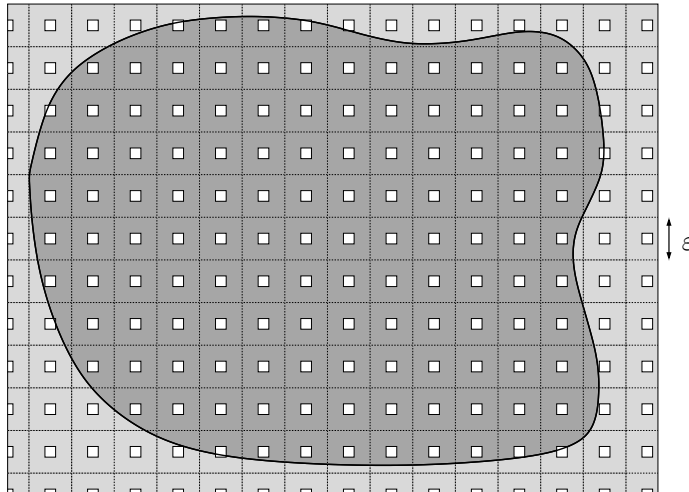


FIGURE 3. In dark grey the stiff matrix $\Omega \cap \varepsilon P$.

for some $c_3 > 0$. In particular, for $i = 1, 2$, we have that $g_0(0^+, e_i) = 3/4$ and $g_0(t, e_i) = 1$ for t larger than some $t_0 > 0$.

The following remarks are in order.

Remark 3. The homogenisation formula for \hat{g} provided in [5, Theorem 4] gives $\hat{g}(e_1) = \hat{g}(e_2) = 3/4$ when computed for the micro-geometry εP as in Theorem 2. Then, since

$$\hat{g}(\nu) \leq g_0(t, \nu) \quad \text{for every } t > 0 \text{ and } \nu \in \mathbb{S}^1,$$

in view of Theorem 2 we may deduce that, for $i = 1, 2$, $\hat{g}(e_i)$ is a sharp lower bound for $g_0(t, e_i)$ whenever $t \rightarrow 0^+$.

Remark 4. In this remark we briefly discuss a few further scalings. To this end let

$$\begin{aligned} F_\varepsilon^{\alpha, \beta}(u) &= \int_{\Omega \cap \varepsilon P} |\nabla u|^2 \, dx + \varepsilon^\alpha \int_{\Omega \setminus \varepsilon P} |\nabla u|^2 \, dx \\ &\quad + \mathcal{H}^1((\Omega \cap \varepsilon P) \cap S_u) + \varepsilon^\beta \mathcal{H}^1((\Omega \setminus \varepsilon P) \cap S_u), \end{aligned} \quad (2.9)$$

where $\alpha, \beta > 0$ and Ω, P are chosen as in Theorem 2. The free-discontinuity functionals in (2.9) are a particular case of those analysed by Braides and Solci in [13]. Then, thanks to [13, Theorem 3] it is possible to deduce that for every $\alpha, \beta > 0$ $(F_\varepsilon^{\alpha, \beta})$ Γ -converges to \hat{F} , where \hat{F} is as in (2.1). This implies, in particular, that a limit of cohesive type cannot be obtained if in our model also the toughness of the soft inclusions vanishes as $\varepsilon \rightarrow 0$.

Choosing in (2.9) $\alpha = 1$, we end up with the sequence

$$\begin{aligned} F_\varepsilon^\beta(u) &= \int_{\Omega \cap \varepsilon P} |\nabla u|^2 \, dx + \varepsilon \int_{\Omega \setminus \varepsilon P} |\nabla u|^2 \, dx \\ &\quad + \mathcal{H}^1((\Omega \cap \varepsilon P) \cap S_u) + \varepsilon^\beta \mathcal{H}^1((\Omega \setminus \varepsilon P) \cap S_u). \end{aligned}$$

Then, our functionals F_ε can be viewed as a “limit” case of F_ε^β when we let $\beta \rightarrow 0^+$. From [13, Theorem 3] we know that for every fixed $\beta > 0$ the Γ -limit of F_ε^β is \hat{F} . On the other hand, in Theorem 2 we prove that when $\beta = 0$ the Γ -limit of $F_\varepsilon = F_\varepsilon^{\beta=0}$ depends on $[u]$ in a nontrivial

way. This fact can be then interpreted as a sort of non-uniformity (or a lack of continuity) of the Γ -limit with respect to the exponent β .

We finally notice that for $\beta = 0$ and $\alpha > 1$ (which in our notation corresponds to choosing $\delta_\varepsilon \ll \varepsilon$) the Γ -limit is again \hat{F} , as it can be seen arguing as in Section 4. Moreover, a heuristic argument also shows that for $\alpha < 1$ (or equivalently $\delta_\varepsilon \gg \varepsilon$) the volume energy density in the Γ -limit does not change being always equal to \hat{f} , while the surface energy density does not depend on the crack-opening and is identically equal to one. Hence in our case $\delta_\varepsilon = \varepsilon$ is the only choice leading to a cohesive limit effect. Indeed, at this scaling, bulk and surface energies of the soft part are comparable at a mesoscopic level.

Remark 5. The choice $n = 2$ has been only made for the sake of clarity as it simplifies the exposition in the proof of Theorem 2. In fact, Theorem 2 can be implemented in higher dimensions still choosing $\delta_\varepsilon = \varepsilon$ and considering the micro-geometry corresponding to

$$P := \mathbb{R}^n \setminus \bigcup_{i \in \mathbb{Z}^n} \overline{Q}_{\frac{1}{4}} + i.$$

3. Γ -CONVERGENCE AND INTEGRAL REPRESENTATION

This section is devoted to the proof of the Γ -convergence of the functionals F_ε .

In order to prove Theorem 1 we use the well-known localisation method of Γ -convergence (for which we refer the reader to [17, Chapters 14-20]) in combination with the integral-representation result [7, Theorem 1]. To this end, we start introducing the localised functionals as below. We denote by $\mathcal{A}(\Omega)$ the class of all open subsets of Ω and for every pair $(u, U) \in L^1(\Omega) \times \mathcal{A}(\Omega)$ we set

$$F_\varepsilon(u, U) := \begin{cases} \int_{U \cap \varepsilon P} |\nabla u|^2 dx + \int_{U \setminus \varepsilon P} |\nabla u|^2 dx + \mathcal{H}^{n-1}(S_u \cap U) & \text{if } u \in SBV^2(U), \\ +\infty & \text{otherwise in } L^1(\Omega). \end{cases} \quad (3.1)$$

Moreover, it is also convenient to introduce the following notation:

$$MS(u, U) := \begin{cases} \int_U |\nabla u|^2 dx + \mathcal{H}^{n-1}(S_u \cap U) & \text{if } u \in GSBV^2(U), \\ +\infty & \text{otherwise in } L^1(\Omega). \end{cases}$$

We notice that the functionals $F_\varepsilon: L^1(\Omega) \times \mathcal{A}(\Omega) \rightarrow [0, +\infty]$ enjoy the following properties: for every $\varepsilon > 0$, $u \in L^1(\Omega)$, and $U \in \mathcal{A}(\Omega)$,

$$\begin{aligned} F_\varepsilon \text{ is increasing: } & F_\varepsilon(u, V) \leq F_\varepsilon(u, U) \text{ for every } V \in \mathcal{A}(\Omega): V \subset U, \\ F_\varepsilon \text{ is local: } & F_\varepsilon(u, U) = F_\varepsilon(v, U) \text{ for every } v \in L^1(\Omega): u = v \text{ } \mathcal{L}^n\text{-a.e. in } U, \\ F_\varepsilon \text{ decreases by truncation: } & F_\varepsilon((u \wedge M) \vee (-M), U) \leq F_\varepsilon(u, U) \text{ for } M > 0. \end{aligned} \quad (3.2)$$

Moreover F_ε is periodic; *i.e.*, for $y_\varepsilon := \lfloor \frac{\cdot}{\varepsilon} \rfloor \varepsilon$ (here the integer part is meant component-wise)

$$F_\varepsilon(u(\cdot - y_\varepsilon), U + y_\varepsilon) = F_\varepsilon(u, U) \text{ for all } y \in \mathbb{R}^n \text{ such that } U + y \subset\subset \Omega \text{ and sufficiently small } \varepsilon, \quad (3.3)$$

and invariant by translations in u ; *i.e.*,

$$F_\varepsilon(u + s, U) = F_\varepsilon(u, U) \text{ for every } s \in \mathbb{R}. \quad (3.4)$$

Now let (ε_k) be a vanishing sequence of strictly positive numbers and define the functionals $F', F'' : L^1(\Omega) \times \mathcal{A}(\Omega) \rightarrow [0, +\infty]$ as

$$F'(\cdot, U) := \Gamma\text{-lim inf}_{k \rightarrow +\infty} F_{\varepsilon_k}(\cdot, U) \quad \text{and} \quad F''(\cdot, U) := \Gamma\text{-lim sup}_{k \rightarrow +\infty} F_{\varepsilon_k}(\cdot, U),$$

for every $U \in \mathcal{A}(\Omega)$. The functionals F', F'' are lower semicontinuous [17, Proposition 6.8] and inherit the properties in (3.2). Specifically, they are increasing [17, Proposition 6.7], local [17, Proposition 16.15], and it is immediate to show that they decrease by truncation. Moreover, in view of [5, Theorem 4], (2.4), and (2.5) we may deduce that the domain of F' and F'' is $GSBV^2(\Omega)$ and that for every $(u, U) \in GSBV^2(\Omega) \times \mathcal{A}(\Omega)$ the following estimate holds:

$$\min\{c_1, c_2\} MS(u, U) \leq F'(u, U) \leq F''(u, U) \leq MS(u, U). \quad (3.5)$$

Now fix $\varepsilon > 0$ and $u \in GSBV^2(\Omega)$; clearly $F_\varepsilon(u, \cdot)$ is the restriction to $\mathcal{A}(\Omega)$ of a Radon measure, and therefore $F_\varepsilon(u, \cdot)$ is in particular inner regular. On the other hand, $F'(u, \cdot), F''(u, \cdot)$ are in general not inner regular, hence we also consider their inner regular envelope; *i.e.*, the two functionals $F'_-, F''_- : L^1(\Omega) \times \mathcal{A}(\Omega) \rightarrow [0, +\infty]$ defined as

$$F'_-(u, U) := \sup \{F'(u, V) : V \subset\subset U, V \in \mathcal{A}(\Omega)\}$$

and

$$F''_-(u, U) := \sup \{F''(u, V) : V \subset\subset U, V \in \mathcal{A}(\Omega)\}.$$

We notice that F'_- and F''_- are both increasing, lower semicontinuous [17, Remark 15.10], and local [17, Remark 15.25]. Moreover, invoking [17, Theorem 16.9] we can find a subsequence (ε_k) converging to zero such that the corresponding functionals F' and F'' satisfy

$$F'_- = F''_- =: F. \quad (3.6)$$

Next, appealing to [17, Proposition 18.6] we show that on $GSBV^2(\Omega)$ we actually have $F = F' = F''$. A preliminary result needed in this direction is the so-called fundamental estimate. The following variant of the fundamental estimate can be derived from the more general [11, Proposition 3.1]. For the reader's convenience we give here a simplified proof relative to our specific case.

Lemma 1 (Fundamental estimate). *For every $\eta > 0$ and for every $U', U'', V \in \mathcal{A}(\Omega)$, with $U' \subset\subset U''$, there exists a constant $M(\eta) > 0$ satisfying the following property: for every $\varepsilon > 0$, for every $u \in L^1(\Omega)$ with $u \in SBV^2(U'')$, and for every $v \in L^1(\Omega)$ with $v \in SBV^2(V)$, there exists a function $\varphi \in C_0^\infty(\Omega)$ with $\varphi = 1$ in a neighbourhood of U' , $\text{spt } \varphi \subset U''$ and $0 \leq \varphi \leq 1$ such that*

$$F_\varepsilon(\varphi u + (1 - \varphi)v, U' \cup V) \leq (1 + \eta) \left(F_\varepsilon(u, U'') + F_\varepsilon(u, V) \right) + M(\eta) \|u - v\|_{L^2(S)}^2,$$

with $S := (U'' \setminus U') \cap V$.

Proof. Let $\eta > 0$, $U', U'', V \in \mathcal{A}(\Omega)$ be fixed as in the statement and let φ be a function in $C_0^\infty(\Omega)$ with $0 \leq \varphi \leq 1$, $\text{spt } \varphi \subset U''$ and $\varphi = 1$ in a neighbourhood of U' .

Let u and v be two functions as in the statement and let $w := \varphi u + (1 - \varphi)v$; clearly, $w \in SBV^2(U' \cap V)$. Then,

$$F_\varepsilon(w, U' \cup V) = F_\varepsilon(u, U') + F_\varepsilon^*(v, V \setminus U'') + F_\varepsilon^*(w, (U'' \setminus U') \cap V), \quad (3.7)$$

where for fixed $u \in L^1(\Omega)$, $F_\varepsilon^*(u, \cdot)$ denotes the measure which extends $F_\varepsilon(u, \cdot)$ to the σ -algebra $\mathcal{B}(\Omega)$ of Borel subsets of Ω in the usual way; *i.e.*,

$$F_\varepsilon^*(u, B) := \inf \{F_\varepsilon(u, U) : U \in \mathcal{A}(\Omega), B \subset U\}.$$

We now estimate the last term in the right-hand side of (3.7). To this end set $S := (U'' \setminus U') \cap V$. For any fixed $\eta \in (0, 1)$ we have

$$\begin{aligned}
F_\varepsilon^*(w, S) &\leq \int_{S \cap \varepsilon P} \left| (1-\eta) \frac{\varphi \nabla u + (1-\varphi) \nabla v}{1-\eta} + \eta \frac{\nabla \varphi (u-v)}{\eta} \right|^2 dx \\
&\quad + \delta_\varepsilon \int_{S \setminus \varepsilon P} \left| (1-\eta) \frac{\varphi \nabla u + (1-\varphi) \nabla v}{1-\eta} + \eta \frac{\nabla \varphi (u-v)}{\eta} \right|^2 dx \\
&\quad + \mathcal{H}^{n-1}(S_u \cap S) + \mathcal{H}^{n-1}(S_v \cap S) \\
&\leq \frac{1}{1-\eta} \left(\int_{S \cap \varepsilon P} |\nabla u|^2 dx + \int_{S \cap \varepsilon P} |\nabla v|^2 dx \right) \\
&\quad + \frac{\delta_\varepsilon}{1-\eta} \left(\int_{S \setminus \varepsilon P} |\nabla u|^2 dx + \int_{S \setminus \varepsilon P} |\nabla v|^2 dx \right) \\
&\quad + \frac{1}{\eta} \int_S |\nabla \varphi|^2 |u-v|^2 dx + \mathcal{H}^{n-1}(S_u \cap S) + \mathcal{H}^{n-1}(S_v \cap S) \\
&\leq \frac{1}{1-\eta} \left(F_\varepsilon^*(u, S) + F_\varepsilon^*(v, S) \right) + \frac{1}{\eta} \int_S |\nabla \varphi|^2 |u-v|^2 dx.
\end{aligned} \tag{3.8}$$

Finally, setting $M := \|\nabla \varphi\|_{L^\infty}$ and combining (3.7) and (3.8) give

$$F_\varepsilon(w, U' \cup V) \leq \frac{1}{1-\eta} \left(F_\varepsilon(u, U'') + F_\varepsilon(v, V) \right) + \frac{M}{\eta} \|u-v\|_{L^2(S)}^2,$$

and hence the thesis. \square

We are now ready to state and prove a compactness result by Γ -convergence. We notice that in the proof of Theorem 3 below the only difference from the general approach developed in [17, Chapter 18] is that in our case the rest in the fundamental estimate is small when computed along sequences converging in $L^2(\Omega)$ (see Lemma 1) while we are interested in the Γ -convergence of F_ε with respect to the $L^1(\Omega)$ -topology.

Theorem 3 (Compactness by Γ -convergence). *Let F be as in (3.6). Then*

- (a) (measure property) *for every $u \in GSBV^2(\Omega)$ the set function $F(u, \cdot)$ is the restriction to $\mathcal{A}(\Omega)$ of a Radon measure on Ω ;*
- (b) (compactness) *for every $U \in \mathcal{A}(\Omega)$*

$$F(\cdot, U) = F'(\cdot, U) = F''(\cdot, U) \quad \text{on } GSBV^2(\Omega);$$

- (c) (translational invariance in x) *for every $u \in L^1(\Omega)$ and $U \in \mathcal{A}(\Omega)$*

$$F(u(\cdot - y), U + y) = F(u, U) \quad \text{for all } y \in \mathbb{R}^n \text{ such that } U + y \subset\subset \Omega;$$

- (d) (translational invariance in u) *for every $u \in L^1(\Omega)$ and $U \in \mathcal{A}(\Omega)$*

$$F(u + s, U) = F(u, U) \quad \text{for all } s \in \mathbb{R}.$$

Proof. The proof of (a) relies on the measure-property criterion of De Giorgi and Letta [17, Theorem 14.23]. The only delicate point here is to show that F is subadditive. This can be handled first appealing to a standard truncation argument and to the fact that on $SBV^2(\Omega) \cap L^\infty(\Omega)$ there is an equivalent sequential characterization of (3.6) (see also [17, Proposition 16.4 and Remark 16.5]) and then arguing as in e.g. [21, Theorem 3.4].

The proof of (b) readily follows from (3.6) once we prove that, for every $u \in GSBV^2(\Omega)$, $F''(u, \cdot)$ is inner regular. Indeed, by definition of inner regular envelope and by monotonicity we always have

$$F = F'' = F'_- \leq F' \leq F'', \quad (3.9)$$

hence if we show that $F'' \leq F''$ (which is equivalent to the inner regularity of F''), from (3.9) we may conclude that $F = F' = F''$ and hence (b).

Appealing to Lemma 1 and to the upper bound in (3.5), we now prove that for every fixed $u \in GSBV^2(\Omega)$ the set function $F''(u, \cdot)$ is inner regular on $\mathcal{A}(\Omega)$. To this end, let $u \in GSBV^2(\Omega)$ and fix $W \in \mathcal{A}(\Omega)$. Since $MS(u, \cdot)$ is a Radon measure, for every $\eta > 0$ there exists a compact set $K \subset W$ such that $MS(u, W \setminus K) < \eta$.

Choose $U, U' \in \mathcal{A}(\Omega)$ satisfying $K \subset U' \subset\subset U \subset\subset W$ and set $V := W \setminus K$. Recalling that F'' is increasing, Lemma 1 easily yields

$$F''(u, W) \leq F''(u, U' \cup V) \leq F''(u, U) + F''(u, V) = F''(u, U) + F''(u, W \setminus K).$$

Moreover, by the definition of F'' and in view of the bound $F'' \leq MS$ we have

$$F''(u, W) \leq F''(u, W) + MS(u, W \setminus K) \leq F''(u, W) + \eta.$$

Hence by the arbitrariness of $\eta > 0$ we get

$$F''(u, W) \leq F''(u, W) \quad \text{for every } W \in \mathcal{A}(\Omega),$$

and thus the inner regularity of $F''(u, \cdot)$ for $u \in GSBV^2(\Omega)$.

Finally, by virtue of (3.3) and (3.4), the proof of (c) and (d) follows by standard arguments (see e.g. [11, Lemma 3.7]) \square

On account of the compactness result Theorem 3 we are now ready to accomplish the proof of Theorem 1, that is to represent the Γ -limit F in an integral form. This last step will be accomplished by means of a standard truncation argument and then appealing to the representation result [7, Theorem 1]. For the reader's convenience, below we rephrase [7, Theorem 1] adapting it to our specific setting.

Theorem 4. *Let $G: SBV^2(\Omega) \times \mathcal{A}(\Omega) \rightarrow [0, +\infty]$ be a functional satisfying for every $(u, U) \in SBV^2(\Omega) \times \mathcal{A}(\Omega)$ the following conditions:*

- (i) $G(\cdot, U)$ is local;
- (ii) $G(\cdot, U)$ is $L^1(\Omega)$ -lower semicontinuous;
- (iii) $G(u, \cdot)$ is the restriction to $\mathcal{A}(\Omega)$ of a Radon measure;
- (iv) $G(\cdot, U)$ is invariant under translation in x and in u ;
- (v) there exists $C > 0$ such that

$$\frac{1}{C} MS(u, U) \leq G(u, U) \leq C(MS(u, U) + 1).$$

Then, there exist Borel functions $f_0: \mathbb{R}^n \rightarrow [0, +\infty)$ and $g_0: \mathbb{R} \times S^{n-1} \rightarrow [0, +\infty)$ such that

$$G(u, U) = \int_U f_0(\nabla u) \, dx + \int_{S_u \cap U} g_0([u], \nu_u) \, d\mathcal{H}^{n-1}$$

for every pair $(u, U) \in SBV^2(\Omega) \times \mathcal{A}(\Omega)$.

Moreover, $g_0(t, \nu) = g_0(-t, -\nu)$ for every $t \in \mathbb{R}$, and $\nu \in S^{n-1}$.

Proof. The proof follows from [7, Theorem 1] by means of a perturbation argument. To this end let $\sigma > 0$ and for every $(u, U) \in SBV^2(\Omega) \times \mathcal{A}(\Omega)$ define the functionals

$$G_\sigma(u, U) := G(u, U) + \sigma \int_{S_u \cap U} (1 + |[u]|) d\mathcal{H}^{n-1}.$$

Clearly

$$\begin{aligned} G_\sigma(u, U) &\geq \frac{1}{C_\sigma} \left(MS(u, U) + \int_{S_u \cap U} |[u]| d\mathcal{H}^{n-1} \right), \\ G_\sigma(u, U) &\leq C_\sigma \left(MS(u, U) + \int_{S_u \cap U} |[u]| d\mathcal{H}^{n-1} + 1 \right) \end{aligned}$$

for $C_\sigma > 0$, thus G_σ satisfies the growth conditions in [7, Theorem 1]. Moreover, it is immediate to check that G_σ satisfies the other hypotheses of [7, Theorem 1] (the lower semicontinuity being a consequence of [1, Theorem 3.7]). Then, we deduce the existence of two Borel functions $f_\sigma: \mathbb{R}^n \rightarrow [0, +\infty)$ and $g_\sigma: \mathbb{R} \times \mathbb{S}^{n-1} \rightarrow [0, +\infty)$ such that

$$G_\sigma(u, U) = \int_U f_\sigma(\nabla u) dx + \int_{S_u \cap U} g_\sigma([u], \nu_u) d\mathcal{H}^{n-1}$$

for every $(u, U) \in SBV^2(\Omega) \times \mathcal{A}(\Omega)$.

By construction the two families of functions (f_σ) and (g_σ) are decreasing as σ decreases (cf. equations (1)–(3), [7, Theorem 1]). Hence, setting $f_0 := \lim_{\sigma \rightarrow 0^+} f_\sigma$ and $g_0 := \lim_{\sigma \rightarrow 0^+} g_\sigma$, by the pointwise convergence of (G_σ) to G and the Monotone Convergence Theorem we deduce

$$G(u, U) = \int_U f_0(\nabla u) dx + \int_{S_u \cap U} g_0([u], \nu_u) d\mathcal{H}^{n-1}$$

for every $(u, U) \in SBV^2(\Omega) \times \mathcal{A}(\Omega)$, and thus the thesis. \square

We are now in a position to prove the Γ -convergence result Theorem 1.

Proof of Theorem 1. Let F be as in (3.6) and let (ε_k) be the corresponding vanishing subsequence. Then, Theorem 3 ensures that

$$F(u, U) := \Gamma\text{-}\lim_{k \rightarrow +\infty} F_{\varepsilon_k}(u, U)$$

for every $(u, U) \in GSBV^2(\Omega) \times \mathcal{A}(\Omega)$ and hence, in particular, that $F(\cdot, U)$ is $L^1(\Omega)$ -lower semicontinuous on $GSBV^2(\Omega)$. Moreover, again appealing to Theorem 3 and to the bound (3.5) ensures that F fulfils the set of hypotheses of Theorem 4. Therefore we deduce the existence of two Borel functions $f_0: \mathbb{R}^n \rightarrow [0, +\infty)$ and $g_0: \mathbb{R} \times \mathbb{S}^{n-1} \rightarrow [0, +\infty)$ such that

$$F(u, U) = \int_U f_0(\nabla u) dx + \int_{S_u \cap U} g_0([u], \nu_u) d\mathcal{H}^{n-1}, \quad (3.10)$$

for every $(u, U) \in SBV^2(\Omega) \times \mathcal{A}(\Omega)$.

Let $y \in \mathbb{R}^n$, $r > 0$, and $\nu \in \mathbb{S}^{n-1}$. We denote by $Q_r^\nu(y)$ the open cube of centre y , side-length $r > 0$, and one face orthogonal to ν . Now fix $y \in \Omega$ and let $r > 0$ be such that $Q_r^\nu(y) \subset \Omega$; then the bounds $c_1|\xi|^2 \leq f_0(\xi) \leq |\xi|^2$ and $c_2 \leq g_0(t, \nu) \leq 1$ readily follow from (2.1), (2.2), (the localised version of) (2.4), and (3.10) by testing $F(\cdot, Q_r^\nu(y))$ on

$$u_\xi(x) = \xi \cdot x \quad \text{and} \quad u_t^\nu(x) = \begin{cases} t & \text{if } (x - y) \cdot \nu > 0, \\ 0 & \text{if } (x - y) \cdot \nu \leq 0, \end{cases}$$

respectively.

Moreover, it is immediate to prove that $g_0(\cdot, \nu)$ is nondecreasing in $(0, +\infty)$ for every fixed $\nu \in \mathbb{S}^{n-1}$ (see e.g. [21, Proof of Theorem 2.2]) while the symmetry property directly follows from [7, Theorem 1]. Finally, the left-continuity of $g_0(\cdot, \nu)$ in $(0, +\infty)$ is a consequence of its monotonicity and of the lower semicontinuity of F .

We now show that the representation formula for F holds on the whole $GSBV^2(\Omega)$. To this end, for $(u, U) \in GSBV^2(\Omega) \times \mathcal{A}(\Omega)$ set

$$\tilde{F}(u, U) := \int_U f_0(\nabla u) \, dx + \int_{S_u \cap U} g_0([u], \nu_u) \, d\mathcal{H}^{n-1},$$

and for every $M > 0$ let $u^M := (u \wedge M) \vee (-M)$. Notice that \tilde{F} decreases by truncation and that it is $L^1(\Omega)$ -lower semicontinuous on $GSBV^2(\Omega)$. Then by the $L^1(\Omega)$ -lower semicontinuity of F on $GSBV^2(\Omega)$ and by (3.10) we have

$$F(u, U) \leq \liminf_{M \rightarrow +\infty} F(u^M, U) = \liminf_{M \rightarrow +\infty} \tilde{F}(u^M, U) \leq \tilde{F}(u, U). \quad (3.11)$$

On the other hand, since F decreases by truncation we may also deduce

$$\tilde{F}(u, U) \leq \liminf_{M \rightarrow +\infty} \tilde{F}(u^M, U) = \liminf_{M \rightarrow +\infty} F(u^M, U) \leq F(u, U). \quad (3.12)$$

Hence gathering (3.11) and (3.12) gives the desired representation result on $GSBV^2(\Omega)$. Then, choosing $U = \Omega$ yields the Γ -convergence result.

To conclude the proof it only remains to show that $f_0 = \hat{f}$. To this end, for any fixed $\delta \in (0, 1)$ consider the functionals $F_\varepsilon^\delta : L^1(\Omega) \rightarrow [0, +\infty]$ given by

$$F_\varepsilon^\delta(u) := \begin{cases} \int_{\Omega \cap \varepsilon P} |\nabla u|^2 \, dx + \delta \int_{\Omega \setminus \varepsilon P} |\nabla u|^2 \, dx & \text{if } u \in H^1(\Omega), \\ +\infty & \text{otherwise in } L^1(\Omega). \end{cases}$$

Then, (F_ε^δ) Γ -converges to the quadratic form

$$F^\delta(u) := \begin{cases} \int_\Omega f^\delta(\nabla u) \, dx & \text{if } u \in H^1(\Omega), \\ +\infty & \text{otherwise in } L^1(\Omega), \end{cases} \quad (3.13)$$

where $f^\delta : \mathbb{R}^n \rightarrow [0, +\infty)$ is defined as

$$f^\delta(\xi) = \min \left\{ \int_Q (\chi_{Q \cap P}(x) + \delta \chi_{Q \setminus P}(x)) |\xi + \nabla w|^2 \, dx : w \in H_{\text{per}}^1(Q) \right\},$$

for every $\xi \in \mathbb{R}^n$ (see [9, Subsection 5.2]). For sufficiently small $\varepsilon > 0$ we have $\hat{F}_\varepsilon \leq F_\varepsilon \leq F_\varepsilon^\delta$ and therefore $\hat{f} \leq f_0 \leq f^\delta$. We show that

$$\lim_{\delta \rightarrow 0^+} f^\delta(\xi) = \hat{f}(\xi) \quad (3.14)$$

for every $\xi \in \mathbb{R}^n$. To this end, we first assume that P is Lipschitz so that

$$\hat{f}(\xi) = \inf \left\{ \int_{Q \cap P} |\xi + \nabla w|^2 \, dx : w \in H_{\text{per}}^1(Q) \right\}.$$

Fix $\xi \in \mathbb{R}^n$; for any given $\eta > 0$, let $w^\eta \in H_{\text{per}}^1(Q)$ be such that

$$\int_{Q \cap P} |\xi + \nabla w^\eta|^2 \, dx \leq \hat{f}(\xi) + \eta.$$

Then

$$f^\delta(\xi) \leq \int_Q (\chi_{Q \cap P}(x) + \delta \chi_{Q \setminus P}(x)) |\xi + \nabla w^\eta|^2 dx \leq \hat{f}(\xi) + \eta + \delta \int_{Q \setminus P} |\xi + \nabla w^\eta|^2 dx.$$

Thus (3.14) follows by letting δ tend to zero and by the arbitrariness of $\eta > 0$.

Finally, if P is not Lipschitz, the fact that $f_0 = \hat{f}$ can be recovered as in [5, equation (44)] by approximating P with an increasing sequence of Lipschitz open sets (P_m) invading P . \square

4. DEPENDENCE ON THE CRACK-OPENING

In the present section we prove the main result of this paper, namely Theorem 2. To this end, we need several ingredients which are collected in the following lemmas.

Henceforth, without loss of generality, we restrict ourselves to the case $\Omega = Q$, cf. (4.2).

Lemma 2. *Let (U_k) be a sequence of Borel subsets of Q such that $\chi_{U_k} \rightharpoonup \theta$ weakly* in $L^\infty(Q)$, with $\theta \in (0, 1]$. Let $(u_k), (v_k)$ be two sequences in $L^1(Q)$ such that $u_k = v_k$ \mathcal{L}^2 -a.e. in U_k . Then, if $u_k \rightarrow u$ and $v_k \rightarrow v$ in $L^1(Q)$, we have $u = v$ \mathcal{L}^2 -a.e. in Q .*

Proof. For every fixed $k \in \mathbb{N}$, $(u_k - v_k)\chi_{U_k} = 0$ \mathcal{L}^2 -a.e. in Q . On the other hand $(u_k - v_k)\chi_{U_k} \rightharpoonup (u - v)\theta$ weakly in $L^1(Q)$, then by virtue of the positivity of θ we immediately infer that $u = v$ \mathcal{L}^2 -a.e. in Q . \square

The following lemma allows us to replace a converging sequence (u_k) by a sequence of functions v_k , periodic and symmetric in the first variable, still converging to the same limit, without essentially increasing the energy F_ε .

Lemma 3. *Set $u_t := t \chi_{Q \cap \{x_2 > 0\}}$ and let $(u_k) \subset L^1(Q)$ be such that $u_k \rightarrow u_t$ in $L^1(Q)$. Assume that (ε_k^{-1}) is a sequence of odd integers (so that Q is union of periodicity cells). Then, for any fixed $\eta > 0$ there exists $(v_k) \subset L^1(Q)$, ε_k -periodic and symmetric in its first variable (i.e., $v_k(x_1, x_2) = v_k(x_1 + \varepsilon_k, x_2)$ and $v_k(x_1, x_2) = v_k(-x_1, x_2)$ for every $(x_1, x_2) \in Q$) such that $v_k \rightarrow u_t$ in $L^1(Q)$ and*

$$\liminf_{k \rightarrow +\infty} F_{\varepsilon_k}(u_k) \geq \liminf_{k \rightarrow +\infty} F_{\varepsilon_k}(v_k) - \eta. \quad (4.1)$$

Proof. Let S_k^i be the open strip $(i\varepsilon_k/2, (i+1)\varepsilon_k/2) \times (-1/2, 1/2)$ and let (λ_k) be a sequence such that

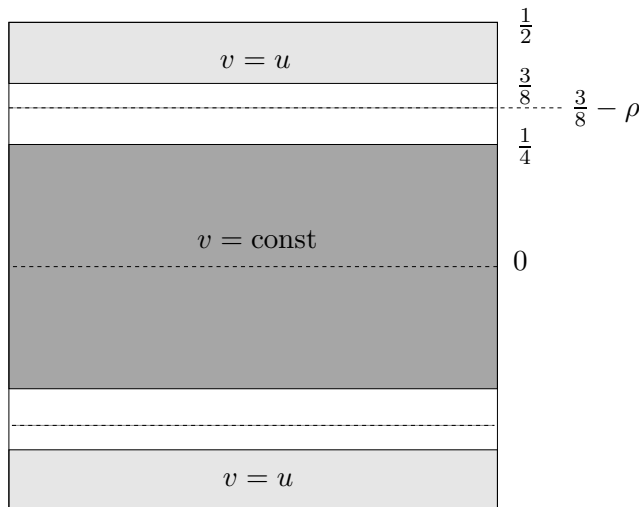
$$\lambda_k \rightarrow +\infty \quad \text{and} \quad \sup_k \left(\lambda_k \int_Q |u_k - u_t| dx \right) \leq \eta.$$

Let $i_k \in \mathbb{Z}$ be a solution to

$$\min \left\{ F_{\varepsilon_k}(u_k, S_k^{i_k}) + \lambda_k \int_{S_k^{i_k}} |u_k - u_t| dx : i = -\varepsilon_k^{-1}, \dots, \varepsilon_k^{-1} - 1 \right\},$$

where $F_{\varepsilon_k}(u_k, S_k^{i_k})$ is defined as in (3.1). In order to define v_k we first extend $u_k|_{S_k^{i_k}}$ to the strip $((i_k - 1)\varepsilon_k/2, (i_k + 1)\varepsilon_k/2) \times (-1/2, 1/2)$ by reflection with respect to $x_1 = i_k\varepsilon_k/2$; we denote by \tilde{u}_k such an extension. Then we extend \tilde{u}_k by ε_k -periodicity in the x_1 -variable to the whole $\mathbb{R} \times (-1/2, 1/2)$. Hence we call v_k the latter extension. Then, (4.1) holds true by definition of v_k , while the penalisation term also ensures that $v_k \rightarrow u_t$ in $L^1(Q)$. \square

Lemma 5 below represents a key ingredient in the proof of Theorem 2 and is a consequence of the elimination lemma [19, Lemma 0.7], which we recall for the reader's convenience.


 FIGURE 4. The function v of Lemma 5.

Definition 1. Let $V \subset \mathbb{R}^n$ be open. We say that $v \in SBV^2(V) \cap L^\infty(V)$ is a local minimiser for $MS(\cdot, V)$ if $MS(v, A) \leq MS(w, A)$ for every open set $A \subset\subset V$, whenever $w \in SBV^2(V) \cap L^\infty(V)$ and $\{w \neq v\} \subset\subset A \subset\subset V$.

Lemma 4 (Elimination property). *Let $V \subset \mathbb{R}^n$ be open. There exists a dimensional constant $\vartheta = \vartheta(n) > 0$, independent of V , such that, if $v \in SBV^2(V) \cap L^\infty(V)$ is a local minimiser for $MS(\cdot, V)$ and $B_{2\rho}(\bar{x}) \subset V$ is any ball with centre $x_0 \in V$ and radius 2ρ with*

$$\mathcal{H}^{n-1}(S_v \cap B_{2\rho}(\bar{x})) < \vartheta(2\rho)^{n-1},$$

then $u \in H^1(B_\rho)$.

Lemma 5. *There exist two constants $\alpha, \beta > 0$ with the following property: Let $u \in SBV^2(Q) \cap L^\infty(Q)$ be symmetric in its first variable and such that $MS(u, U) \leq \beta$, where $U := (-1/2, 1/2) \times (-3/8, 3/8)$. Then there exists a function $v \in SBV^2(Q) \cap L^\infty(Q)$ still symmetric and such that*

- (i) v is constant in $(-1/2, 1/2) \times (-1/4, 1/4)$;
- (ii) $v = u$ in $Q \setminus U$;
- (iii) $\|v\|_{L^\infty(Q)} \leq \|u\|_{L^\infty(Q)}$;
- (iv) $\alpha \int_Q |\nabla v|^2 dx + \mathcal{H}^1(S_v \cap Q) \leq MS(u, Q)$.

Proof. We first substitute u with a solution w to the minimisation problem

$$\min \{MS(w, Q) : w \in SBV^2(Q) \cap L^\infty(Q), w = u \text{ in } Q \setminus U\}.$$

We clearly have $\|w\|_{L^\infty(Q)} \leq \|u\|_{L^\infty(Q)}$. Since u is symmetric in x_1 , we can always assume that w inherits the same property. Moreover, we can extend w by reflection to $(-3/2, 3/2) \times (-1/2, 1/2)$ keeping the minimality with respect to all competitors coinciding with u in $(-3/2, 3/2) \times ((-1/2, 1/2) \setminus (-3/8, 3/8))$, where u is extended by reflection, too. This allows us to apply the elimination property Lemma 4. More precisely, let $\rho \in (0, 1/8)$ be fixed and $\beta := 2\rho\vartheta$, where $\vartheta = \vartheta(2)$ is as in Lemma 4. If $B_{2\rho}(\bar{x}) \subset (-3/2, 3/2) \times (-3/8, 3/8)$, then our assumptions imply $\mathcal{H}^1(S_w \cap B_{2\rho}(\bar{x})) \leq \beta$, hence in its turn the elimination property ensures that $w \in H^1(B_\rho(\bar{x}))$. By moving \bar{x} in $(-3/2, 3/2) \times [-3/8+2\rho, 3/8-2\rho]$, we infer that $w \in H^1((-1/2, 1/2) \times (-3/8+\rho, 3/8-\rho))$.

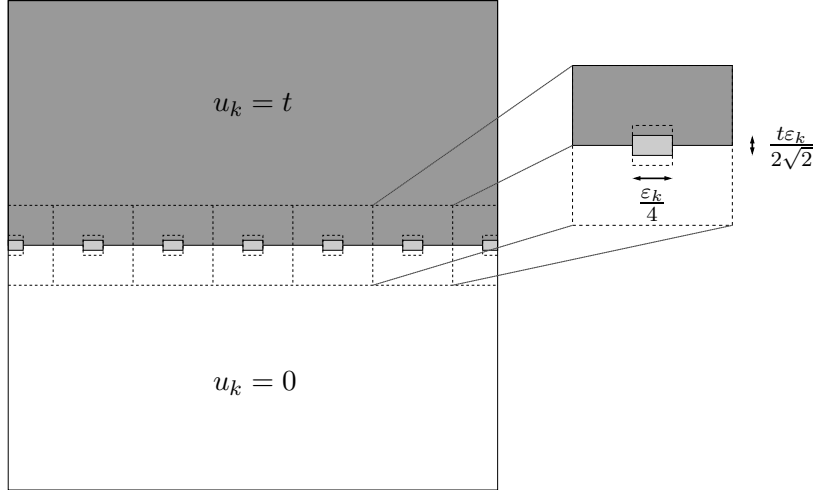


FIGURE 5. The “bridging” sequence (u_k) . The discontinuity set S_{u_k} is given by the union of the horizontal interface between the white and the dark grey zone and the vertical part of the boundary of the light grey rectangles. The light grey rectangles (where u_k is affine) are contained in the soft inclusions and act as “bridging elements” in the deformed configuration.

The desired function v is finally obtained by means of a standard cut-off argument (see Figure 4). Specifically, let φ be a smooth function defined on Q and depending only on the second variable, with $\varphi(x_1, x_2) = 1$ for $|x_2| > 5/16$ and $\varphi(x_1, x_2) = 0$ for $|x_2| < 1/4$. We define $v := (w-m)\varphi + m$, where m is the mean value of w on $(-1/2, 1/2) \times (-5/16, 5/16)$. Then (i)–(iii) hold by construction, while (iv) follows by the Poincaré-Wirtinger inequality, recalling that w is a Sobolev function in $(-1/2, 1/2) \times (-5/16, 5/16)$. \square

Remark 6. *It is worth mentioning that the previous lemma can be generalised to any dimension n with some $\beta = \beta(n)$ (compare with Lemma 7 in Section 5 which is a purely two-dimensional argument). Moreover, by using the scaled functions $x \mapsto u(\varepsilon x)/\sqrt{\varepsilon}$, Lemma 5 can be stated in the cube Q_ε up to replacing β by $\varepsilon^{n-1}\beta$.*

We are now in a position to prove Theorem 2. By using approximate displacements u_k mimicking a “bridging effect” in the microstructure (see e.g. Figure 5), we will show that the surface energy density g_0 depends on the crack-opening in a nontrivial way.

Proof of Theorem 2. We start noticing that in view of the symmetry of the microgeometry the functionals F_ε are invariant under ninety-degree rotations. Then it is immediate to check that $g_0(t, e_1) = g_0(t, e_2)$ for every $t \geq 0$. Thus it is enough to prove (2.8) when $i = 2$. By Theorem 1 we have that

$$g_0(t, e_2) = F(u_t, Q), \quad \text{where } u_t := t \chi_{Q \cap \{x_2 > 0\}}. \quad (4.2)$$

We divide the proof into two main parts.

Estimate from above. By (2.4) we immediately deduce that $g_0(t, e_2) \leq 1$ for every $t > 0$. Then, to prove that $g_0(\cdot, e_2)$ satisfies the estimate from above in (2.8), we need to show that $g_0(t, e_2) \leq 3/4 + \sqrt{2}t$ whenever $t \leq 1/(4\sqrt{2})$. To this end, let $t \leq 1/(4\sqrt{2})$ and let $R \subset \mathbb{R}^2$ be the open rectangle defined as

$$R := \left(-\frac{1}{8}, \frac{1}{8}\right) \times \left(-\frac{t}{4\sqrt{2}}, \frac{t}{4\sqrt{2}}\right).$$

With this choice of t we also have $R \subset Q_{1/4}$. Set

$$R_k := Q \cap \varepsilon_k \bigcup_{i \in \mathbb{Z}} (R + (i, 0))$$

and let $(u_k) \subset SBV^2(Q)$ be the sequence of “bridging” functions defined as

$$u_k(x) := \begin{cases} t & \text{if } x \in Q \setminus R_k \text{ and } x_2 \geq 0, \\ \frac{t}{2} + \frac{2\sqrt{2}}{\varepsilon_k} x_2 & \text{if } x \in R_k, \\ 0 & \text{if } x \in Q \setminus R_k \text{ and } x_2 < 0, \end{cases}$$

(see Figure 5). We clearly have $u_k \rightarrow u_t$ in $L^1(Q)$; moreover

$$\int_{R_k} |\nabla u_k|^2 dx \leq (\lfloor \frac{1}{\varepsilon_k} \rfloor + 1) \left(\frac{t}{\sqrt{2}} \right) \quad \text{and} \quad \mathcal{H}^1(S_{u_k}) \leq \varepsilon_k (\lfloor \frac{1}{\varepsilon_k} \rfloor + 1) \left(\frac{3}{4} + \frac{t}{\sqrt{2}} \right).$$

Thus we readily deduce

$$g_0(t, e_2) = F(u_t, Q) \leq \limsup_{k \rightarrow +\infty} F_{\varepsilon_k}(u_k, Q) \leq \frac{3}{4} + \sqrt{2}t,$$

and hence the estimate from above.

Estimate from below. Let $u_k \rightarrow u_t$ in $L^1(Q)$; we want to prove that

$$\liminf_k F_{\varepsilon_k}(u_k, Q) \geq \min\left\{ \frac{3}{4} + c_3 t^2, 1 \right\} \quad (4.3)$$

for some $c_3 > 0$.

The strategy of the proof is to show that the energy essentially concentrates in a horizontal layer of thickness proportional to ε_k , giving a contribution controlled from below as in (4.3); this is shown in Step 4. In order to prove this, we suitably modify (u_k) obtaining a new sequence (w_k) that is Sobolev outside the layer as above and whose gradient has equibounded L^2 -norm in the same layer. The construction of the sequence (w_k) is performed in Steps 1, 2, and 3 and is based on consecutive modifications of (u_k) , which have the property of keeping the limit u_t . Indeed, at each step the sequence remains unchanged in “large” portions of each periodicity cell, so that the convergence to u_t is preserved thanks to Lemma 2. Moreover the various modifications of (u_k) are such that their energy changes in a way that can be suitably controlled in terms of $F_{\varepsilon_k}(u_k, Q)$.

We start noticing that we may assume

$$\liminf_k F_{\varepsilon_k}(u_k, Q) \leq 1, \quad (4.4)$$

otherwise there is nothing to prove. Up to an extraction of a subsequence, we can additionally assume that the liminf in (4.4) is actually a limit. In this way we can modify the sequence (u_k) , possibly extracting further subsequences, keeping the estimates.

Since F_{ε_k} decreases by truncation, we can also suppose that $0 \leq u_k \leq t$ in Q . Moreover, we assume that ε_k^{-1} is an odd integer and, in view of Lemma 3, that u_k is ε_k -periodic and symmetric in its first variable.

We divide the proof into four steps.

Step 1. Appealing to Lemma 5, in this step we modify u_k replacing it by constant values in suitably chosen regions of the domain Q .

We decompose Q into squares of side-length ε_k by setting $Q_{\varepsilon_k}^{i,j} := Q_{\varepsilon_k} + p_k^{i,j}$, where

$$p_k^{i,j} := \left(-\frac{1}{2} + (i - \frac{1}{2})\varepsilon_k, -\frac{1}{2} + (j - \frac{1}{2})\varepsilon_k \right), \quad i, j \in \left\{ 1, \dots, \frac{1}{\varepsilon_k} \right\}.$$

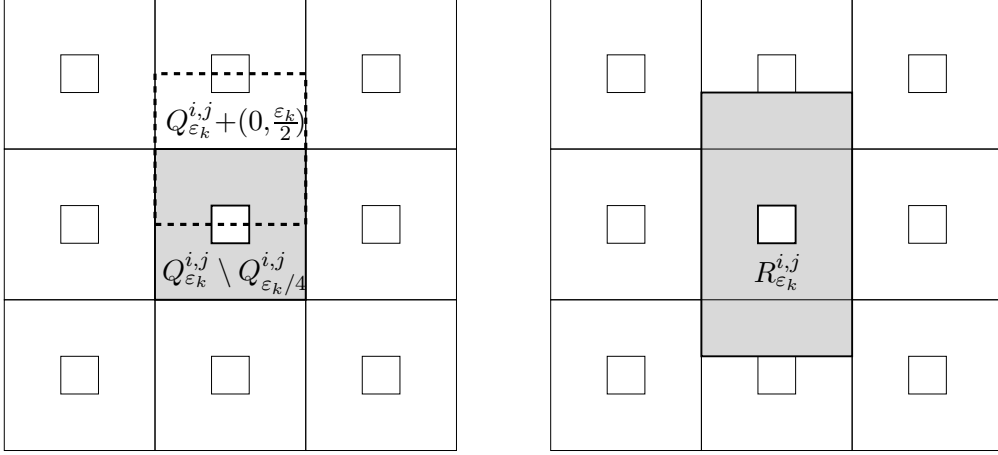


FIGURE 6. In grey the sets $Q_{\varepsilon_k}^{i,j} \setminus Q_{\varepsilon_k/4}^{i,j}$ and $R_{\varepsilon_k}^{i,j}$. The dashed line indicates the boundary of the translated square $Q_{\varepsilon_k}^{i,j} + (0, \varepsilon_k/2)$.

We also consider the smaller squares $Q_{\varepsilon_k/4}^{i,j} := Q_{\varepsilon_k/4} + p_k^{i,j}$ and the strips

$$T_{\varepsilon_k}^j := \left(\left(-\frac{1}{2}, \frac{1}{2}\right) \times \left(-\frac{\varepsilon_k}{2}, \frac{\varepsilon_k}{2}\right) \right) + \left(0, -\frac{1}{2} + \left(j - \frac{1}{2}\right)\varepsilon_k\right),$$

$$T_{\frac{\varepsilon_k}{2}}^j := \left(\left(-\frac{1}{2}, \frac{1}{2}\right) \times \left(-\frac{\varepsilon_k}{4}, \frac{\varepsilon_k}{4}\right) \right) + \left(0, -\frac{1}{2} + \left(j - \frac{1}{2}\right)\varepsilon_k\right).$$

Moreover, we define the sets $R_{\varepsilon_k}^{i,j} := \varepsilon_k R + p_k^{i,j}$ (see Figure 6), where

$$R := \left(\left(-1/2, 1/2\right) \times \left(-7/8, 7/8\right) \right) \setminus \overline{Q}_{1/4}.$$

Let $\mathcal{J}_k \subset \{2, \dots, \varepsilon_k^{-1} - 1\}$ be the set of indices j such that

$$MS(u_k, R_{\varepsilon_k}^{1,j}) \leq \beta \varepsilon_k,$$

where $\beta > 0$ will be chosen later. Thanks to the periodicity of u_k , for $j \in \mathcal{J}_k$ we have

$$MS(u_k, R_{\varepsilon_k}^{i,j}) \leq \beta \varepsilon_k,$$

for every $i \in \{1, \dots, \varepsilon_k^{-1}\}$.

For β small enough, by scaling and translation (see also Remark 6) we can apply Lemma 5 in each square $Q_{\varepsilon_k}^{i,j} \pm (0, \varepsilon_k/2)$, for $i \in \{1, \dots, \varepsilon_k^{-1}\}$ and $j \in \mathcal{J}_k$, thus obtaining a new sequence (v_k) which is equal to (u_k) (at least) in $Q_{\varepsilon_k/4}^{i,j}$ and is constant in the translated strips $T_{\varepsilon_k/2}^j + (0, \varepsilon_k/2)$, $T_{\varepsilon_k/2}^j - (0, \varepsilon_k/2)$ (the dark grey part in Figure 7(i)). We call these two constant values b_k^j and a_k^j , respectively. Moreover, Lemma 5 yields the existence of a constant $\alpha > 0$ such that

$$\alpha \int_{Q_k} |\nabla v_k|^2 dx + \mathcal{H}^1(S_{v_k} \cap Q_k) \leq MS(u_k, Q_k), \quad (4.5)$$

where $Q_k := Q \cap \varepsilon_k P$ is the stiff matrix, cf. (2.7). Note that $v_k = u_k$ in $Q \setminus \varepsilon_k P$.

Due to the lack of control on ∇u_k in the soft inclusions $Q \setminus \varepsilon_k P$, it is not obvious that $v_k \rightarrow u_t$ in $L^1(Q)$. Then, to prove the desired convergence we employ a result from [5], originally stated for perforated domains. In fact, by [5, Theorem 1] one can find a sequence (\bar{v}_k) such that $\bar{v}_k = v_k$

in Q_k and $\bar{v}_k \rightarrow u_t$ strongly in $L^1(Q)$ (up to subsequences not relabelled). Hence,

$$\int_Q |v_k - u_t| dx = \int_{Q_k} |\bar{v}_k - u_t| dx + \int_{Q \setminus \varepsilon_k P} |u_k - u_t| dx \rightarrow 0,$$

and the convergence to u_t is preserved.

Step 2. In this step we modify v_k in each strip $T_{\varepsilon_k/2}^j$, for $j \in \mathcal{J}_k$, by replacing it with the affine interpolation between the values a_k^j and b_k^j .

Denote by v_k^{aff} the function defined in $\bigcup_{j \in \mathcal{J}_k} T_{\varepsilon_k/2}^j$ (light grey part in Figure 8(i)) which is constant in its first variable and is the affine interpolation between a_k^j and b_k^j in its second variable. We set

$$\tilde{v}_k := \begin{cases} v_k^{\text{aff}} & \text{in } \bigcup_{j \in \mathcal{J}_k} T_{\varepsilon_k/2}^j, \\ v_k & \text{otherwise in } Q. \end{cases}$$

For each $x_1 \in (-1/2, 1/2)$ let $S_{x_1}^j$ be the vertical one-dimensional slice of $T_{\varepsilon_k/2}^j$ through x_1 ; *i.e.*,

$$S_{x_1}^j := (\{x_1\} \times [-\frac{\varepsilon_k}{4}, \frac{\varepsilon_k}{4}]) + (0, -\frac{1}{2} + (j - \frac{1}{2})\varepsilon_k).$$

By [2, Theorems 3.28, 3.107, and 3.108], for \mathcal{L}^1 -a.e. x_1 such that $S_{x_1}^j \cap S_{v_k} = \emptyset$ the function $v_k(x_1, \cdot)$ is absolutely continuous and therefore

$$|b_k^j - a_k^j| \leq \int_{S_{x_1}^j} |\partial_{x_2} v_k(x_1, s)| ds.$$

Appealing to Jensen's inequality, we get

$$\int_{T_{\varepsilon_k/2}^j} |\nabla v_k^{\text{aff}}|^2 dx = \frac{2}{\varepsilon_k} |b_k^j - a_k^j|^2 \leq \int_{S_{x_1}^j} |\partial_{x_2} v_k(x_1, s)|^2 ds \quad (4.6)$$

for x_1 as above. Let

$$\hat{P} := \mathbb{R}^2 \setminus \left(\left(\bigcup_{i \in \mathbb{Z}} (-\frac{1}{8}, \frac{1}{8}) + i \right) \times \mathbb{R} \right);$$

we notice that the vertical one-dimensional slices of \hat{P} do not intersect any of the soft inclusions. We now assume $\beta \leq 1/4$, so that at least half of the vertical one-dimensional slices in $T_{\varepsilon_k/2}^j \cap \varepsilon_k \hat{P}$ (see Figure 7(ii)) do not intersect S_{v_k} , indeed

$$\mathcal{H}^1(\{x_1 : S_{x_1}^j \subset \varepsilon_k \hat{P} \text{ and } S_{x_1}^j \cap S_{v_k} = \emptyset\}) \geq \frac{3}{4} - \beta \geq \frac{1}{2},$$

where we have used the fact that for $j \in \mathcal{J}_k$

$$\mathcal{H}^1(S_{v_k} \cap T_{\varepsilon_k/2}^j \cap \varepsilon_k \hat{P}) \leq \mathcal{H}^1(S_{u_k} \cap T_{\varepsilon_k/2}^j \cap \varepsilon_k \hat{P}) \leq \beta.$$

Integrating (4.6) on $\{x_1 : S_{x_1}^j \subset \varepsilon_k \hat{P} \text{ and } S_{x_1}^j \cap S_{v_k} = \emptyset\}$ we find

$$\frac{1}{2} \int_{T_{\varepsilon_k/2}^j} |\nabla v_k^{\text{aff}}|^2 dx \leq \int_{T_{\varepsilon_k/2}^j \cap \varepsilon_k \hat{P}} |\nabla v_k|^2 dx.$$

Define

$$C := \bigcup \left\{ \overline{T_{\varepsilon_k}^j} : j \in \mathcal{J}_k \right\}$$

and consider the open set $A := \text{int}(C)$. By construction, we have

$$\frac{1}{2} \int_A |\nabla \tilde{v}_k|^2 dx \leq \int_{A \cap \varepsilon_k P} |\nabla v_k|^2 dx. \quad (4.7)$$

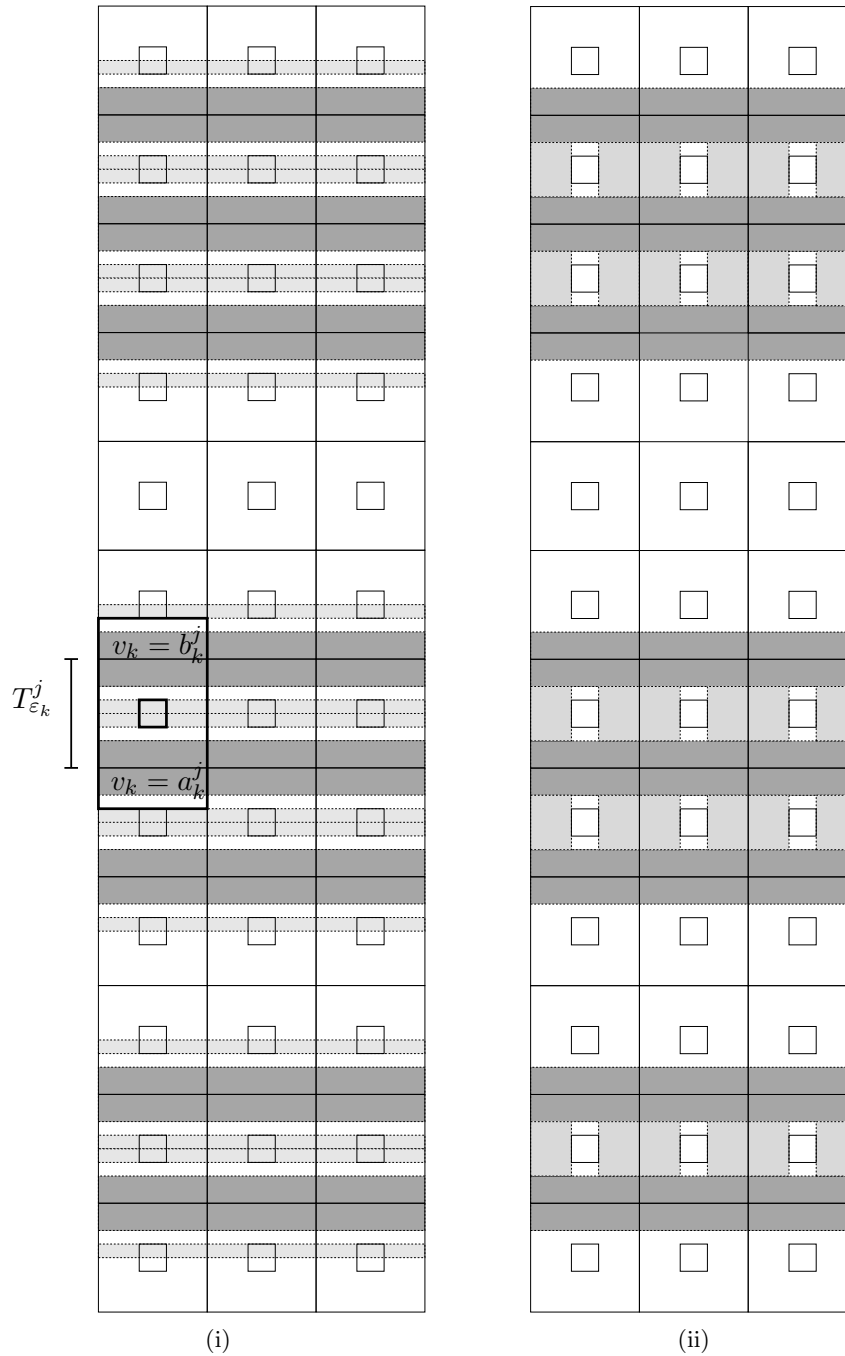


FIGURE 7. (i) In black the boundary of the set $R_{\varepsilon_k}^{i,j}$. The function v_k is defined in Step 1 and assumes constant values in each dark grey strip. (ii) In light grey the intersection of $\varepsilon_k \hat{P}$ with the strips $\bigcup_{j \in \mathcal{J}_k} T_{\varepsilon_k/2}^j$.

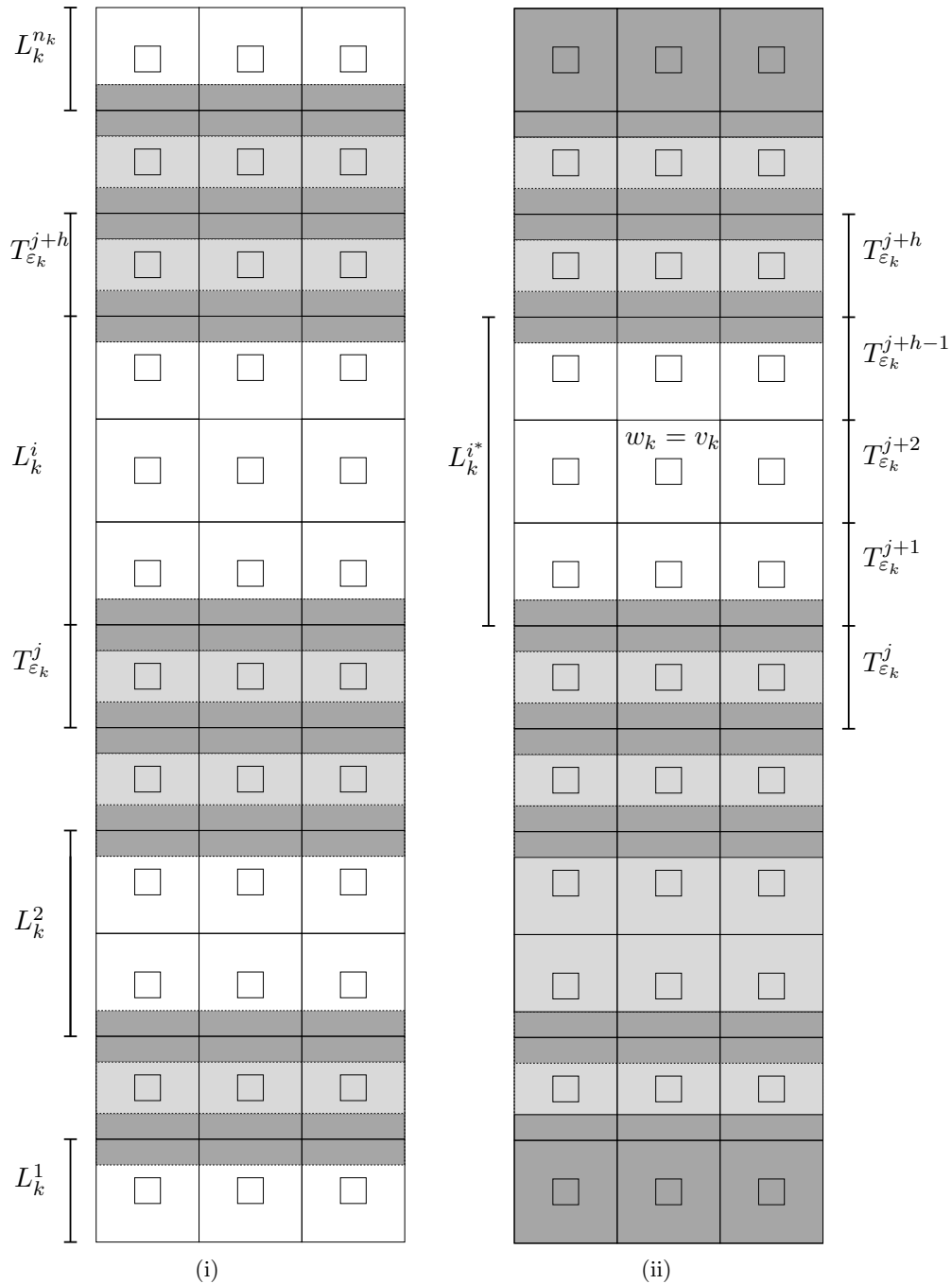


FIGURE 8. (i) The function \tilde{v}_k is defined in Step 2 and is affine in the light grey strips. (ii) The function w_k is defined in Step 3; it assumes constant values in each dark grey strip and is affine in the light grey strips. Note that the surface energy concentrates in the strip $L_{\epsilon_k}^{i*}$.

We want to show that $\tilde{v}_k \rightarrow u_t$ in $L^1(Q)$. To this end we introduce the auxiliary sequence (\hat{v}_k) defined as

$$\hat{v}_k := \begin{cases} \tilde{v}_k & \text{in } A \\ 0 & \text{in } Q \setminus A. \end{cases}$$

Denote by $L_k^1, \dots, L_k^{n_k}$ the connected components of $Q \setminus A$ (see Figure 8(i)), enumerated going from the bottom to the top of Q . By definition of \mathcal{J}_k , we have that $T_{\varepsilon_k}^1 \subset L_k^1$ and $T_{\varepsilon_k}^{\varepsilon_k^{-1}} \subset L_k^{n_k}$. Let \mathcal{J}_k^c be the complement of \mathcal{J}_k in $\{1, \dots, \varepsilon_k^{-1}\}$. For $i \in \{1, \dots, \varepsilon_k^{-1}\}$ and $j \in \mathcal{J}_k^c$ we have

$$MS(u_k, R_{\varepsilon_k}^{i,j}) > \beta \varepsilon_k.$$

Thus, recalling that $\liminf_k F_{\varepsilon_k}(u_k) \leq 1$, for k large enough we deduce

$$n_k \leq \#\mathcal{J}_k^c \leq \frac{2F_{\varepsilon_k}(u_k)}{\beta} \leq \frac{3}{\beta}, \quad (4.8)$$

where $\#\mathcal{J}_k^c$ denotes the number of elements of \mathcal{J}_k^c and the prefactor 2 comes from the partial superposition of the sets $R_{\varepsilon_k}^{i,h}$. We clearly have $(\hat{v}_k) \subset SBV^2(Q)$ with discontinuities concentrated along the boundaries of L_k^h . Gathering (4.5), (4.7), and (4.8), we have that $MS(\hat{v}_k, Q)$ is bounded. Therefore, appealing to the compactness result [2, Theorem 4.8], we deduce that (\hat{v}_k) is pre-compact in $L^1(Q)$. Since by definition of \hat{v}_k

$$\tilde{v}_k = \begin{cases} \hat{v}_k & \text{in } A \\ v_k & \text{in } Q \setminus A, \end{cases}$$

we have that also (\tilde{v}_k) is pre-compact in $L^1(Q)$. Further, since $\tilde{v}_k = v_k$ on $T_{\varepsilon_k/2}^j \pm (0, \varepsilon_k/2)$, invoking Lemma 2 gives $\tilde{v}_k \rightarrow u_t$ in $L^1(Q)$.

Step 3. We now introduce the final sequence (w_k) as a further modification of the sequence (\tilde{v}_k) ; specifically, we set $w_k := \tilde{v}_k$ in A and define it in $Q \setminus A$ as explained below. We start observing that since

$$\mathcal{L}^2(Q \setminus A) = \varepsilon_k \#\mathcal{J}_k^c \leq \frac{3\varepsilon_k}{\beta} \rightarrow 0,$$

this last modification will keep the limit u_t . Therefore we only need to check that we can also control the energy along the new sequence in terms of $F_{\varepsilon_k}(u_k, Q)$.

We start defining w_k in L_k^1 and $L_k^{n_k}$. To this end let $m_k := \min \mathcal{J}_k$ and $M_k := \max \mathcal{J}_k$ and set $w_k := a_k^{m_k}$ in L_k^1 and $w_k := b_k^{M_k}$ in $L_k^{n_k}$. Since in this way we have not added any new discontinuities, this first modification of \tilde{v}_k does not increase the energy.

Consider now an intermediate strip L_k^i , $i \in \{2, \dots, n_k - 1\}$. Observe that here $\tilde{v}_k = v_k$. By definition, L_k^i is adjacent to two strips $T_{\varepsilon_k}^j$ and $T_{\varepsilon_k}^{j+h}$ for some $j, j+h \in \mathcal{J}_k$ (see Figure 8(i)). If $\mathcal{H}^1(S_{v_k} \cap L_k^i) < 1/2$, we argue as in Step 2 and define w_k in L_k^i as the affine interpolation between the values b_k^j and a_k^{j+h} , obtaining

$$\frac{1}{4} \int_{L_k^i} |\nabla w_k|^2 dx \leq \int_{L_k^i \cap \varepsilon_k \hat{P}} |\nabla v_k|^2 dx. \quad (4.9)$$

If $\mathcal{H}^1(S_{v_k} \cap L_k^i) \geq 1/2$, we simply set $w_k := v_k$ in L_k^i . We notice that if there are two or more strips L_k^i such that $\mathcal{H}^1(S_{v_k} \cap L_k^i) \geq 1/2$, then

$$\int_{Q_k} |\nabla u_k|^2 dx + \mathcal{H}^1(S_{u_k} \cap Q) \geq 1$$

and the lower bound (4.3) holds true. On the other hand, there must be at least one of such strips; otherwise, w_k would be piecewise affine (thus Sobolev) in the whole Q , so in view of (4.5), (4.7), and (4.9) we should get

$$\frac{\alpha}{4} \int_Q |\nabla w_k|^2 dx \leq F_{\varepsilon_k}(u_k, Q),$$

and then deduce that (w_k) converges weakly in $H^1(Q)$, thus violating the $L^1(Q)$ -convergence to u_t . Therefore, we can assume that there is precisely one index i^* for which $\mathcal{H}^1(S_{v_k} \cap L_k^{i^*}) \geq 1/2$.

Summarising, in this step we have defined a sequence (w_k) that is piecewise affine in $Q \setminus L_k^{i^*}$ (see Figure 8(ii)). In particular $w_k = b_k^j$ in $T_{\varepsilon_k/2}^j + (0, \varepsilon_k/2)$ and $w_k = a_k^{j+h}$ in $T_{\varepsilon_k/2}^{j+h} - (0, \varepsilon_k/2)$. Moreover, by construction $w_k \rightarrow u_t$ strongly in $L^1(Q)$.

Step 4. In this step we show that the energy of w_k concentrates in the strip $L_k^{i^*}$ and provides us with a lower bound on $F_{\varepsilon_k}(u_k, Q)$ as in (4.3).

To this end, we start observing that since $\hat{P} \subset P$ and $w_k = v_k$ in $L_k^{i^*}$, for k large enough (so that $\varepsilon_k < \alpha$), by (4.5) we have

$$F_{\varepsilon_k}(u_k, Q) \geq \alpha \int_{L_k^{i^*} \cap \varepsilon_k \hat{P}} |\nabla w_k|^2 dx + \varepsilon_k \int_{L_k^{i^*} \setminus \varepsilon_k \hat{P}} |\nabla w_k|^2 dx + \mathcal{H}^1(S_{w_k} \cap L_k^{i^*}).$$

Since $j+1, \dots, j+h-1 \in \mathcal{J}_k^c$ (see Figure 8(ii)), similarly to (4.8), we get

$$h-1 \leq \frac{2\mathcal{H}^1(S_{w_k} \cap Q)}{\beta} \leq \frac{3}{\beta}. \quad (4.10)$$

Moreover, since $w_k \rightarrow u_t$ strongly in $L^1(Q)$ and outside the strip $L_k^{i^*}$ the function w_k is piecewise affine with gradient uniformly bounded in $L^2(Q; \mathbb{R}^2)$, necessarily $L_k^{i^*}$ converges to $\{x_2 = 0\} \cap Q$ in the sense of the Hausdorff distance and $b_k^j - a_k^{j+h} \rightarrow t$ as $k \rightarrow +\infty$; hence in particular

$$|b_k^j - a_k^{j+h}| \geq \frac{t}{2}, \quad (4.11)$$

for k large enough. Comparing the bulk energy of w_k in $L_k^{i^*} \cap \varepsilon_k \hat{P}$ with that of the affine interpolation between the values b_k^j and a_k^{j+h} we get

$$\int_{L_k^{i^*} \cap \varepsilon_k \hat{P}} |\nabla w_k|^2 dx \geq \left(\frac{3}{4} - \mathcal{H}^1(S_{w_k} \cap L_k^{i^*} \cap \varepsilon_k \hat{P})\right) \frac{|b_k^j - a_k^{j+h}|^2}{\varepsilon_k(h-1)}. \quad (4.12)$$

Since

$$\varepsilon_k(h-1) = \mathcal{L}^2(L_k^{i^*}) \leq \mathcal{L}^2(Q \setminus A) \rightarrow 0,$$

gathering (4.4), (4.11), and (4.12) we have

$$\liminf_{k \rightarrow +\infty} \mathcal{H}^1(S_{w_k} \cap L_k^{i^*} \cap \varepsilon_k \hat{P}) \geq \frac{3}{4}. \quad (4.13)$$

We now estimate the energy in $L_k^{i^*} \setminus \varepsilon_k \hat{P}$. When we come to minimise

$$\hat{F}(w_k) := \varepsilon_k \int_{L_k^{i^*} \setminus \varepsilon_k \hat{P}} |\nabla w_k|^2 dx + \mathcal{H}^1(S_{w_k} \cap (L_k^{i^*} \setminus \varepsilon_k \hat{P})),$$

by an easy one-dimensional argument we find that in each connected component of $L_k^{i^*} \setminus \varepsilon_k \hat{P}$ the minimum is attained either at the affine interpolation between the two constants, or at a

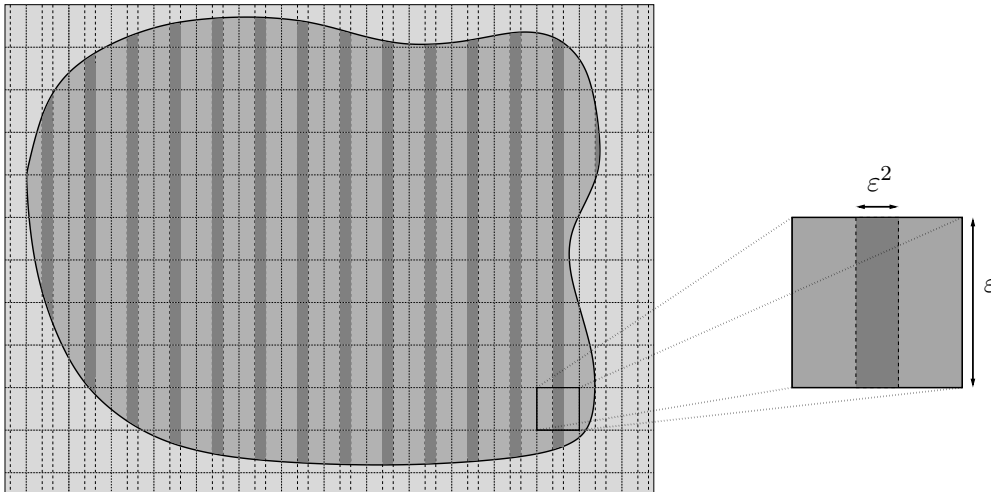


FIGURE 9. In darker grey the reinforcement $\Omega \setminus \varepsilon D_\varepsilon$.

function that is piecewise constant, takes values in $\{b_k^j, a_k^{j+h}\}$, and jumps along a horizontal line. Computing the energy of these two competitors leads to the estimate

$$\hat{F}(w_k) \geq \min \left\{ \frac{|b_k^j - a_k^{j+h}|^2}{4(h-1)}, \frac{1}{4} \right\}. \quad (4.14)$$

Hence, finally, combining (4.10)–(4.14) gives

$$\liminf_k F_{\varepsilon_k}(u_k) \geq \frac{3}{4} + \min \left\{ \frac{\beta}{48} t^2, \frac{1}{4} \right\} = \min \left\{ \frac{3}{4} + \frac{\beta}{48} t^2, 1 \right\}$$

and thus (4.3) with $c_3 = \beta/48$.

We now remove the assumption that ε_k^{-1} is an odd integer. To this end, let m_k be the largest integer such that $(2m_k + 1)\varepsilon_k \leq 1$. Since $Q_{(2m_k+1)\varepsilon_k}$ can be decomposed into periodicity cells entirely contained in Q , we can repeat in $Q_{(2m_k+1)\varepsilon_k}$ the construction leading to the sequence (w_k) . Finally, we extend w_k by reflection in the vanishing frame $Q \setminus Q_{(2m_k+1)\varepsilon_k}$ and conclude. \square

5. HOMOGENISATION OF A FIBRE-REINFORCED BRITTLE COMPOSITE

The object of this last section is the homogenisation of a laminar composite made of a brittle constituent reinforced by periodically distributed tougher fibres (see Figure 9). Although from a modelling point of view this problem is quite different from the one analysed in the previous sections, we show that the same mathematical techniques apply also in this case (with minor modifications) and lead again to a homogenised model of cohesive type.

Let Ω be a bounded open subset of \mathbb{R}^2 with Lipschitz boundary. For any $\varepsilon > 0$ we consider the energy-functional $E_\varepsilon: L^1(\Omega) \rightarrow [0, +\infty]$ defined as

$$E_\varepsilon(u) := \begin{cases} \int_\Omega |\nabla u|^2 dx + \mathcal{H}^1(S_u \cap \Omega \cap \varepsilon D_\varepsilon) + \frac{1}{\varepsilon} \mathcal{H}^1(S_u \cap (\Omega \setminus \varepsilon D_\varepsilon)) & \text{if } u \in SBV^2(\Omega), \\ +\infty & \text{otherwise in } L^1(\Omega), \end{cases} \quad (5.1)$$

where

$$D_\varepsilon := \mathbb{R}^2 \setminus \bigcup_{i \in \mathbb{Z}} (A_\varepsilon + (i, 0)), \quad A_\varepsilon := \left(-\frac{\varepsilon}{2}, \frac{\varepsilon}{2}\right) \times \mathbb{R}.$$

Here, $\Omega \setminus \varepsilon D_\varepsilon$ (constituted by vertical strips of thickness ε^2) models the so-called *fibre-reinforcement*. Therefore, unlike the previous case, the two constituents $\Omega \cap \varepsilon D_\varepsilon$ and $\Omega \setminus \varepsilon D_\varepsilon$ now differ in their toughness, while they share (for simplicity) the same elastic modulus, equal to one; however, the same qualitative behaviour described below is retrieved when the two constituents have different elastic moduli (both independent of ε).

We notice that following [4, 21] one could also model the tough fibres as a constituent with infinite toughness (thus imposing a restriction on the location of the jump set of the admissible displacements). We rather find less artificial to model the reinforcement as a material whose toughness diverges as ε tends to zero. Further, in order to obtain a nontrivial limit model we require that the toughness is equal to the inverse of the length-scale of the microstructure. Indeed, different choices lead to a Γ -limit where the surface energy density either is identically equal to one or grows linearly with the crack-opening. (This can be seen for instance arguing as in the estimate from above in the proof of Theorem 5.) When the toughness equals $1/\varepsilon$ we prove instead that the surface energy density depends on the crack-opening and becomes constant when the crack-opening is large (see (5.4)).

The following result holds true.

Theorem 5. *For every decreasing sequence of positive numbers converging to zero, there exists a subsequence (ε_k) such that (E_{ε_k}) Γ -converges to a functional $E: L^1(\Omega) \rightarrow [0, +\infty]$ of the form*

$$E(u) := \begin{cases} \int_{\Omega} |\nabla u|^2 dx + \int_{S_u} g([u], \nu_u) d\mathcal{H}^1 & \text{if } u \in GSBV^2(\Omega), \\ +\infty & \text{otherwise in } L^1(\Omega), \end{cases} \quad (5.2)$$

where $g: \mathbb{R} \times \mathbb{S}^1 \rightarrow [0, +\infty)$ is a Borel function satisfying the following properties:

(i) for any fixed $\nu \in \mathbb{S}^1$, $g(\cdot, \nu)$ is nondecreasing in $(0, +\infty)$ and satisfies the symmetry condition $g(-t, -\nu) = g(t, \nu)$; moreover, for every $t \neq 0$ and $\nu \in \mathbb{S}^1$

$$1 \leq g(t, \nu) \leq 2; \quad (5.3)$$

(ii) for every $t > 0$

$$\min\{1 + ct^2, 2\} \leq g(t, e_2) \leq \min\{1 + 2\sqrt{2}t, 2\}, \quad (5.4)$$

for some $c > 0$. Hence in particular $g(0^+, e_2) = 1$ and $g(t, e_2) = 2$ for t larger than some $t_0 > 0$.

We remark that (5.1) shares some similarities with the energy-functionals considered in [4, 21] where the authors analyse the asymptotic behaviour of a brittle material reinforced by a reticulated (connected) unbreakable structure, modelled as a constituent with infinite toughness. In [4, Theorem 3.1] and [21, Theorem 2.2] it is shown that a composite as above gives rise to a homogenised model of cohesive type (with an activation threshold). Moreover, the effect of having (suitably chosen) regions with infinite toughness is that the limit surface energy density grows linearly in the crack-opening, which is though unsatisfactory from a physical point of view. In this perspective, Theorem 5 provides us with a homogenised model which is physically more relevant than the one obtained in [4, 21].

The proof of Theorem 5 closely follows those of Theorems 1 and 2, for this reason in the present section we give just a sketch of this proof detailing only the parts which are substantially different.

Γ -convergence and integral representation. The Γ -convergence of the functionals E_ε can be derived also in this case by combining the localisation method and the integral representation result [7, Theorem 1]. The main difference with the previous case is that now we trivially have $MS \leq E_\varepsilon$ while E_ε does not satisfy an estimate from above in terms of the Mumford-Shah

functional. Hence, to fully fit the assumptions needed to apply the localisation method of Γ -convergence, we set $E'' := \Gamma\text{-lim sup}_{\varepsilon \rightarrow 0} E_\varepsilon$ and prove that $E'' \leq 2MS$. To this end, it is convenient to introduce the following notation.

Let us fix an open rectangle R containing $\bar{\Omega}$ and let $\mathcal{W}(R)$ be the space of all functions $w \in SBV^2(R) \cap L^\infty(R)$ enjoying the following properties:

- $S_w \subset L$, with L finite union of pairwise disjoint closed segments contained in R ;
- $w \in W^{1,\infty}(R \setminus L)$.

Moreover, we denote by $\mathcal{A}_L(\Omega)$ the class of all open subsets of Ω with Lipschitz boundary.

To obtain the desired estimate on E'' we need the following approximation lemma which is a consequence of [16, Theorem 3.9].

Lemma 6. *Let $U \in \mathcal{A}_L(\Omega)$ and let $u \in SBV^2(U) \cap L^\infty(U)$. Then u has an extension $v \in SBV^2(R) \cap L^\infty(R)$ with compact support in R , such that*

$$\mathcal{H}^1(S_v \cap \partial U) = 0,$$

and $\|v\|_{L^\infty(R)} = \|u\|_{L^\infty(U)}$. Moreover, there exist a sequence $(w_j) \subset \mathcal{W}(R)$ converging to v in $L^1(R)$, and a sequence (L_j) of finite unions of pairwise disjoint closed segments contained in R and such that $S_{w_j} \subset L_j$, with the following properties:

$$\|w_j\|_{L^\infty(R)} \leq \|v\|_{L^\infty(R)} = \|u\|_{L^\infty(U)},$$

$$\nabla w_j \rightarrow \nabla v \quad \text{strongly in } L^2(R; \mathbb{R}^2), \text{ hence } \nabla w_j \rightarrow \nabla u \quad \text{strongly in } L^2(U; \mathbb{R}^2),$$

$$\mathcal{H}^1(L_j \cap U) \rightarrow \mathcal{H}^1(S_v \cap U) = \mathcal{H}^1(S_u \cap U).$$

Proof. To prove the first assertion we can use locally a reflection argument in a curvilinear coordinate system for which the boundary is flat. The global extension can be obtained, as usual, through a partition of unity. Then, the existence of the approximating sequence (w_j) is a consequence of the density result [16, Theorem 3.9]. \square

We are in a position to prove the following proposition.

Proposition 1. *For every $u \in SBV^2(\Omega) \cap L^\infty(\Omega)$ and for every $U \in \mathcal{A}_L(\Omega)$ we have*

$$E''(u, U) \leq \int_U |\nabla u|^2 dx + 2\mathcal{H}^1(S_u \cap U). \quad (5.5)$$

Proof. We fix $U \in \mathcal{A}_L(\Omega)$; in view of Lemma 6 and of the locality of E'' it is enough to prove

$$E''(u, U) \leq \int_U |\nabla u|^2 dx + 2\mathcal{H}^1(L \cap U),$$

for $u \in \mathcal{W}(R)$.

We want to construct a sequence $(u_k) \subset SBV^2(U)$ converging to u in $L^1(\Omega)$ and such that

$$\limsup_{k \rightarrow +\infty} E_{\varepsilon_k}(u_k, U) \leq \int_U |\nabla u|^2 dx + 2\mathcal{H}^1(L \cap U).$$

Since U has Lipschitz boundary, we can slightly modify u near each connected component of \bar{S}_u to find a L that intersects ∂U in a finite number of points. This can be done, for instance, by slightly shifting these connected components taking into account the area formula for ∂U .

Now we explicitly construct (u_k) when L is a single closed segment; then, the general case follows easily. Let $\nu \in \mathbb{S}^1$ denote the normal to L . We need to distinguish two cases.

Case 1: $\nu = e_1$; i.e., L is a vertical segment.

To fix the ideas let $c \in \mathbb{R}$ and suppose that $L \subset \{x_1 = c\}$. Let R' be an open rectangle such that $\overline{\Omega} \cup L \subseteq R' \subset\subset R$. In this case we obtain u_k as a suitable translation of u . The latter is needed in order to prevent the possibility that L entirely falls in a tough vertical fibre, thus paying too much in energy. To this end, for k large enough and for $x \in R'$ we set

$$u_k(x) := u\left(x + \left(c - \left\lfloor \frac{c}{\varepsilon_k} \right\rfloor \varepsilon_k - \varepsilon_k^2\right) e_1\right),$$

so that S_{u_k} is contained in the vertical line $x_1 = \lfloor \frac{c}{\varepsilon_k} \rfloor \varepsilon_k + \varepsilon_k^2$. Hence, $(u_k) \subset SBV^2(R')$, $S_{u_k} \cap \Omega \subseteq \Omega \cap \varepsilon_k D_{\varepsilon_k}$, u_k is bounded in $L^\infty(R')$, $u_k \rightarrow u$ in $L^1(\Omega)$. Moreover,

$$\begin{aligned} \lim_{k \rightarrow +\infty} E_k(u_k, U) &= \lim_{k \rightarrow +\infty} \int_U |\nabla u_k|^2 dx + \mathcal{H}^1(S_{u_k} \cap U \cap \varepsilon_k D_{\varepsilon_k}) \\ &\leq \int_U |\nabla u|^2 dx + \mathcal{H}^1(L \cap U), \end{aligned}$$

which gives the desired estimate in this case.

Case 2: $\nu \neq e_1$.

Let ϑ be the angle (measured anticlockwise) between e_1 and ν^\perp ; thanks to the symmetries of the problem, it is enough to consider the case $0 \leq \vartheta < \pi/2$.

Let $N_k \in \mathbb{N}$ be the number of tough fibres intersecting $L \cap U$; notice that since L intersects ∂U in a finite number of points, we have

$$\frac{\varepsilon_k}{\cos \vartheta} N_k \leq \mathcal{H}^1(L \cap U) + o(1) \quad \text{as } k \rightarrow +\infty. \quad (5.6)$$

Then, setting $u_k := u$ for every $k \in \mathbb{N}$, in view of (5.6) we find

$$\begin{aligned} &\lim_{k \rightarrow +\infty} E_k(u_k, U) \\ &= \lim_{k \rightarrow +\infty} \left(\int_U |\nabla u_k|^2 dx + \mathcal{H}^1(S_{u_k} \cap U \cap \varepsilon_k D_{\varepsilon_k}) + \frac{1}{\varepsilon_k} \mathcal{H}^1(S_{u_k} \cap (U \setminus \varepsilon_k D_{\varepsilon_k})) \right) \\ &= \int_U |\nabla u|^2 dx + \lim_{k \rightarrow +\infty} \left(N_k \frac{\varepsilon_k}{\cos \vartheta} + \frac{1}{\varepsilon_k} N_k \frac{\varepsilon_k^2}{\cos \vartheta} \right) \\ &\leq \int_U |\nabla u|^2 dx + 2 \mathcal{H}^1(L \cap U), \end{aligned}$$

and thus the thesis. \square

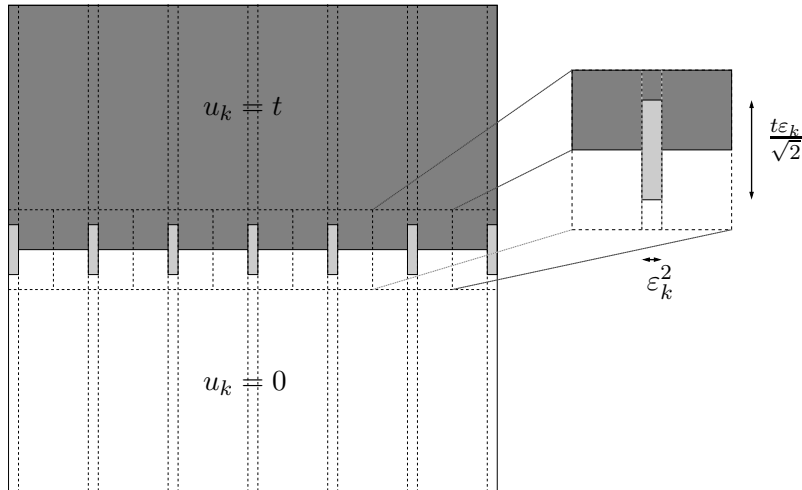
Thus by virtue of Proposition 1, arguing as in Section 3 we can deduce that E_ε Γ -converges to the functional E as in (5.2). We also remark that the bound from above in (5.3) follows from (5.5) choosing

$$u(x) = \begin{cases} t & \text{if } x \cdot \nu > 0 \\ 0 & \text{if } x \cdot \nu \leq 0 \end{cases} \quad \text{and } U = Q.$$

Dependence on the crack-opening. We now turn to the proof of (ii) Theorem 5. We start proving the estimate from above in (5.4).

Estimate from above. We want to show that $g(t, e_2) \leq 1 + 2\sqrt{2}t$. To this end, let $R_\varepsilon \subset A_\varepsilon$ be the open rectangle defined as

$$R_\varepsilon := \left(-\frac{\varepsilon}{2}, \frac{\varepsilon}{2}\right) \times \left(-\frac{t}{2\sqrt{2}}, \frac{t}{2\sqrt{2}}\right).$$

FIGURE 10. The optimal sequence u_k .

Set

$$R_k := Q \cap \varepsilon_k \left(\bigcup_{i \in \mathbb{Z}} R_{\varepsilon_k} + (i, 0) \right)$$

and let $(u_k) \subset SBV^2(Q)$ be the sequence of functions defined as

$$u_k(x) := \begin{cases} t & \text{if } x \in Q \setminus R_k \text{ and } x_2 \geq 0, \\ \frac{t}{2} + \frac{\sqrt{2}}{\varepsilon_k} x_2 & \text{if } x \in R_k, \\ 0 & \text{if } x \in Q \setminus R_k \text{ and } x_2 < 0, \end{cases}$$

(see Figure 10). We clearly have $u_k \rightarrow u_t$ in $L^1(Q)$; moreover

$$\int_{R_k} |\nabla u_k|^2 dx \leq \varepsilon_k \left(\lfloor \frac{1}{\varepsilon_k} \rfloor + 1 \right) (\sqrt{2} t) \quad \text{and} \quad \mathcal{H}^1(S_{u_k}) \leq \varepsilon_k \left(\lfloor \frac{1}{\varepsilon_k} \rfloor + 1 \right) (1 + \sqrt{2} t).$$

Since $S_{u_k} \subset Q \cap \varepsilon_k D_{\varepsilon_k}$, we readily deduce

$$g(t, e_2) = E(u_t, Q) \leq \limsup_{k \rightarrow +\infty} E_{\varepsilon_k}(u_k, Q) \leq 1 + 2\sqrt{2}t,$$

and hence the estimate from above.

Estimate from below. We now discuss the more delicate estimate from below in (5.4). Here we follow the steps of the proof of Theorem 2, from which we borrow some notation.

We notice that in this case the consecutive modifications of a sequence with equibounded energy are in general easier with respect to those in Theorem 2. For instance, since E_ε is bounded from below by the Mumford-Shah functional now we can readily deduce that the (analogues of the) two sequences (v_k) and (\tilde{v}_k) are precompact in $SBV^2(Q)$. On the other hand, in the present case, when we come to define (the analogous of) v_k we cannot directly apply Lemma 5. Indeed, to keep the control on the energy of v_k , now we need to avoid the introduction of new jumps falling in the reinforced fibres. For this reason we need a suitably improved variant of Lemma 5.

The following lemma is a slight modification of the patching lemma [4, Lemma 3.3], to which we refer for an idea of the proof, and it holds true in dimension two only.

Lemma 7. *There exist two constants $\alpha, \beta > 0$ with the following property: Let $u \in SBV^2(Q) \cap L^\infty(Q)$ be symmetric in the first variable and such that $\mathcal{H}^1(S_u \cap U) \leq \beta$, where $U := (-1/2, 1/2) \times (-3/8, 3/8)$. Then there exists a function $v \in SBV^2(Q) \cap L^\infty(Q)$ still symmetric and such that*

- (i) v is constant in $(-1/2, 1/2) \times (-1/4, 1/4)$;
- (ii) $v = u$ in $Q \setminus U$;
- (iii) $\|v\|_{L^\infty(Q)} \leq \|u\|_{L^\infty(Q)}$;
- (iv) $\alpha \int_Q |\nabla v|^2 dx \leq \int_Q |\nabla u|^2 dx$;
- (v) $S_v \subset S_u$.

On account of Lemma 7 we are now ready to show that for any given $(u_k) \subset SBV^2(Q)$ such that $u_k \rightarrow u_t$ in $L^1(Q)$, we have

$$\liminf_k E_{\varepsilon_k}(u_k) \geq \min\{1 + ct^2, 2\}, \quad (5.7)$$

for some $c > 0$.

We may assume that $\liminf_k E_{\varepsilon_k}(u_k) \leq 2$, otherwise there is nothing to prove. As in Theorem 2, we can additionally assume that the liminf is actually a limit, that $0 \leq u_k \leq t$ in Q , that ε_k^{-1} is an odd integer, and, in view of Lemma 3, that u_k is ε_k -periodic and symmetric in its first variable. Also the general strategy of the proof remains the same: we are going to modify (u_k) obtaining two intermediate sequences (v_k) and (\tilde{v}_k) , and then finally the more regular sequence (w_k) whose energy concentrates in a horizontal layer of thickness proportional to ε_k .

Let $R_{\varepsilon_k}^{i,j} := \varepsilon_k R + p_k^{i,j}$, where $R := (-1/2, 1/2) \times (-7/8, 7/8)$, and let $\mathcal{J}_k \subset \{2, \dots, \varepsilon_k^{-1} - 1\}$ be the set of indices j such that

$$\mathcal{H}^1(S_{u_k} \cap R_{\varepsilon_k}^{1,j}) \leq \beta.$$

Therefore, proceeding as in Theorem 2 Step 1, we may now appeal to Lemma 7 to define a new sequence (v_k) without introducing any new jumps, and thus keeping the control on the energy. Moreover, since $E_\varepsilon \geq MS$ the sequence (v_k) is precompact in $SBV^2(Q)$.

Then, arguing as in Theorem 2 Step 2, we further modify (v_k) constructing a new sequence (\tilde{v}_k) . We also get the analogue of (4.7), with the only difference that now we take into account all of the vertical one-dimensional slices and not only a subset of them.

Afterwards, we modify (\tilde{v}_k) in those strips L_k^i such that $\mathcal{H}^1(S_{v_k} \cap L_k^i) < 3/4$ thus obtaining the final sequence (w_k) . The latter is defined replacing \tilde{v}_k by a suitably chosen affine function, as in Theorem 2 Step 3. As before, the strong $L^1(Q)$ -convergence to u_t is preserved.

We notice that in this case there is at least one strip L_k^i such that $\mathcal{H}^1(S_{v_k} \cap L_k^i) \geq 3/4$ (otherwise w_k would be a Sobolev sequence with gradient uniformly bounded in $L^2(Q; \mathbb{R}^2)$, thus violating the convergence to u_t). Moreover, if there are three or more of such strips, then $\mathcal{H}^1(S_{u_k} \cap Q) \geq 2$ and the lower bound (5.7) holds trivially true. We are then left with the case where there are one or two strips such that $\mathcal{H}^1(S_{v_k} \cap L_k^i) = \mathcal{H}^1(S_{w_k} \cap L_k^i) \geq 3/4$. Since $w_k \rightarrow u_t$ strongly in $L^1(Q)$, at least in one of the two strips as above we have

$$|b_k^j - a_k^{j+h}| \geq \frac{t}{3},$$

for k large enough. We denote this strip by $L_k^{i^*}$. As in Theorem 2 Step 4, we now estimate from below $E_{\varepsilon_k}(u_k, Q)$ with the energy of w_k in $L_k^{i^*}$. We find

$$E_{\varepsilon_k}(u_k, Q) \geq \alpha \int_{L_k^{i^*}} |\nabla w_k|^2 dx + \mathcal{H}^1(S_{w_k} \cap L_k^{i^*} \cap \varepsilon_k D_{\varepsilon_k}) + \frac{1}{\varepsilon_k} \mathcal{H}^1((S_{w_k} \cap L_k^{i^*}) \setminus \varepsilon_k D_{\varepsilon_k}),$$

where α is the constant given by Lemma 7.

Comparing the bulk energy of w_k in $L_k^{i*} \cap \varepsilon_k D_{\varepsilon_k}$ with that of the affine interpolation between the values b_k^j and a_k^{j+h} , and using the same argument as in (4.12) we obtain

$$\liminf_{k \rightarrow +\infty} \mathcal{H}^1(S_{w_k} \cap L_k^{i*} \cap \varepsilon_k D_{\varepsilon_k}) \geq 1.$$

On the other hand, comparing the energy of w_k in $L_k^{i*} \setminus \varepsilon_k D_{\varepsilon_k}$ with those of the affine interpolation and of the piecewise constant function taking values in $\{b_k^j, a_k^{j+h}\}$ we get

$$\alpha \int_{L_k^{i*} \setminus \varepsilon_k D_{\varepsilon_k}} |\nabla w_k|^2 dx + \frac{1}{\varepsilon_k} \mathcal{H}^1((S_{w_k} \cap L_k^{i*}) \setminus \varepsilon_k D_{\varepsilon_k}) \geq \min \left\{ \frac{\alpha |b_k^j - a_k^{j+h}|^2}{(h-1)}, 1 \right\}.$$

Gathering the previous inequalities and the analogue of (4.10) gives

$$\liminf_k E_{\varepsilon_k}(u_k) \geq \min \left\{ 1 + \frac{\alpha\beta}{27} t^2, 2 \right\},$$

whence (5.7) follows. This concludes the proof of Theorem 5.

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(Marco Barchiesi) DIPARTIMENTO DI MATEMATICA ED APPLICAZIONI, UNIVERSITÀ DI NAPOLI “FEDERICO II”,
VIA CINTIA, 80126 NAPOLI, ITALY

E-mail address: barchies@gmail.com

(Giuliano Lazzaroni) SISSA, VIA BONOMEA 265, 34136 TRIESTE, ITALY

E-mail address: giuliano.lazzaroni@sissa.it

(Caterina Ida Zeppieri) INSTITUT FÜR NUMERISCHE UND ANGEWANDTE MATHEMATIK, UNIVERSITÄT MÜNSTER,
EINSTEINSTR. 62, 48149 MÜNSTER, GERMANY

E-mail address: caterina.zeppieri@uni-muenster.de